

Current Quantum Gravity Theories, Experimental Evidence, Philosophical Implications

carlo roveli
IHES 2017

1. We have quantum gravity theories

2. We have have empirical evidence

3. There is some understanding on some issues of major philosophical relevance

i: nature of physical space

ii: nature of physical time

iii: relation between physical and experiential time

iv: connection between relational aspects of GR and QM

1. We do have quantum gravity theories

Many directions of investigation

loop quantum gravity,
string theory,
asymptotic safety,
Hořava–Lifshitz theory,
supergravity,
AdS-CFT-like dualities
twistor theory,
causal set theory,
entropic gravity,
emergent gravity,
non-commutative geometry,
group field theory,
Penrose nonlinear quantum dynamics
causal dynamical triangulations,
shape dynamics,
't Hooft theory
non-quantization of geometry
...

Many directions of investigation

loop quantum gravity,
string theory,
asymptotic safety,
Hořava–Lifshitz theory,
supergravity,
AdS-CFT-like dualities
twistor theory,
causal set theory,
entropic gravity,
emergent gravity,
non-commutative geometry,
group field theory,
Penrose nonlinear quantum dynamics
causal dynamical triangulations,
shape dynamics,
't Hooft theory
non-quantization of geometry
...

Many directions of investigation

Vastly different numbers of researchers involved

loop quantum gravity,
string theory,
asymptotic safety,
Hořava–Lifshitz theory,
supergravity,
AdS-CFT-like dualities
twistor theory,
causal set theory,
entropic gravity,
emergent gravity,
non-commutative geometry,
group field theory,
Penrose nonlinear quantum dynamics
causal dynamical triangulations,
shape dynamics,
't Hooft theory
non-quantization of geometry
...

Many directions of investigation

Vastly different numbers of researchers involved

Most are highly incomplete

loop quantum gravity,
string theory,
asymptotic safety,
Hořava–Lifshitz theory,
supergravity,
AdS-CFT-like dualities
twistor theory,
causal set theory,
entropic gravity,
emergent gravity,
non-commutative geometry,
group field theory,
Penrose nonlinear quantum dynamics
causal dynamical triangulations,
shape dynamics,
't Hooft theory
non-quantization of geometry
...

Many directions of investigation

Vastly different numbers of researchers involved

Most are highly incomplete

Several are only vaguely connected
to the actual problem of quantum gravity

loop quantum gravity,
string theory,
asymptotic safety,
Hořava–Lifshitz theory,
supergravity,
AdS-CFT-like dualities
twistor theory,
causal set theory,
entropic gravity,
emergent gravity,
non-commutative geometry,
group field theory,
Penrose nonlinear quantum dynamics
causal dynamical triangulations,
shape dynamics,
't Hooft theory
non-quantization of geometry
...

Many directions of investigation

Vastly different numbers of researchers involved

Most are highly incomplete

Several are only vaguely connected
to the actual problem of quantum gravity

Several are related, boundaries are fluid

loop quantum gravity,
string theory,
asymptotic safety,
Hořava–Lifshitz theory,
supergravity,
AdS-CFT-like dualities
twistor theory,
causal set theory,
entropic gravity,
emergent gravity,
non-commutative geometry,
group field theory,
Penrose nonlinear quantum dynamics
causal dynamical triangulations,
shape dynamics,
't Hooft theory
non-quantization of geometry
...

Many directions of investigation

Vastly different numbers of researchers involved

Most are highly incomplete

Several are only vaguely connected
to the actual problem of quantum gravity

Several are related, boundaries are fluid

Many offer useful insights

loop quantum gravity,
string theory,
asymptotic safety,
Hořava–Lifshitz theory,
supergravity,
AdS-CFT-like dualities
twistor theory,
causal set theory,
entropic gravity,
emergent gravity,
non-commutative geometry,
group field theory,
Penrose nonlinear quantum dynamics
causal dynamical triangulations,
shape dynamics,
't Hooft theory
non-quantization of geometry
...

Many directions of investigation

Vastly different numbers of researchers involved

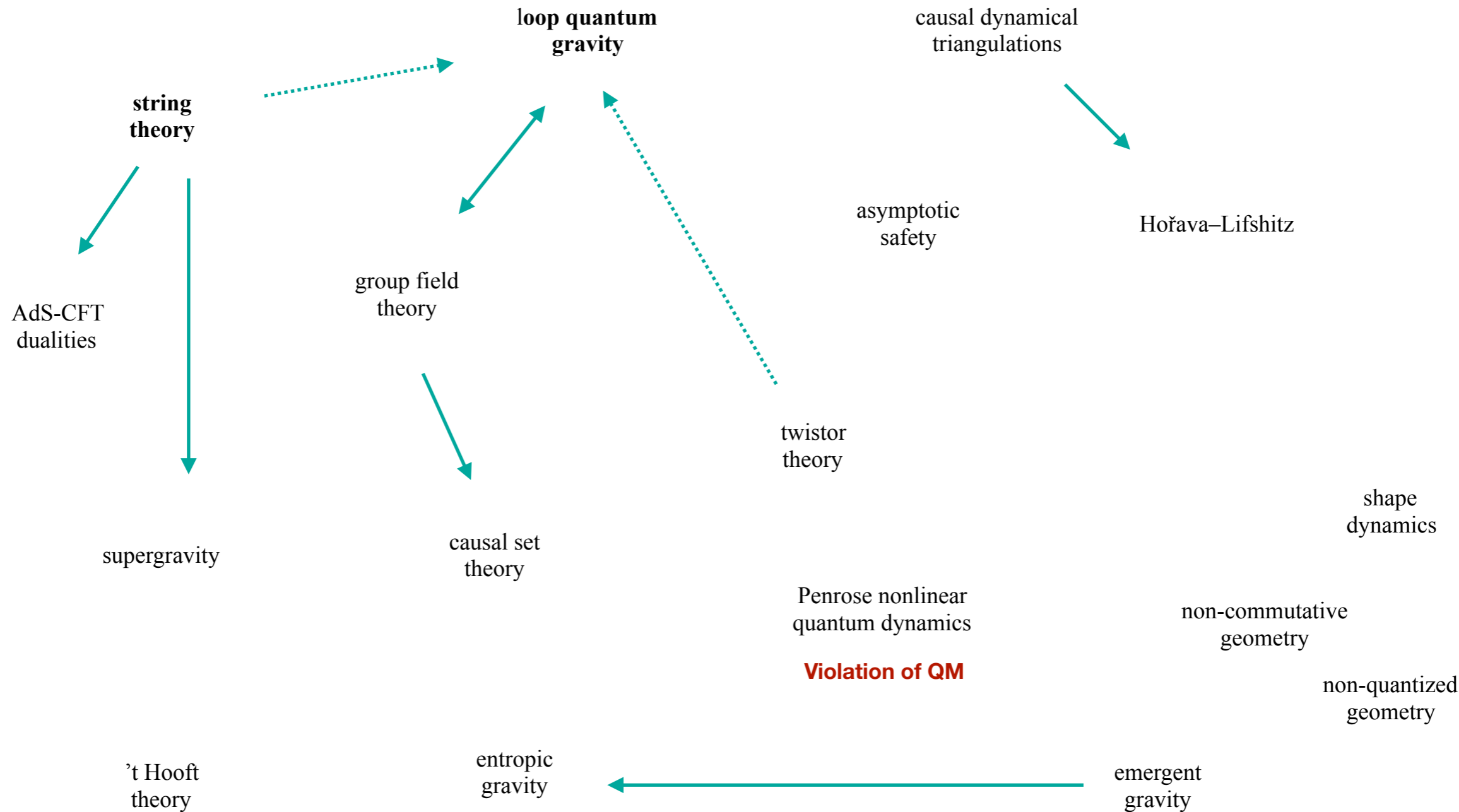
Most are highly incomplete

Several are only vaguely connected
to the actual problem of quantum gravity

Several are related, boundaries are fluid

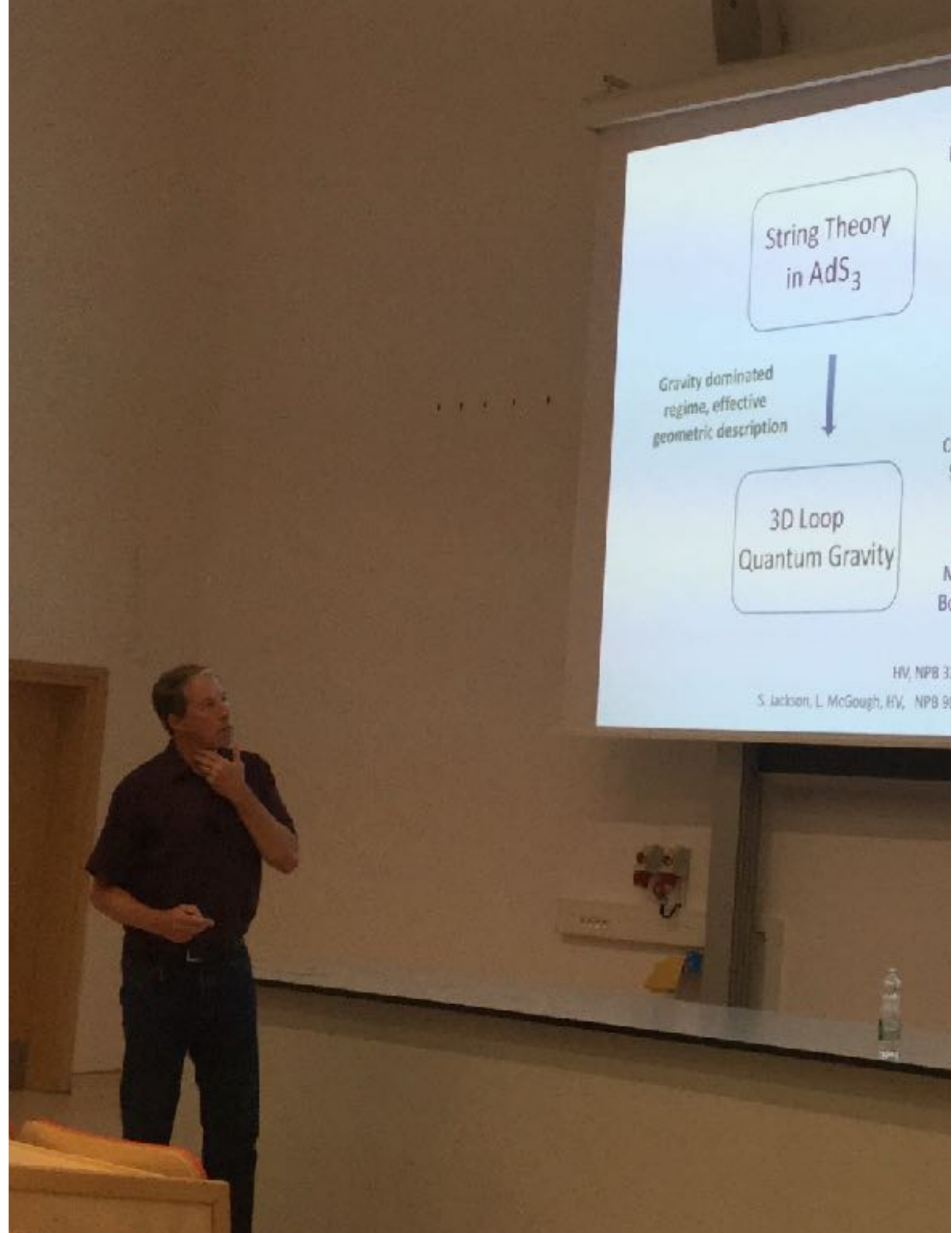
Many offer useful insights

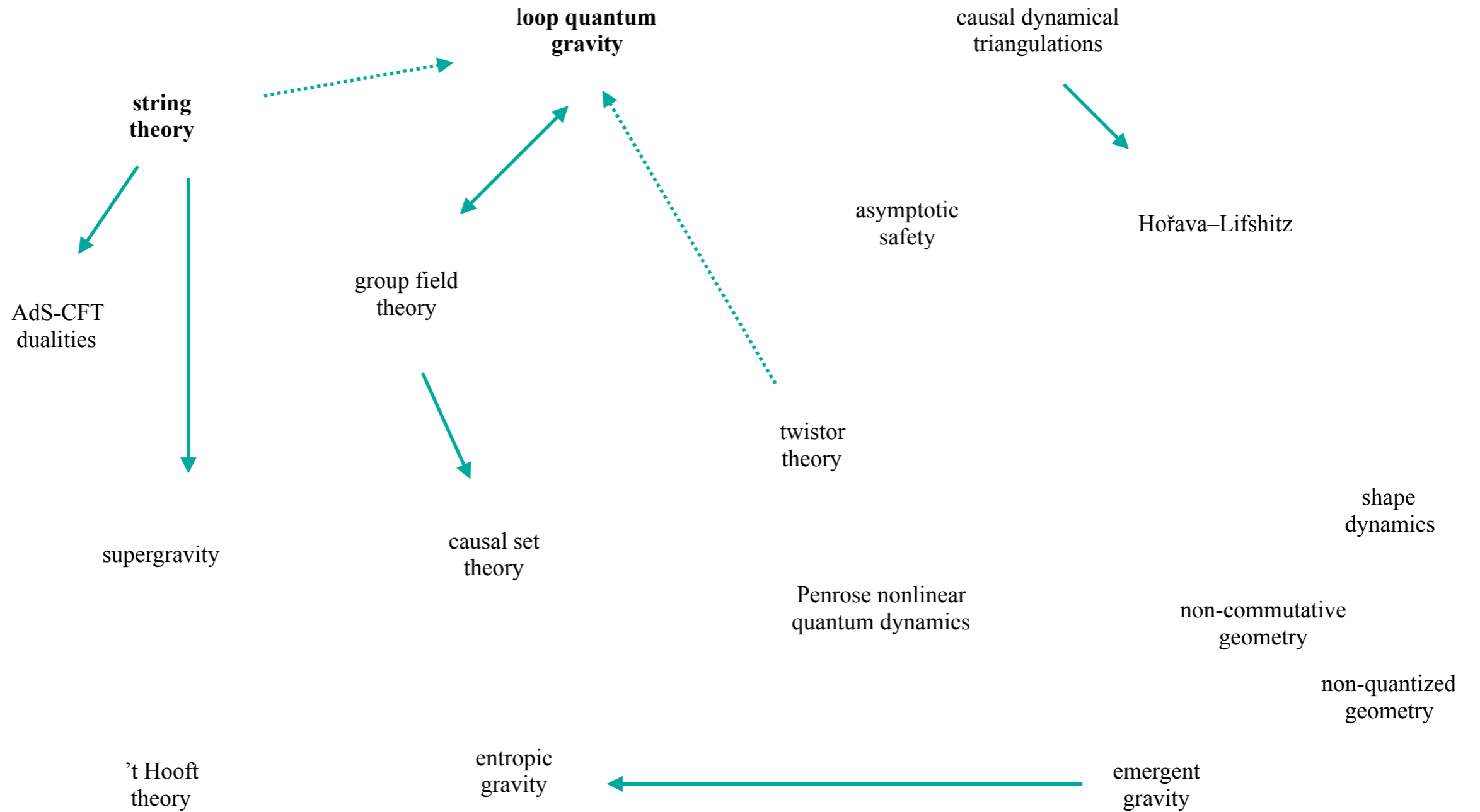
A few offer rather complete
tentative theories of quantum gravity

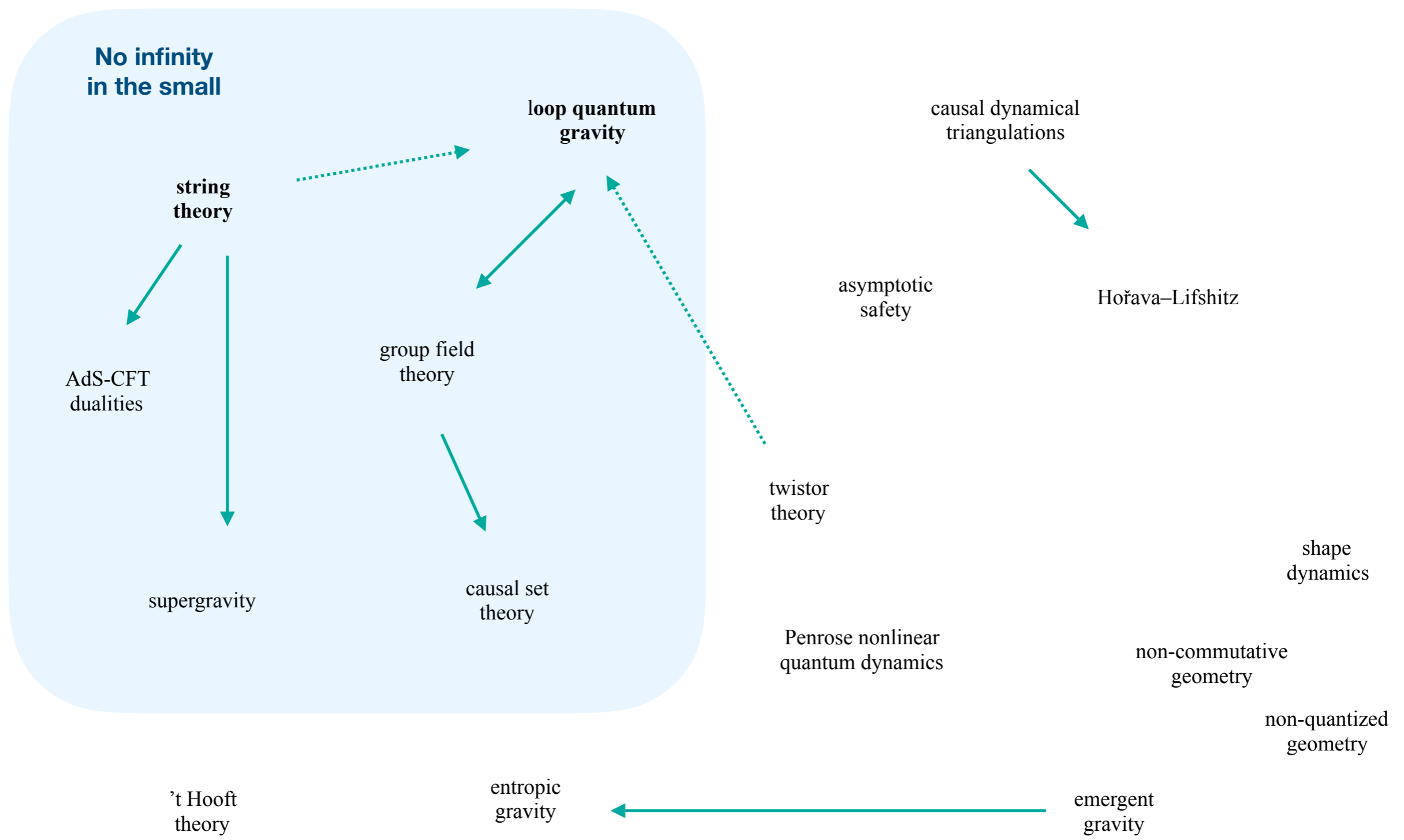


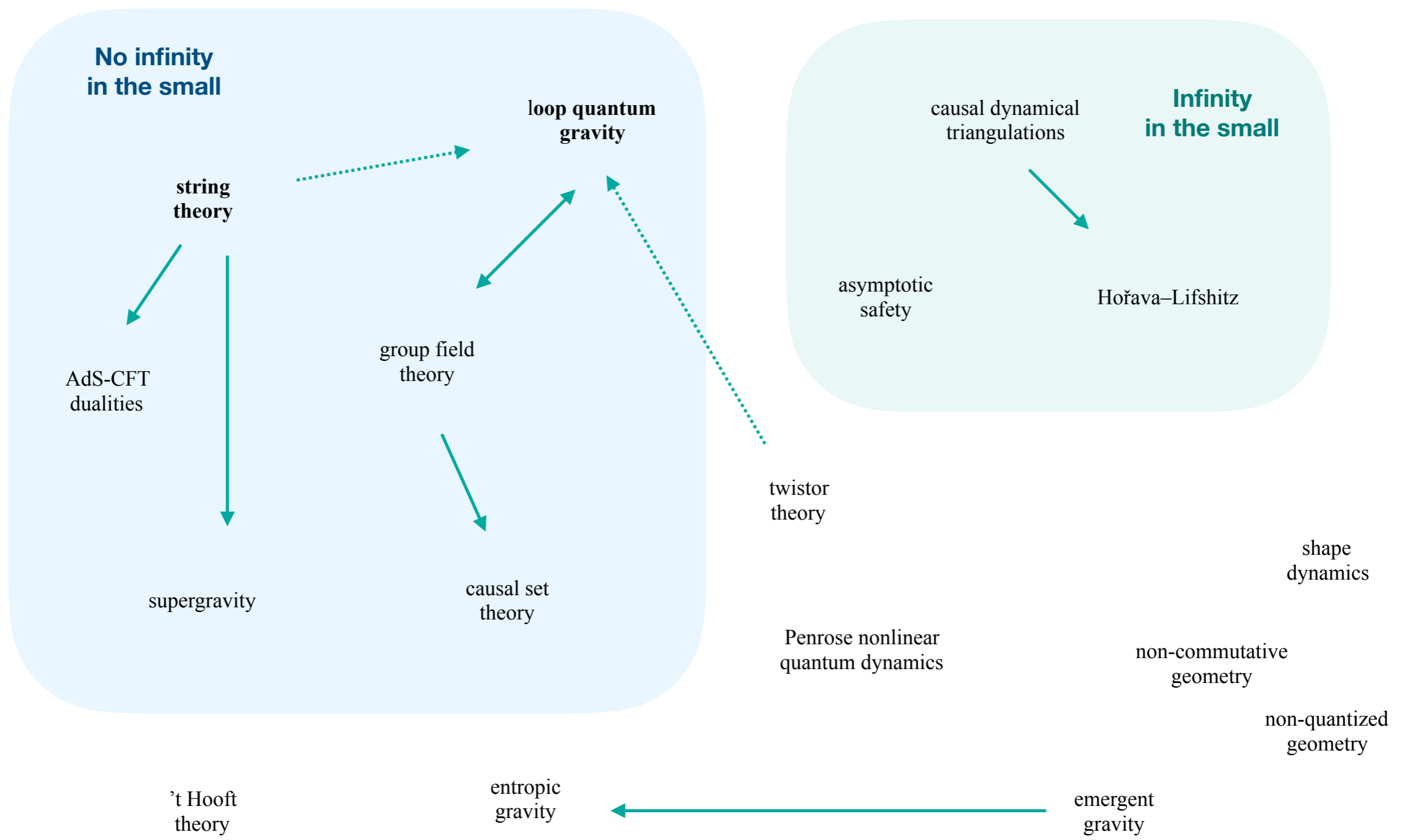
Several are related

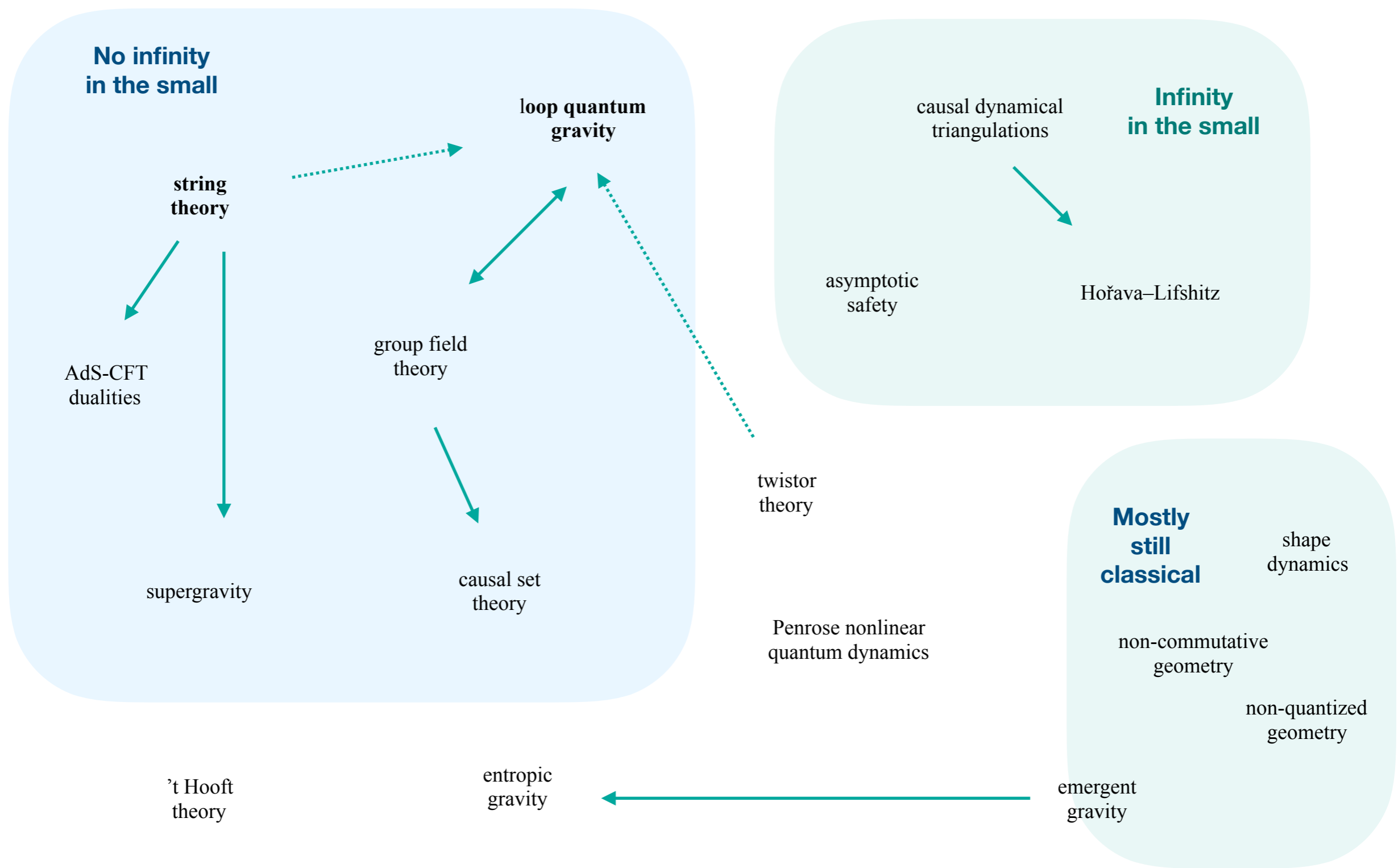
Herman Verlinde
at LOOP17 in Warsaw

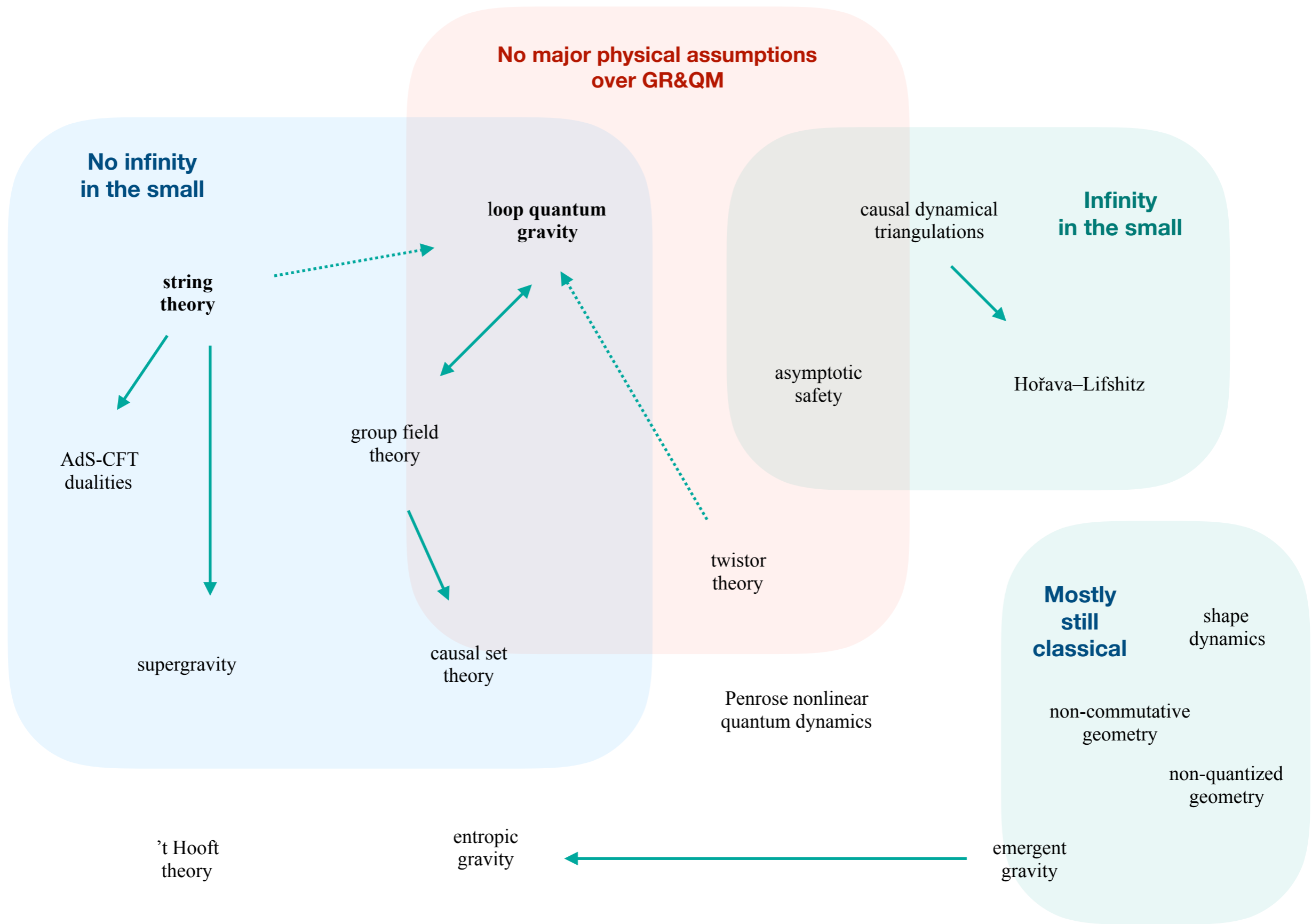












**No major physical assumptions
over GR&QM**

**No infinity
in the small**

**Infinity
in the small**

string
theory

loop quantum
gravity

causal dynamical
triangulations

AdS-CFT
dualities

group field
theory

asymptotic
safety

Hořava-Lifshitz

supergravity

causal set
theory

twistor
theory

**Mostly
still
classical**

shape
dynamics

Penrose nonlinear
quantum dynamics

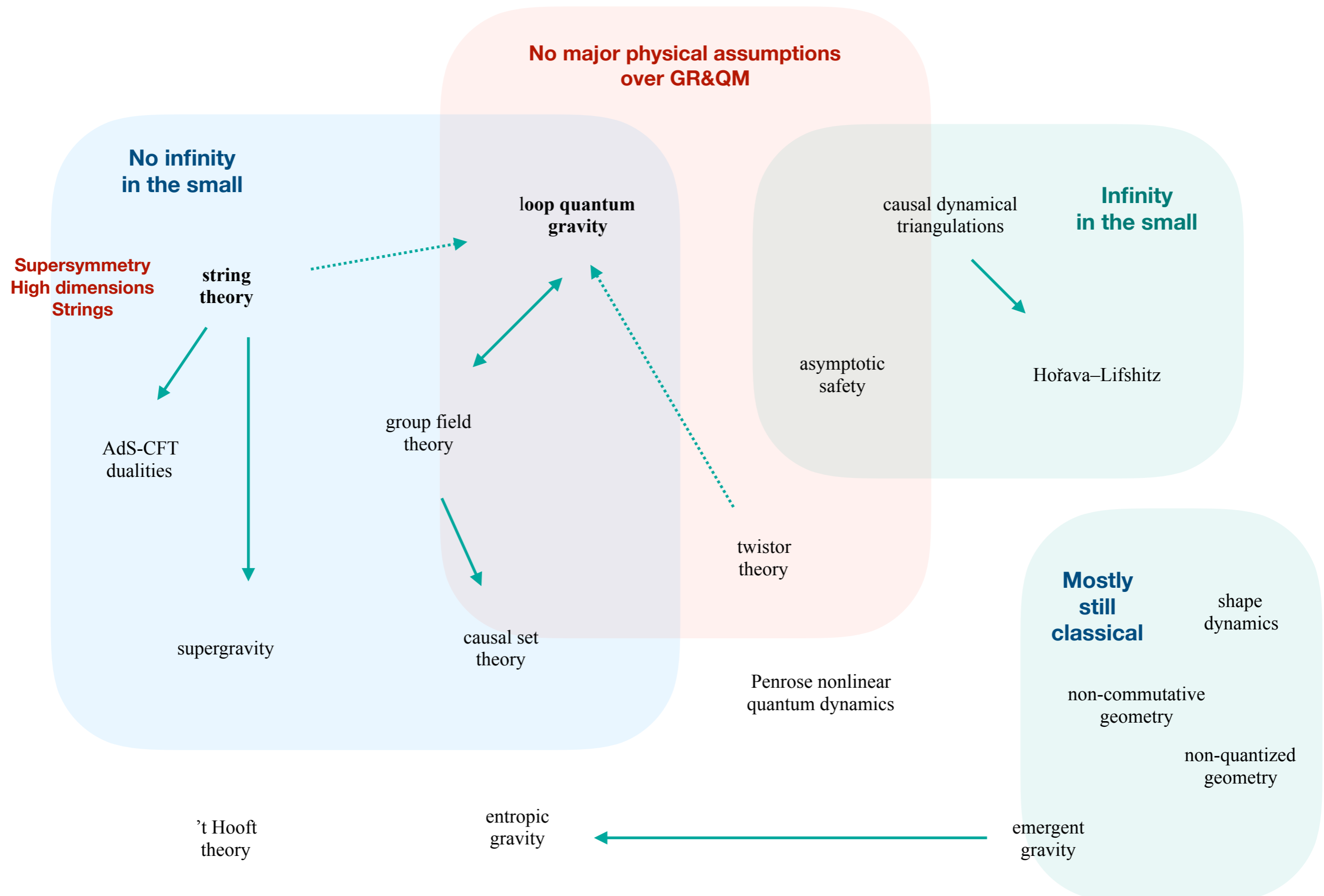
non-commutative
geometry

non-quantized
geometry

't Hooft
theory

entropic
gravity

emergent
gravity



**No major physical assumptions
over GR&QM**

**No infinity
in the small**

**Infinity
in the small**

**Supersymmetry
High dimensions
Strings**

string
theory

loop quantum
gravity

causal dynamical
triangulations

AdS-CFT
dualities

group field
theory

asymptotic
safety

Hořava-Lifshitz

supergravity

causal set
theory

twistor
theory

**Mostly
still
classical**

shape
dynamics

Penrose nonlinear
quantum dynamics

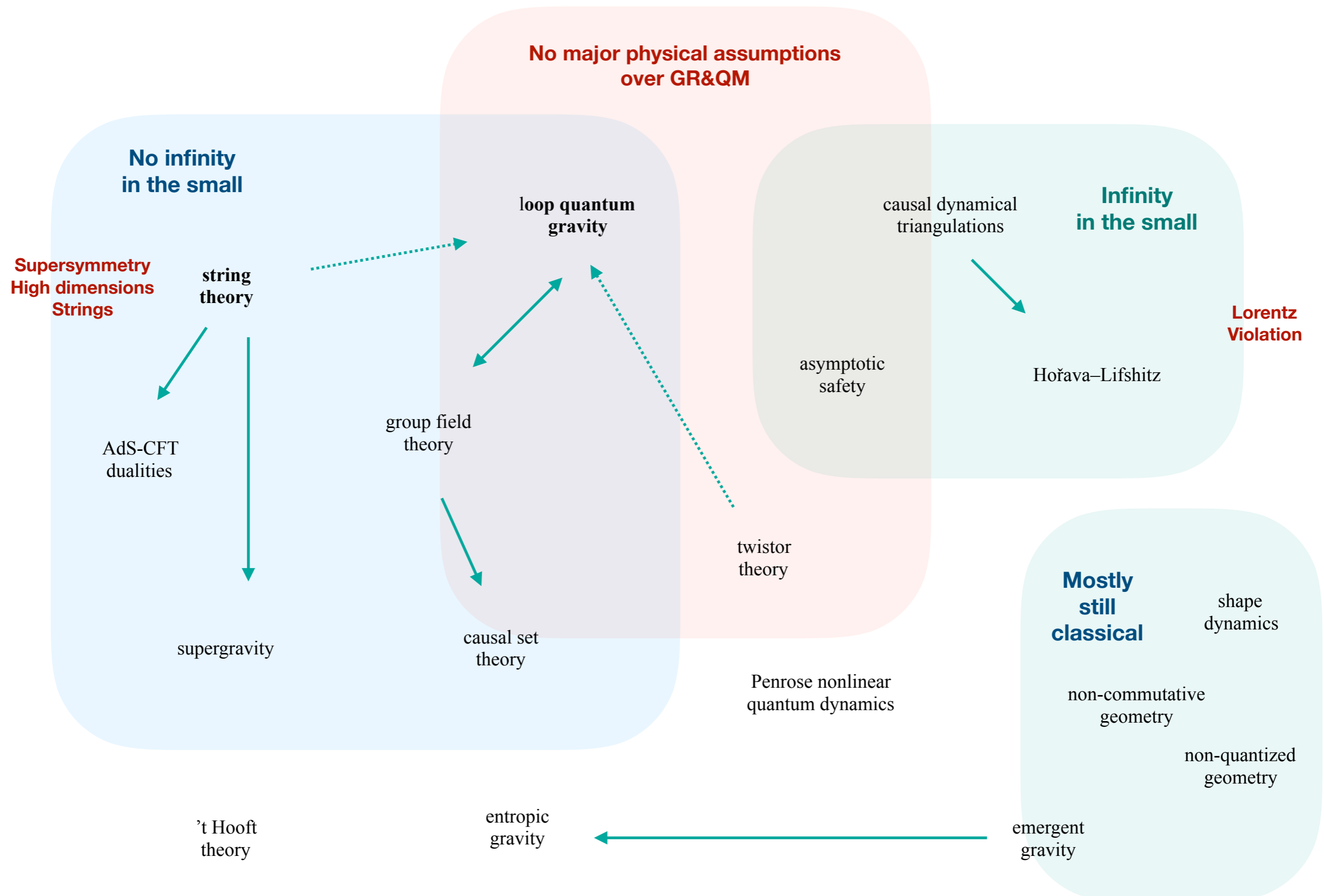
non-commutative
geometry

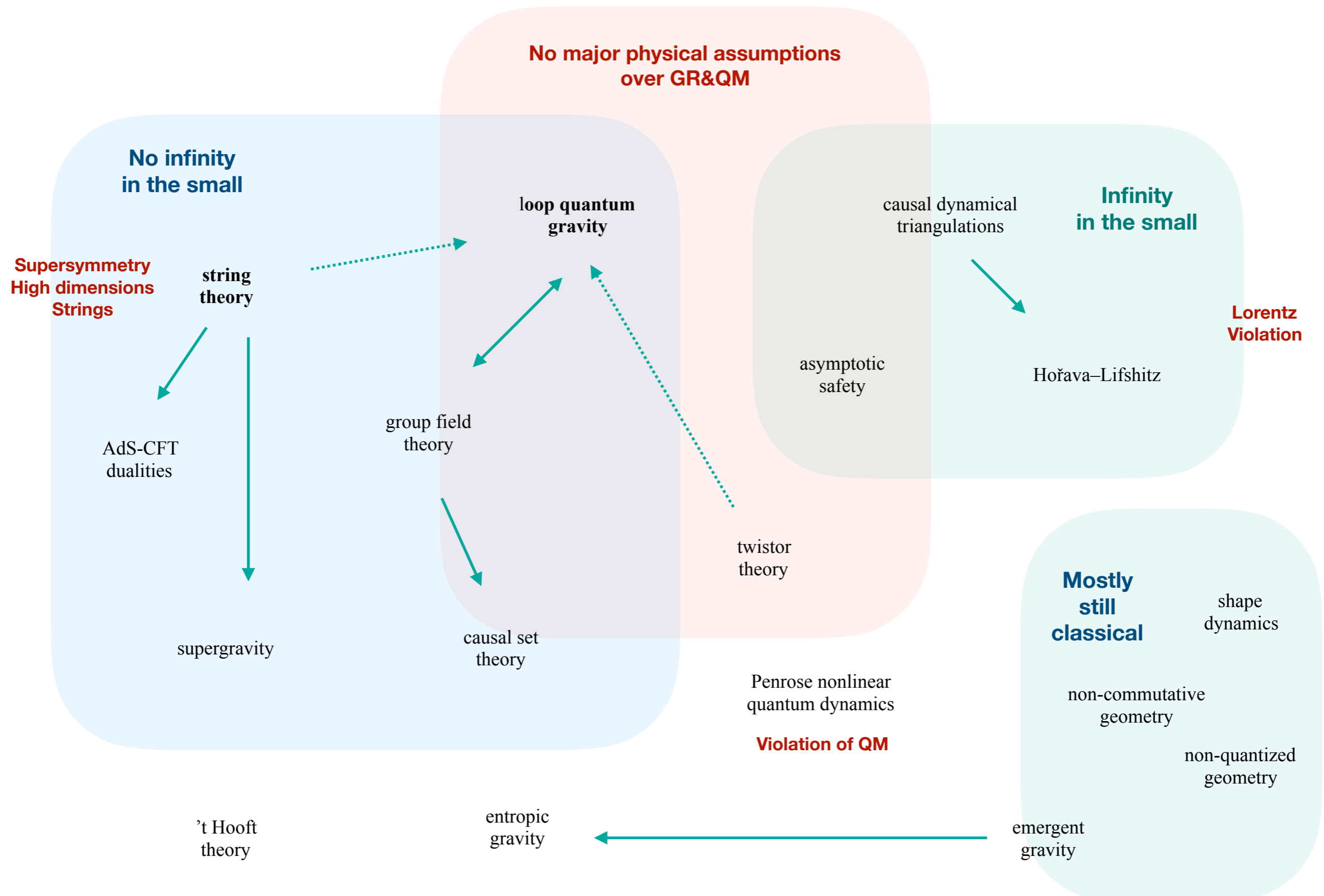
non-quantized
geometry

't Hooft
theory

entropic
gravity

emergent
gravity





Discriminatory questions:

Is **Lorentz symmetry** violated at the Planck scale or not?

Are there **supersymmetric** particles or not?

Is **Quantum Mechanics violated** in the presence of gravity or not?

Are there physical degrees of freedom at any **arbitrary small scale** or not?

Is geometry **discrete** i the small?

	Lorentz violations at Planck scale	Infinite d.o.f. at Planck scale	Supersymmetry	QM violations	Geometry is discrete?
Strings	No	No	Yes	No	No
Loops	No	No	No	No	Yes
Hojava Lifshitz	Yes	Yes	No	No	No
Asymptotic safety	No	Yes	No	No	No
Nonlinear quantum dynamics	No	Yes	No	Yes	No
Your favorite

2. We do have empirical evidence

Empirical evidence: 1: Lorentz invariance

Empirical evidence: 1: Lorentz invariance

Violation of Lorentz invariance → Renormalizability

Observation has already ruled out theories

$$E^2 = k^2 + c_n \frac{k^n}{m_{Pl}^{n-2}}$$

Order	photon	e^-/e^+	Protons	Neutrinos ^a
n=2	N.A.	$O(10^{-16})$	$O(10^{-20})$ (CR)	$O(10^{-8} \div 10^{-10})$
n=3	$O(10^{-16})$ (GRB)	$O(10^{-16})$ (CR)	$O(10^{-14})$ (CR)	$O(40)$
n=4	$O(10^{-8})$ (CR)	$O(10^{-8})$ (CR)	$O(10^{-6})$ (CR)	$O(10^{-7})^*$ (CR)

Table 2. Summary of typical strengths of the available constraints on the SME at different n orders for rotational invariant, neutrino flavour independent LIV operators. GRB=gamma rays burst, CR=cosmic rays. ^a From neutrino oscillations we have constraints on the difference of LIV coefficients of different flavors up to $O(10^{-28})$ on dim 4, $O(10^{-8})$ and expected up to $O(10^{-14})$ on dim 5 (ICE3), expected up to $O(10^{-4})$ on dim 6 op. * Expected constraint from future experiments.

S. Liberati, *Class. Quant. Grav.* **30**, 133001 (2013)

Lorentz violating solutions of QG are under empirical stress

Is Lorentz invariance compatible with discreteness ?

Yes!

Classical discreteness breaks Lorentz invariance.

Quantum discreteness does not !

Cfr rotational invariance:

If a *classical* vector component can take only discrete values only,
then $SO(3)$ is broken.

But if *quantum* vector can have discrete eigenvalues in a $SO(3)$ invariant theory

$$L_z|m\rangle = \hbar m|m\rangle$$

$$L_z(\theta)|m\rangle_\theta = R(\theta)L_zR(\theta)^{-1}|m\rangle_\theta = \hbar m|m\rangle_\theta$$

$$[L_z, L_z(\theta)] \neq 0$$

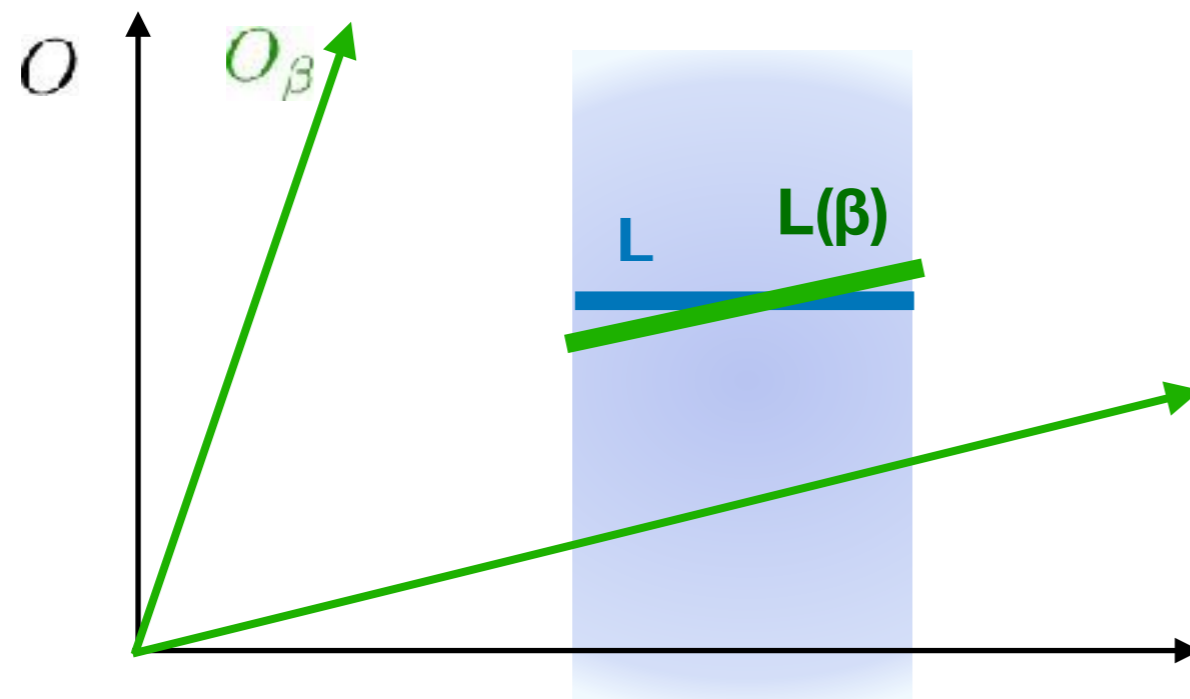
$$|m\rangle_\theta = R(\theta)|m\rangle = \sum_n R_{mn}(\theta)|n\rangle$$

$$L_z|m\rangle = \hbar m|m\rangle$$

$$L_z(\theta)|m\rangle_\theta = R(\theta)L_zR(\theta)^{-1}|m\rangle_\theta = \hbar m|m\rangle_\theta$$

$$[L_z, L_z(\theta)] \neq 0$$

$$|m\rangle_\theta = R(\theta)|m\rangle = \sum_n R_{mn}(\theta)|n\rangle$$



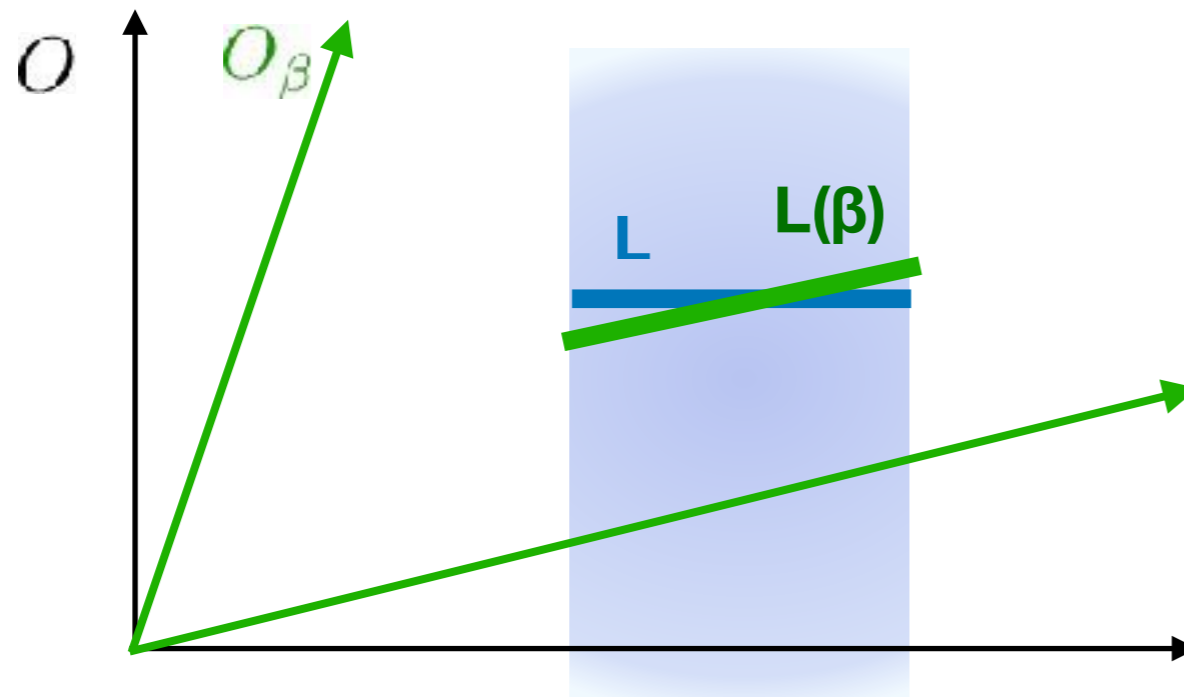
$$[L, L(\beta)] \neq 0$$

$$L_z|m\rangle = \hbar m|m\rangle$$

$$L_z(\theta)|m\rangle_\theta = R(\theta)L_zR(\theta)^{-1}|m\rangle_\theta = \hbar m|m\rangle_\theta$$

$$[L_z, L_z(\theta)] \neq 0$$

$$|m\rangle_\theta = R(\theta)|m\rangle = \sum_n R_{mn}(\theta)|n\rangle$$



$$[L, L(\beta)] \neq 0$$

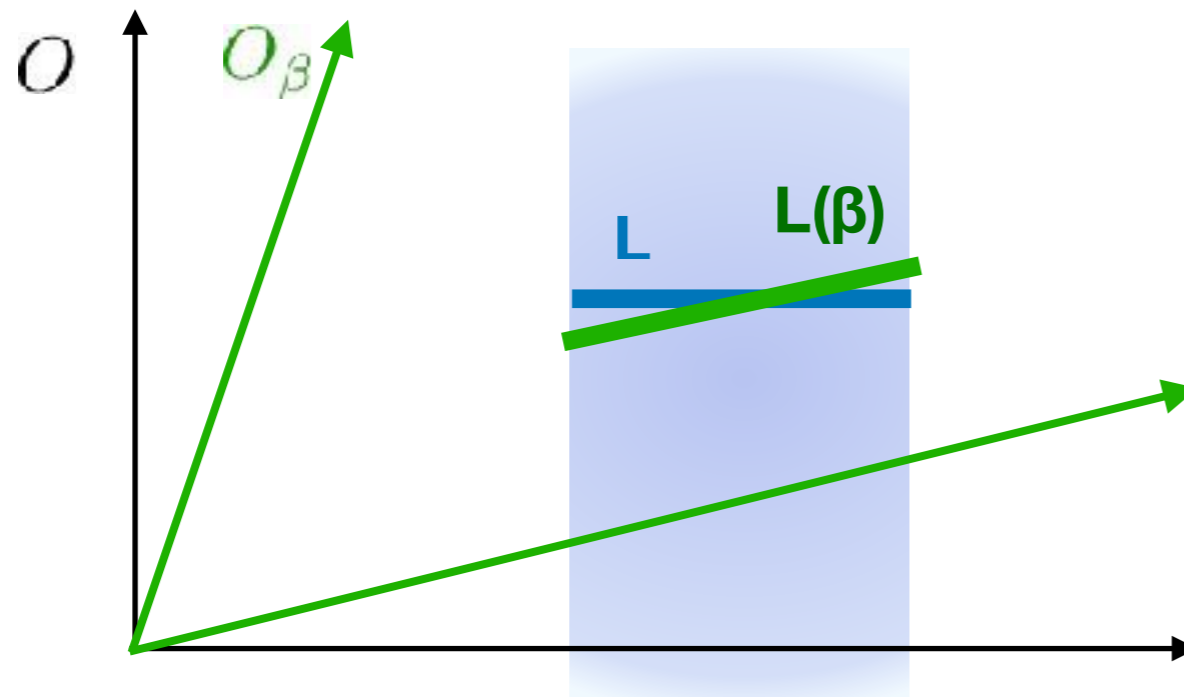
Lorentz invariance and quantum discreteness are compatible

$$L_z|m\rangle = \hbar m|m\rangle$$

$$L_z(\theta)|m\rangle_\theta = R(\theta)L_zR(\theta)^{-1}|m\rangle_\theta = \hbar m|m\rangle_\theta$$

$$[L_z, L_z(\theta)] \neq 0$$

$$|m\rangle_\theta = R(\theta)|m\rangle = \sum_n R_{mn}(\theta)|n\rangle$$



$$[L, L(\beta)] \neq 0$$

Lorentz invariance and quantum discreteness are compatible

=> Geometry is quantum geometry

Empirical evidence: 2: Supersymmetry

Empirical evidence: 2: Supersymmetry

arXiv.org > hep-ex > arXiv:1708.02794

Search or Article II
(Help | Advanced search)

High Energy Physics - Experiment

Search for new phenomena with large jet multiplicities and missing transverse momentum using large-radius jets and flavour-tagging at ATLAS in 13 TeV pp collisions

ATLAS Collaboration

(Submitted on 9 Aug 2017)

A search is presented for particles that decay producing a large jet multiplicity and invisible particles. The event selection applies a veto on the presence of isolated electrons or muons and additional requirements on the number of b-tagged jets and the scalar sum of masses of large-radius jets. Having explored the full ATLAS 2015–2016 dataset of LHC proton–proton collisions at $\sqrt{s} = 13$ TeV, which corresponds to 36.1 fb^{-1} of integrated luminosity, **no evidence is found for physics beyond the Standard Model.** The results are interpreted in the context of simplified models inspired by R-parity-conserving and R-parity-violating supersymmetry, where gluinos are pair-produced. More generic models within the **phenomenological minimal supersymmetric Standard Model** are also considered.

Comments: Comments: 53 pages in total, author list starting page 37, 7 figures, 5 tables, submitted to JHEP, All figures including auxiliary figures are available at [this http URL](#).

Subjects: High Energy Physics - Experiment (hep-ex)

Report number: CERN-EP-2017-138

Cite as: arXiv:1708.02794 [hep-ex]

Empirical evidence: 2: Supersymmetry

arXiv.org > hep-ex > arXiv:1708.02794 Search or Article II
(Help | Advanced search)

High Energy Physics - Experiment

Search for new phenomena with large jet multiplicities and missing transverse momentum using large-radius jets and flavour-tagging at ATLAS in 13 TeV pp collisions

ATLAS Collaboration
(Submitted on 9 Aug 2017)

A search is presented for particles that decay producing a large jet multiplicity and invisible particles. The event selection applies a veto on the presence of isolated electrons or muons and additional requirements on the number of b-tagged jets and the scalar sum of masses of large-radius jets. Having explored the full ATLAS 2015–2016 dataset of LHC proton–proton collisions at $\sqrt{s} = 13$ TeV, which corresponds to 36.1 fb^{-1} of integrated luminosity, **no evidence is found for physics beyond the Standard Model.** The results are interpreted in the context of simplified models inspired by R-parity-conserving and R-parity-violating supersymmetry, where gluinos are pair-produced. More generic models within the **phenomenological minimal supersymmetric Standard Model** are also considered.

Comments: Comments: 53 pages in total, author list starting page 37, 7 figures, 5 tables, submitted to JHEP, All figures including auxiliary figures are available at [this http URL](#).

Subjects: High Energy Physics - Experiment (hep-ex)

Report number: CERN-EP-2017-138

Cite as: arXiv:1708.02794 [hep-ex]

Once again, no sign of supersymmetry

Empirical evidence: 2: Supersymmetry

arXiv.org > hep-ex > arXiv:1708.02794 Search or Article II
(Help | Advanced search)

High Energy Physics - Experiment

Search for new phenomena with large jet multiplicities and missing transverse momentum using large-radius jets and flavour-tagging at ATLAS in 13 TeV pp collisions

ATLAS Collaboration
(Submitted on 9 Aug 2017)

A search is presented for particles that decay producing a large jet multiplicity and invisible particles. The event selection applies a veto on the presence of isolated electrons or muons and additional requirements on the number of b-tagged jets and the scalar sum of masses of large-radius jets. Having explored the full ATLAS 2015–2016 dataset of LHC proton–proton collisions at $\sqrt{s} = 13$ TeV, which corresponds to 36.1 fb^{-1} of integrated luminosity, **no evidence is found for physics beyond the Standard Model.** The results are interpreted in the context of simplified models inspired by R-parity-conserving and R-parity-violating supersymmetry, where gluinos are pair-produced. More generic models within the **phenomenological minimal supersymmetric Standard Model** are also considered.

Comments: Comments: 53 pages in total, author list starting page 37, 7 figures, 5 tables, submitted to JHEP, All figures including auxiliary figures are available at [this http URL](#).

Subjects: High Energy Physics - Experiment (hep-ex)

Report number: CERN-EP-2017-138

Cite as: arXiv:1708.02794 [hep-ex]

Once again, no sign of supersymmetry

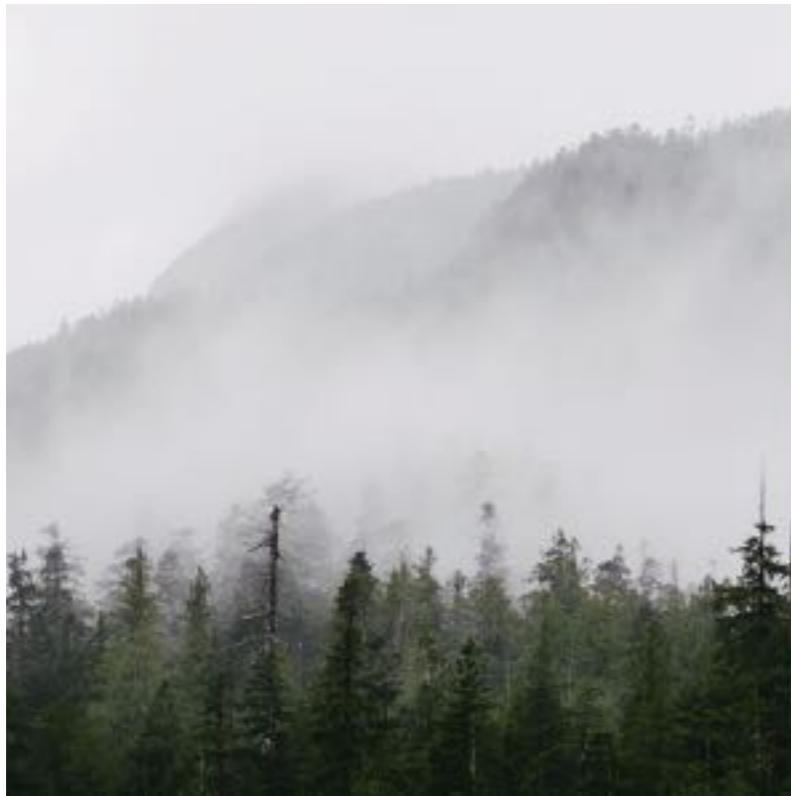
Solution of QG using supersymmetry are under empirical stress

A point about philosophy of science:

- Popper's *falsification*: theories are either "OK" or "proved wrong".



Karl Popper



A point about philosophy of science:

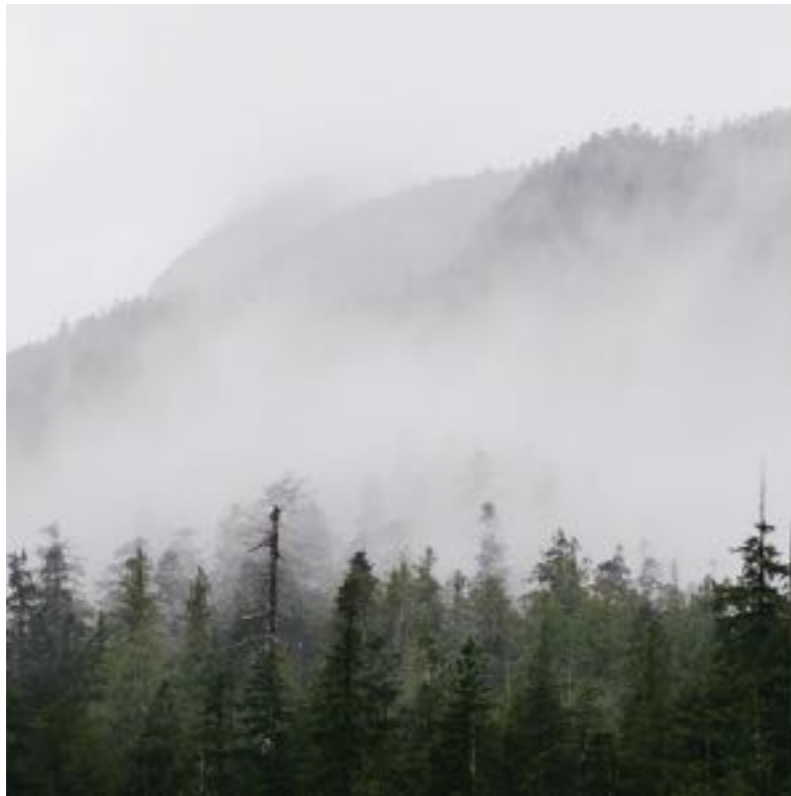
- Popper's *falsification*: theories are either "OK" or "proved wrong".
- Bayesian "*confirmation*": we have "degrees of confidence" in theories; these are lowered, or enhanced, by empirical (dis-)confirmation.



Karl Popper



Bruno De Finetti



A point about philosophy of science:

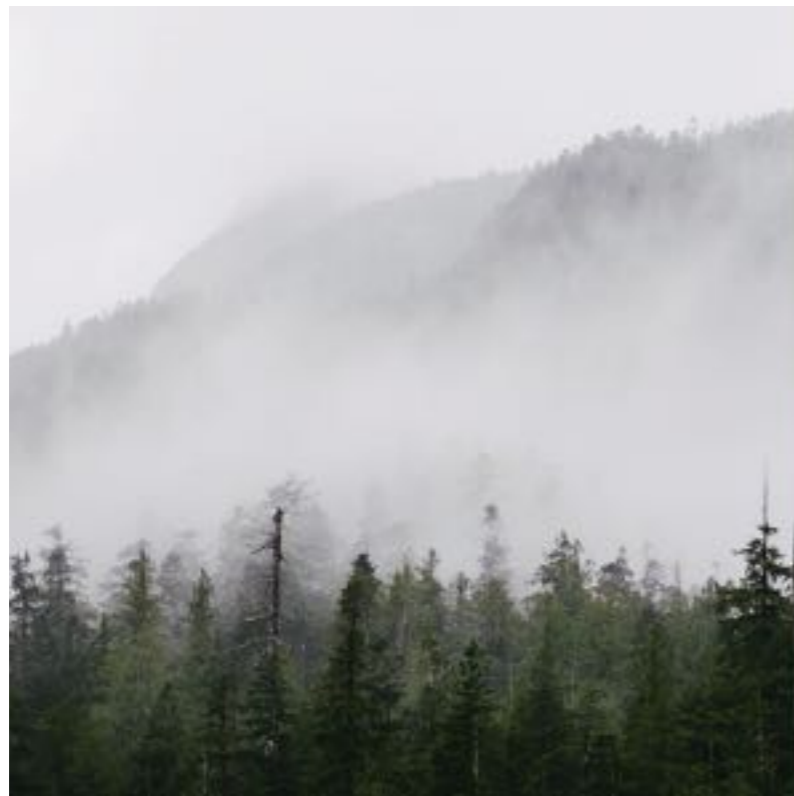
- Popper's *falsification*: theories are either "OK" or "proved wrong".
- Bayesian "*confirmation*": we have "degrees of confidence" in theories; these are lowered, or enhanced, by empirical (dis-)confirmation.



Karl Popper



Bruno De Finetti



Not seeing a giraffe in the French Alps during a hike, does **not** prove that there are no giraffes in French Alps

A point about philosophy of science:

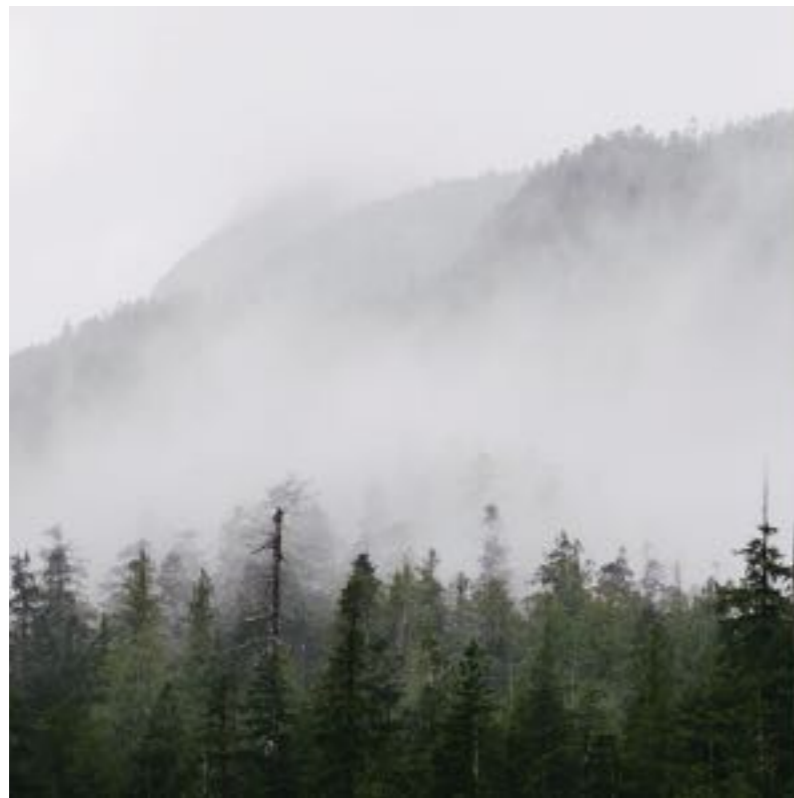
- Popper's *falsification*: theories are either "OK" or "proved wrong".
- Bayesian "*confirmation*": we have "degrees of confidence" in theories; these are lowered, or enhanced, by empirical (dis-)confirmation.



Karl Popper



Bruno De Finetti



Not seeing a giraffe in the French Alps during a hike, does **not** prove that there are no giraffes in French Alps

A point about philosophy of science:

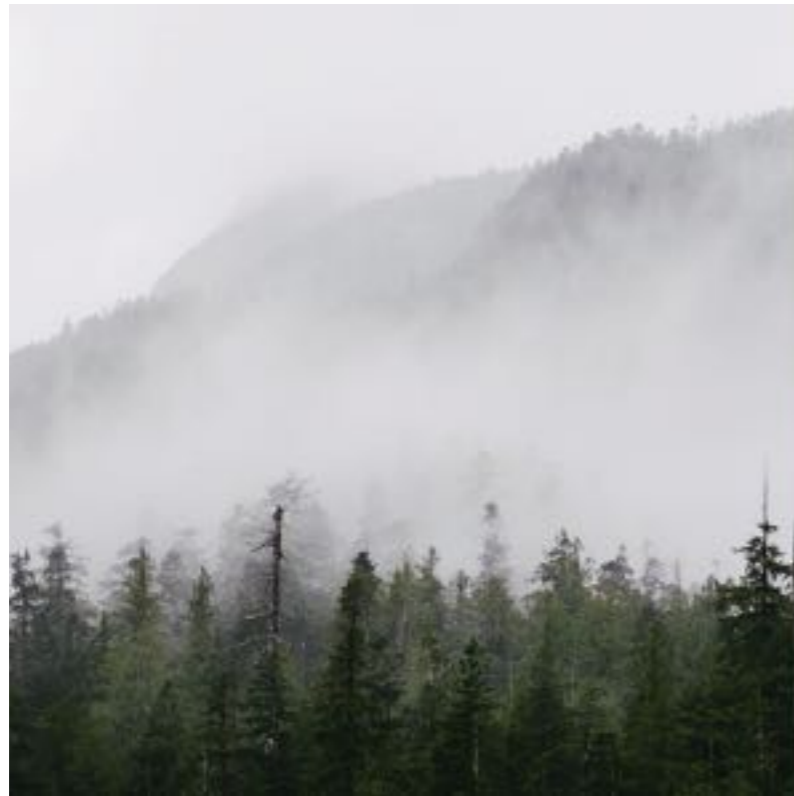
- Popper's *falsification*: theories are either "OK" or "proved wrong".
- Bayesian "*confirmation*": we have "degrees of confidence" in theories; these are lowered, or enhanced, by empirical (dis-)confirmation.



Karl Popper



Bruno De Finetti



Not seeing a giraffe in the French Alps during a hike, does **not** prove that there are no giraffes in French Alps

But if for thirty years nobody sees a giraffe...

A point about philosophy of science:

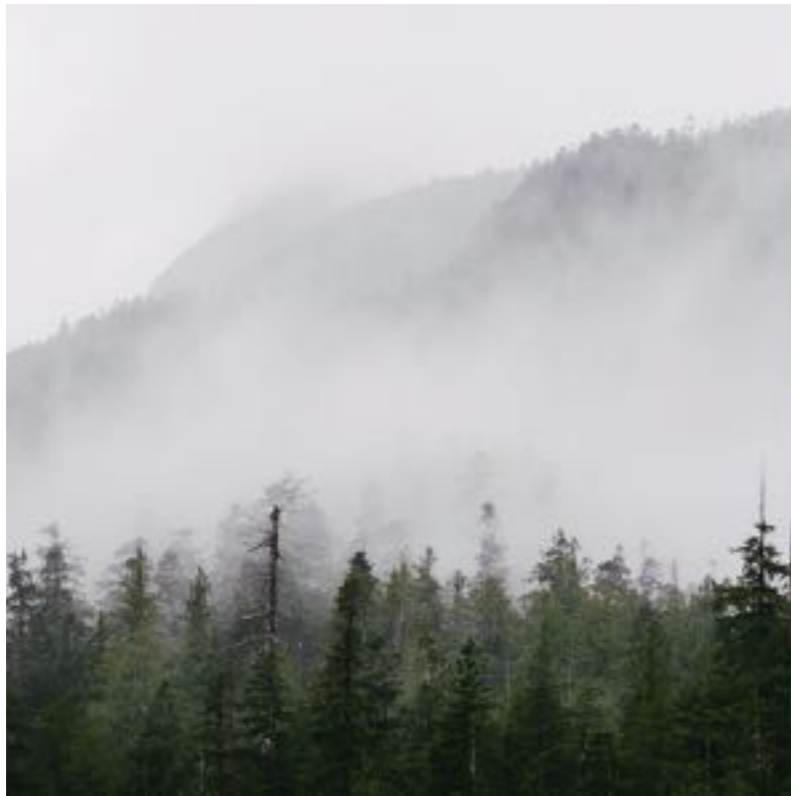
- Popper's *falsification*: theories are either "OK" or "proved wrong".
- Bayesian "*confirmation*": we have "degrees of confidence" in theories; these are lowered, or enhanced, by empirical (dis-)confirmation.



Karl Popper



Bruno De Finetti



Not seeing a giraffe in the French Alps during a hike, does **not** prove that there are no giraffes in French Alps

But if for thirty years nobody sees a giraffe...

And we have now heard that supersymmetry is "going to be seen soon" for more than thirty years....

A point about philosophy of science:

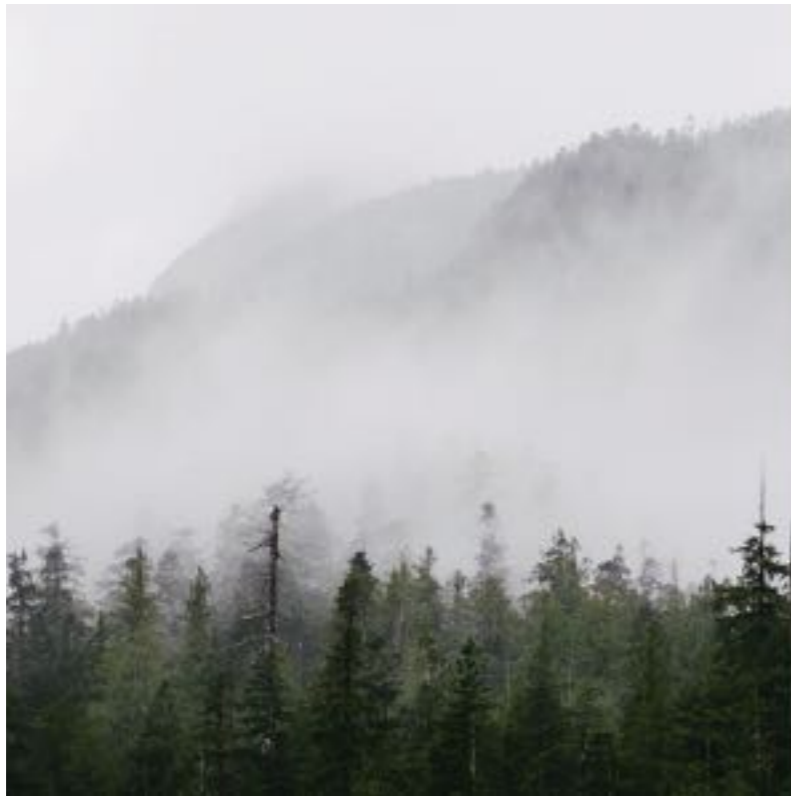
- Popper's *falsification*: theories are either "OK" or "proved wrong".
- Bayesian "*confirmation*": we have "degrees of confidence" in theories; these are lowered, or enhanced, by empirical (dis-)confirmation.



Karl Popper



Bruno De Finetti



Not seeing a giraffe in the French Alps during a hike, does **not** prove that there are no giraffes in French Alps

But if for thirty years nobody sees a giraffe...

And we have now heard that supersymmetry is "going to be seen soon" for more than thirty years....

Solution of QG using supersymmetry are under empirical stress

Empirical evidence: 3: Lab experiments (future)

Empirical evidence: 3: Lab experiments

Analog systems

Planck scale effects
in the lab

Violations of QM
suggested by QG

Quantum property
of the metric

Empirical evidence: 3: Lab experiments

Analog systems



Planck scale effects
in the lab

Violations of QM
suggested by QG

Quantum property
of the metric

Empirical evidence: 3: Lab experiments

Analog systems



Test the **consequences** of an assumption.
Not the **assumption** themselves.

Planck scale effects
in the lab

Violations of QM
suggested by QG

Quantum property
of the metric

Empirical evidence: 3: Lab experiments

Analog systems



Test the **consequences** of an assumption.
Not the **assumption** themselves.

Planck scale effects
in the lab

NOT predicted by most QG theories

Violations of QM
suggested by QG

Quantum property
of the metric

Empirical evidence: 3: Lab experiments

Analog systems



Test the **consequences** of an assumption.
Not the **assumption** themselves.

Planck scale effects
in the lab

NOT predicted by most QG theories

Violations of QM
suggested by QG



Quantum property
of the metric

Empirical evidence: 3: Lab experiments

Analog systems



Test the **consequences** of an assumption.
Not the **assumption** themselves.

Planck scale effects
in the lab

NOT predicted by most QG theories

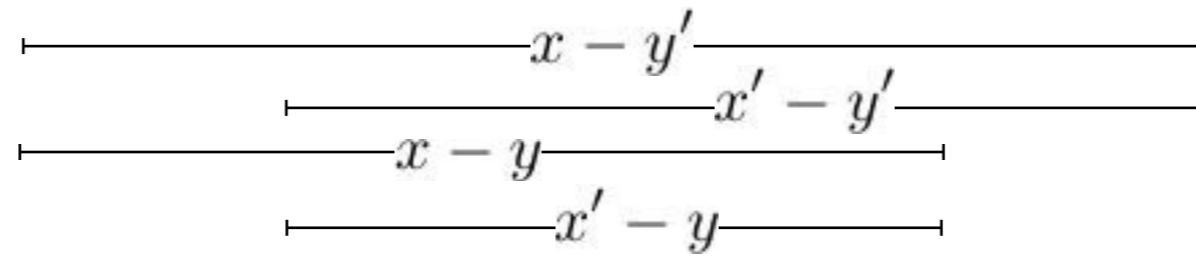
Violations of QM
suggested by QG



Quantum property
of the metric

Can falsify the hypothesis that the
gravitational field is classical.

Is the metric a quantum entity?



$$|\psi\rangle = (|1x\rangle + |1x'\rangle) \otimes (|2y\rangle + |2y'\rangle)$$

$$\begin{aligned}
 |\psi\rangle &= |1x\rangle \otimes |2y\rangle + |1x'\rangle \otimes |2y\rangle + |1x\rangle \otimes |2y'\rangle + |1x'\rangle \otimes |2y'\rangle \\
 |\psi(t)\rangle &= |1x\rangle \otimes |2y\rangle e^{\frac{i}{\hbar} \frac{Gm^2}{x-y} t} + |1x'\rangle \otimes |2y\rangle e^{\frac{i}{\hbar} \frac{Gm^2}{x'-y} t} + |1x\rangle \otimes |2y'\rangle e^{\frac{i}{\hbar} \frac{Gm^2}{x-y'} t} + |1x'\rangle \otimes |2y'\rangle e^{\frac{i}{\hbar} \frac{Gm^2}{x'-y'} t}
 \end{aligned}$$

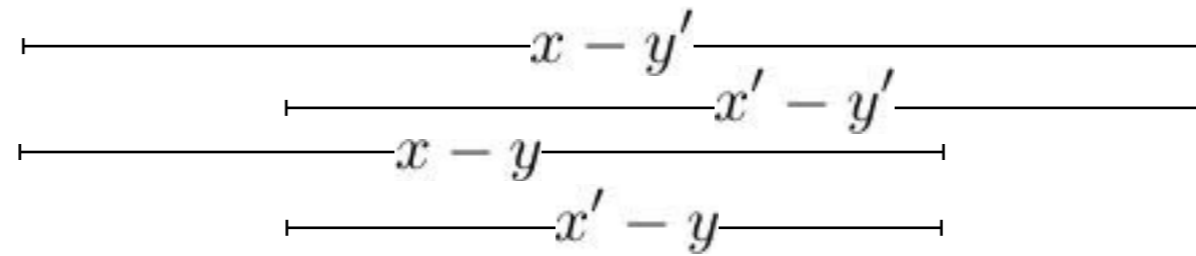
$$\frac{Gm^2}{\hbar \Delta x} t \sim \pi$$

S Bose, A Mazumdar, GW Morley, H Ulbricht, M Toroš, M Paternostro, A Geraci, P Barker, MS Kim, G Milburn: *A Spin Entanglement Witness for Quantum Gravity*, 2017.

C Marletto, V Vedral: *An entanglement-based test of quantum gravity using two massive particles*, 2017.

Is the metric a quantum entity?

Can falsify the hypothesis that the gravitational field is classical.



$$|\psi\rangle = (|1x\rangle + |1x'\rangle) \otimes (|2y\rangle + |2y'\rangle)$$

$$\begin{aligned}
 |\psi\rangle &= |1x\rangle \otimes |2y\rangle + |1x'\rangle \otimes |2y\rangle + |1x\rangle \otimes |2y'\rangle + |1x'\rangle \otimes |2y'\rangle \\
 |\psi(t)\rangle &= |1x\rangle \otimes |2y\rangle e^{\frac{i}{\hbar} \frac{Gm^2}{x-y} t} + |1x'\rangle \otimes |2y\rangle e^{\frac{i}{\hbar} \frac{Gm^2}{x'-y} t} + |1x\rangle \otimes |2y'\rangle e^{\frac{i}{\hbar} \frac{Gm^2}{x-y'} t} + |1x'\rangle \otimes |2y'\rangle e^{\frac{i}{\hbar} \frac{Gm^2}{x'-y'} t}
 \end{aligned}$$

$$\frac{Gm^2}{\hbar \Delta x} t \sim \pi$$

S Bose, A Mazumdar, GW Morley, H Ulbricht, M Toroš, M Paternostro, A Geraci, P Barker, MS Kim, G Milburn: *A Spin Entanglement Witness for Quantum Gravity*, 2017.

C Marletto, V Vedral: *An entanglement-based test of quantum gravity using two massive particles*, 2017.

Empirical evidence: 4: The Sky (future)

Empirical evidence: 4: The Sky

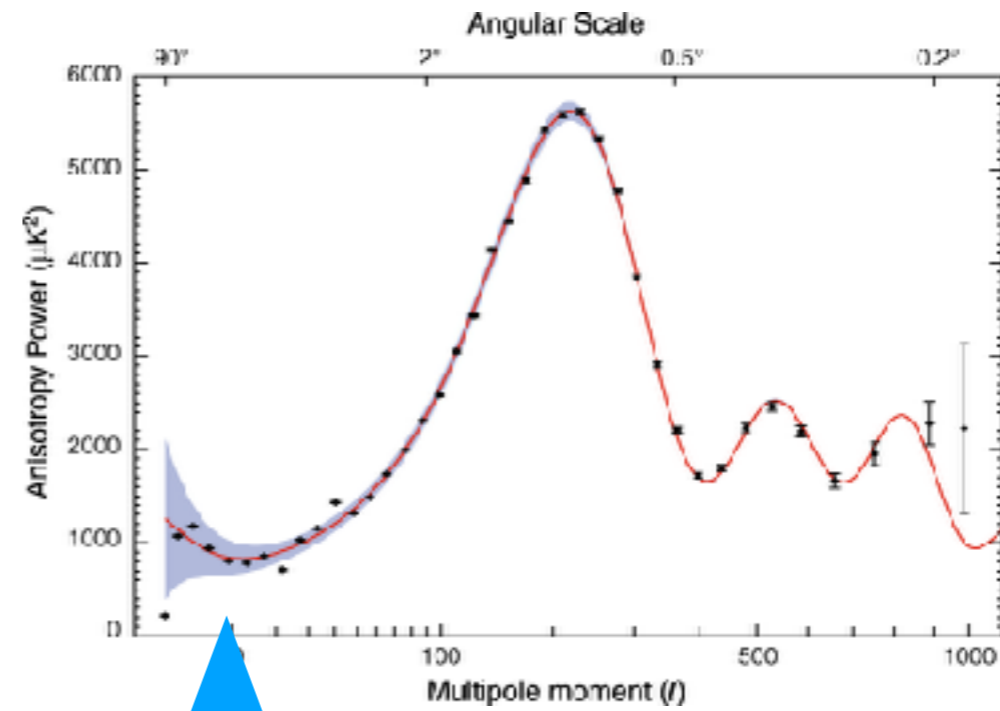
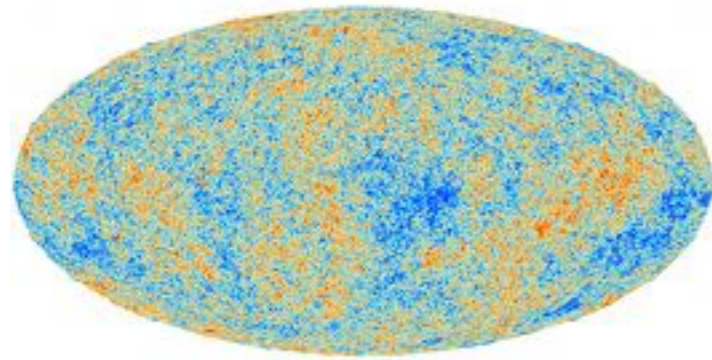
a) Early Universe:

“Quantum cosmology”

b) Black holes:

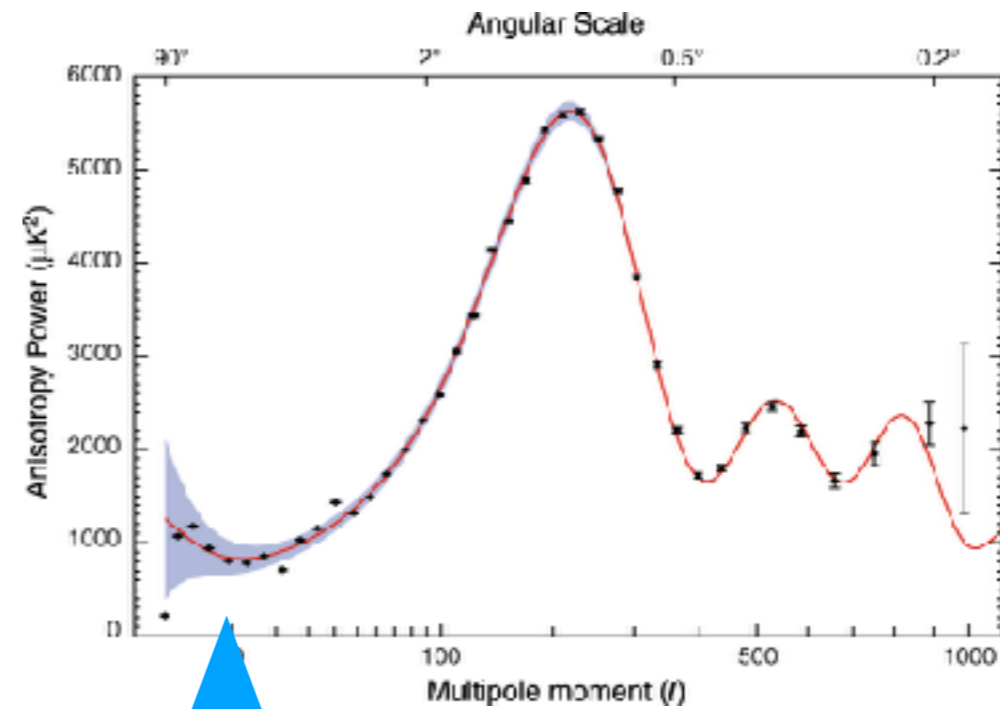
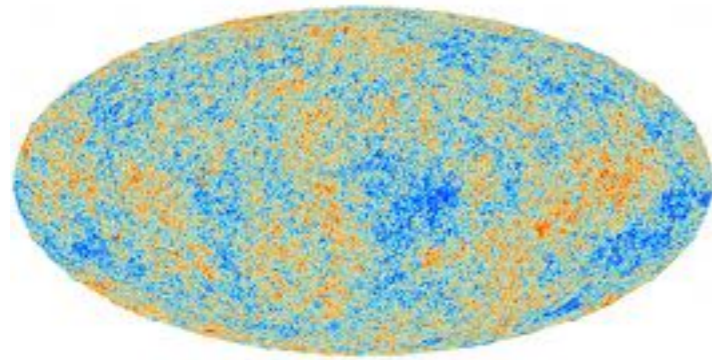
Disruption of the photon ring

Planck Stars



Large activity to describe the physics of the very early universe, and find traces in the CMB

Notice: this is all physics of **few** degrees of freedom!

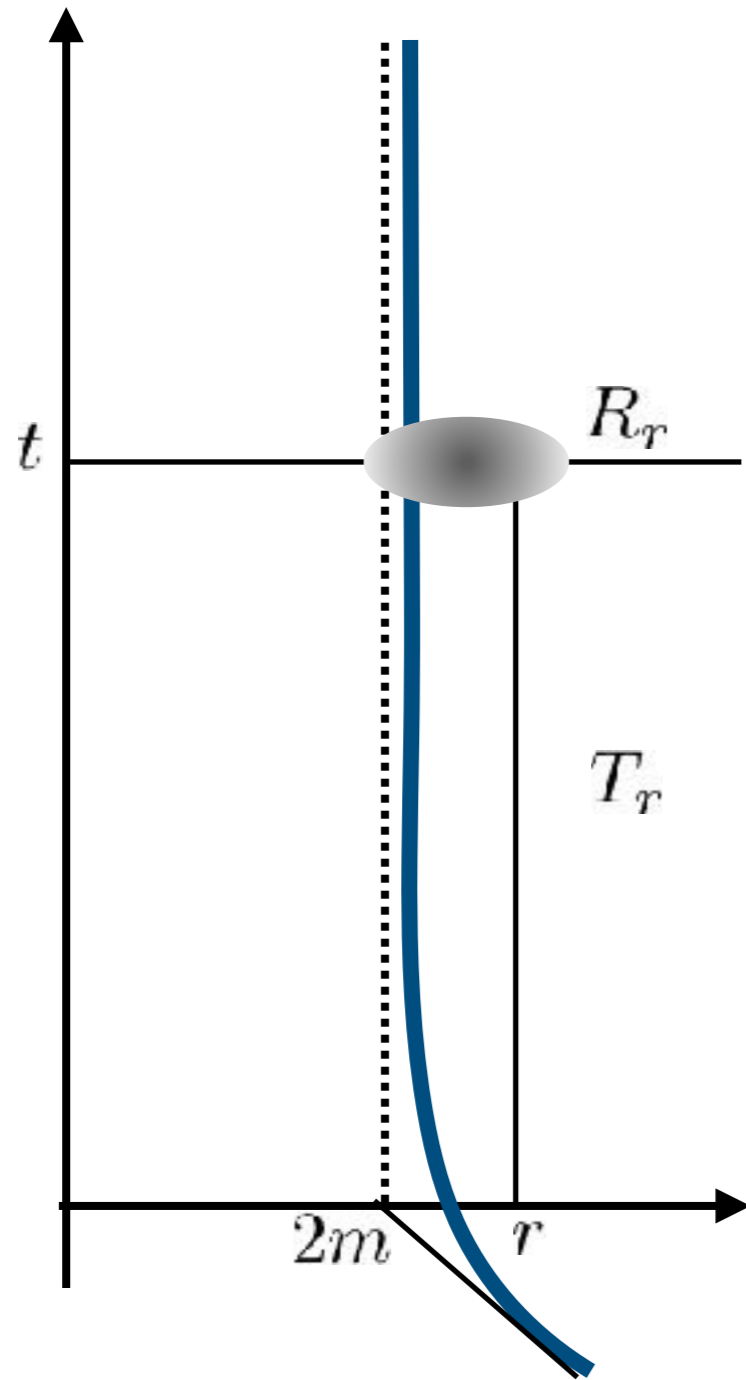


Large activity to describe the physics of the very early universe, and find traces in the CMB

Notice: this is all physics of **few** degrees of freedom!

Great effort to find testable consequences of the theories in course

Small effects pile up over time



$$R_r T_r \sim \frac{m}{r^3} \left(1 - \frac{2m}{r}\right) t$$

$$r_{(\max R_r T_r)} = \frac{7}{3} m$$

$$R_r T_r \sim \frac{l}{l_{Pl}} \rightarrow T \sim m^2$$

- Wide quantum fluctuations of the metric

Giddings

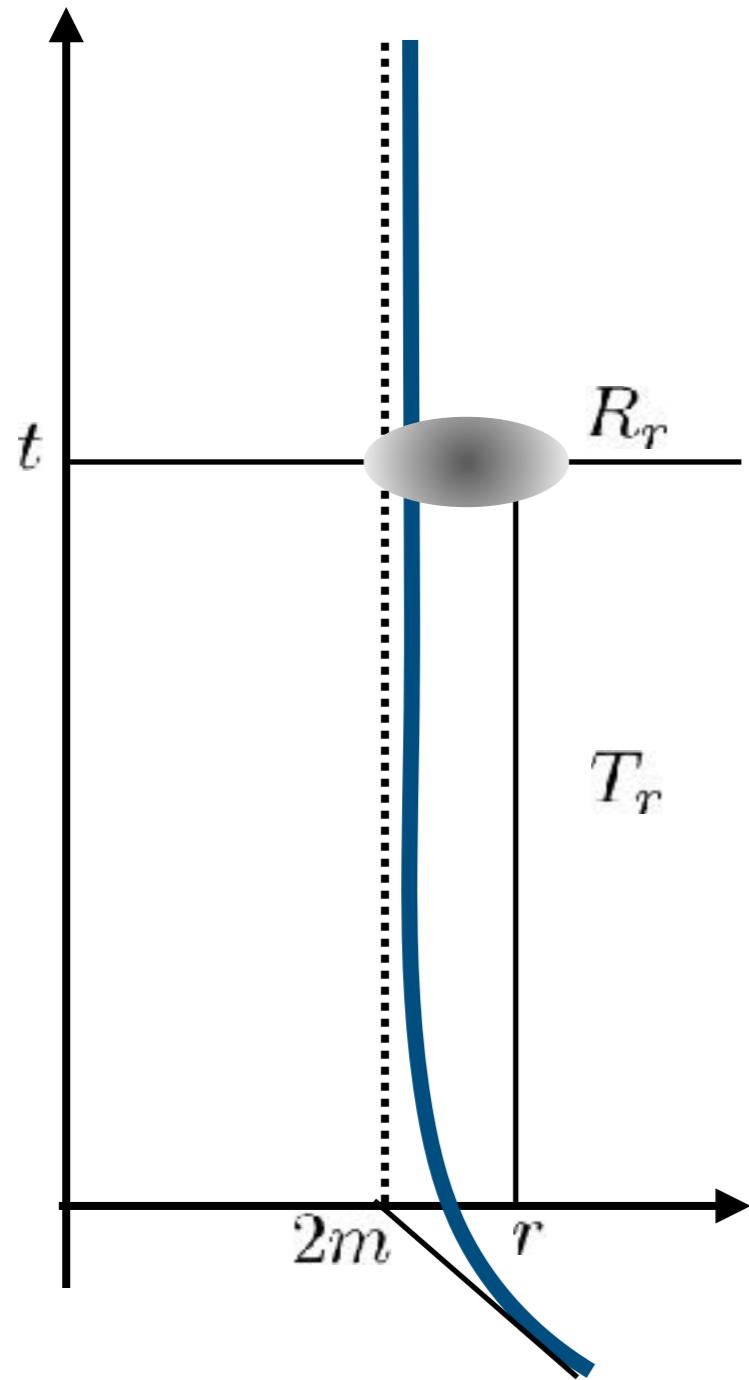
- Boson condensate of low energy gravitons

Dvali

- Fluctuations of the causal structure allowing black hole to decay

Haggard, Barrau, Vidotto, CR

Small effects pile up over time



$$R_r T_r \sim \frac{m}{r^3} \left(1 - \frac{2m}{r}\right) t$$

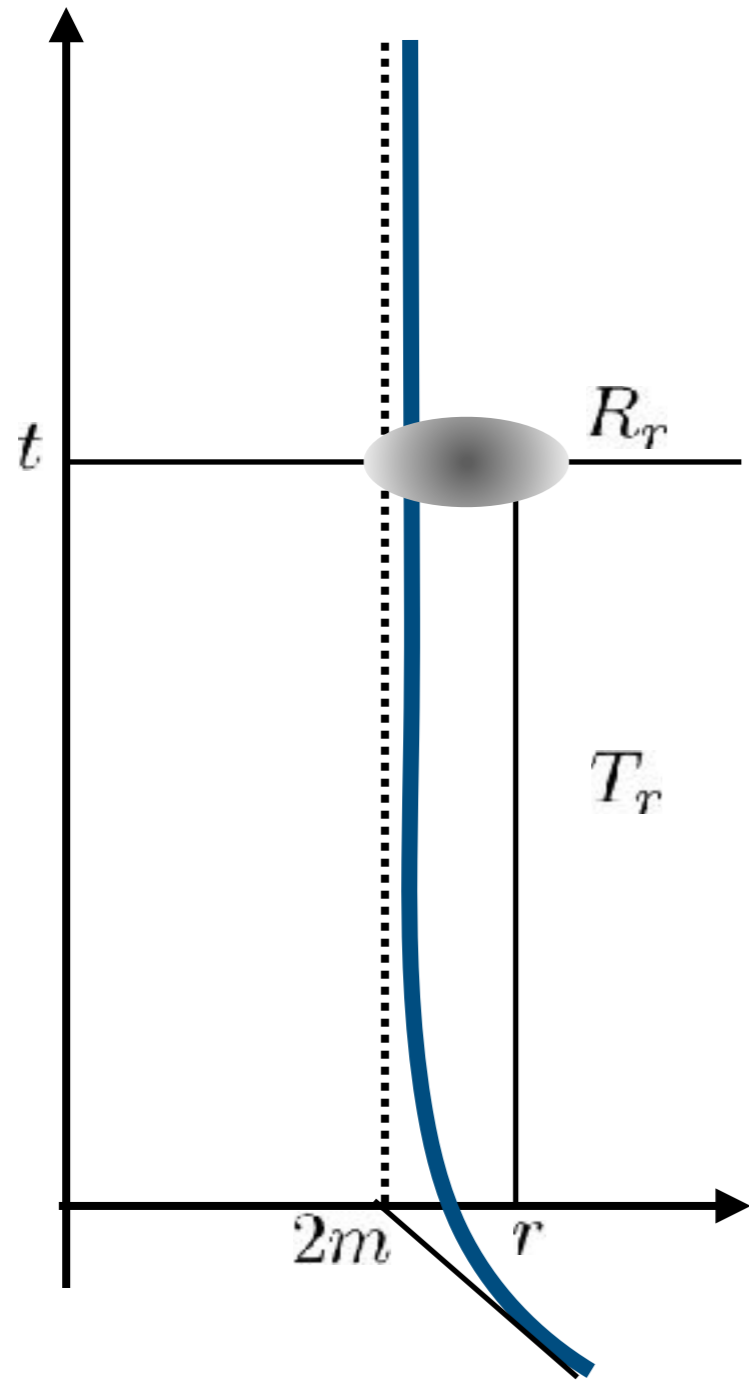
$$r_{(\max R_r T_r)} = \frac{7}{3} m$$

$$R_r T_r \sim \frac{l}{l_{Pl}} \rightarrow T \sim m^2$$

- Wide quantum fluctuations of the metric
Giddings
- Boson condensate of low energy gravitons
Dvali
- Fluctuations of the causal structure allowing black hole to decay

Haggard, Barrau, Vidotto, CR

Small effects pile up over time



$$R_r T_r \sim \frac{m}{r^3} \left(1 - \frac{2m}{r}\right) t$$

$$r_{(\max R_r T_r)} = \frac{7}{3} m$$

$$R_r T_r \sim \frac{l}{l_{Pl}} \rightarrow T \sim m^2$$

- Wide quantum fluctuations of the metric

Giddings

- Boson condensate of low energy gravitons

Dvali

- Fluctuations of the causal structure allowing black hole to decay

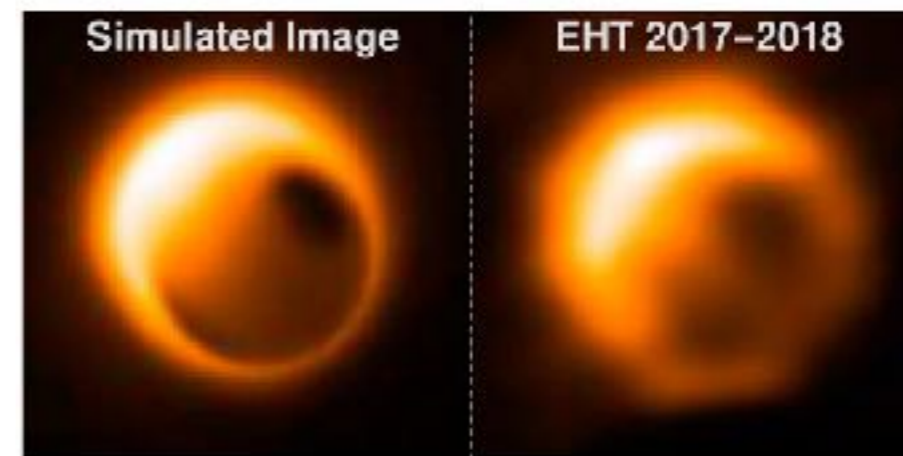
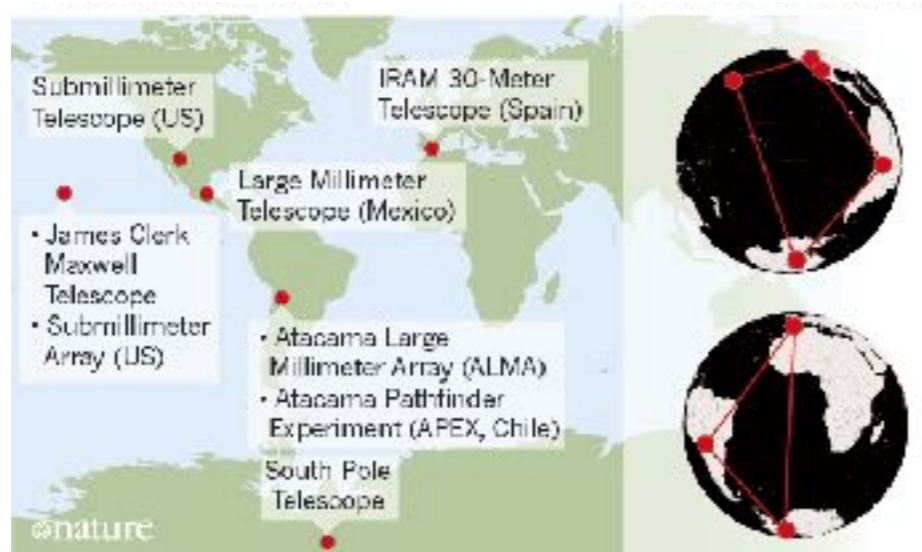
Haggard, Barrau, Vidotto, CR

- Wide quantum fluctuations of the metric

Theoretical reason: to bring information out of the hole

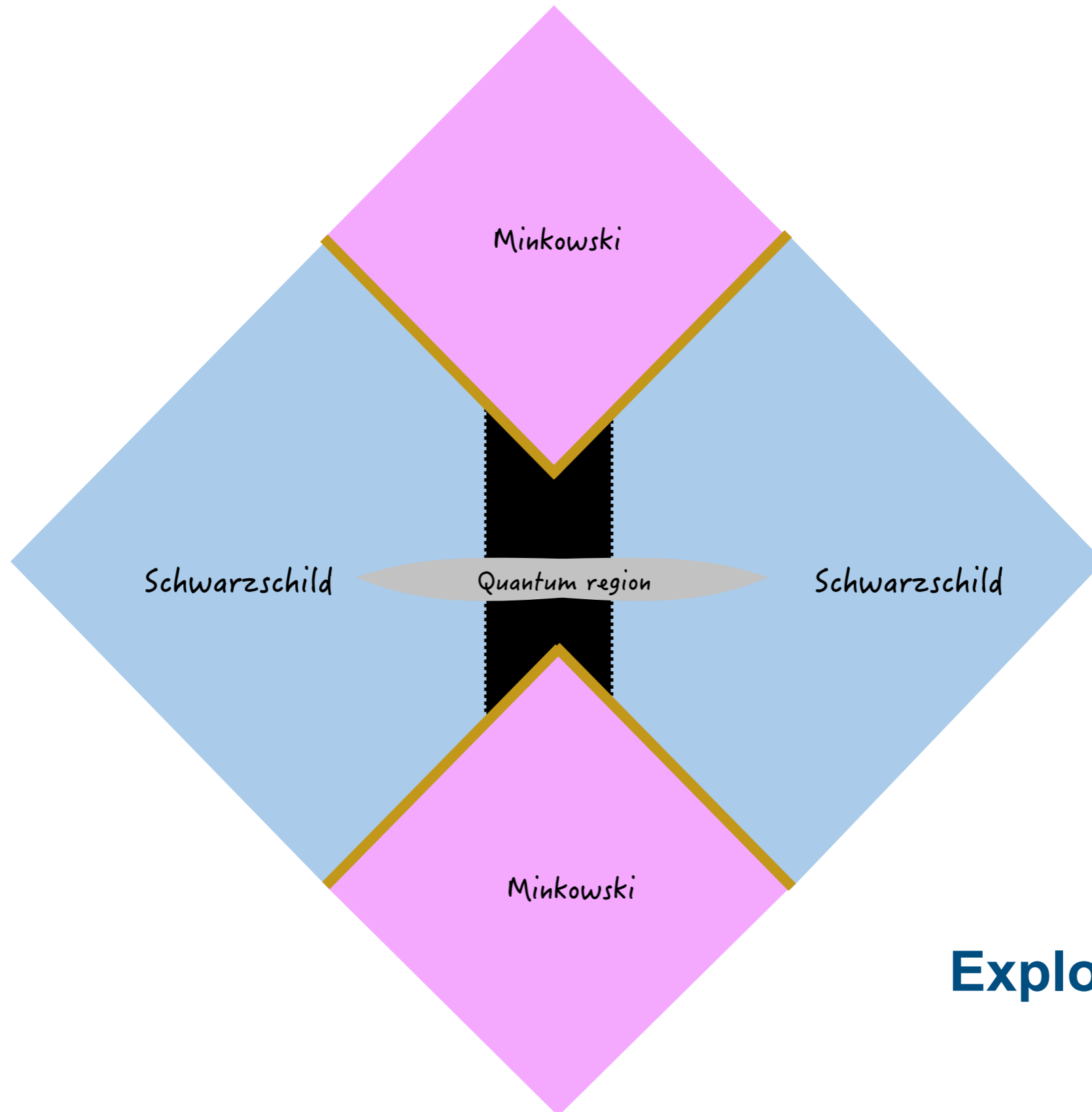
Observable consequence: Event Horizon Telescope

Possibly visible distortion of the photon ring

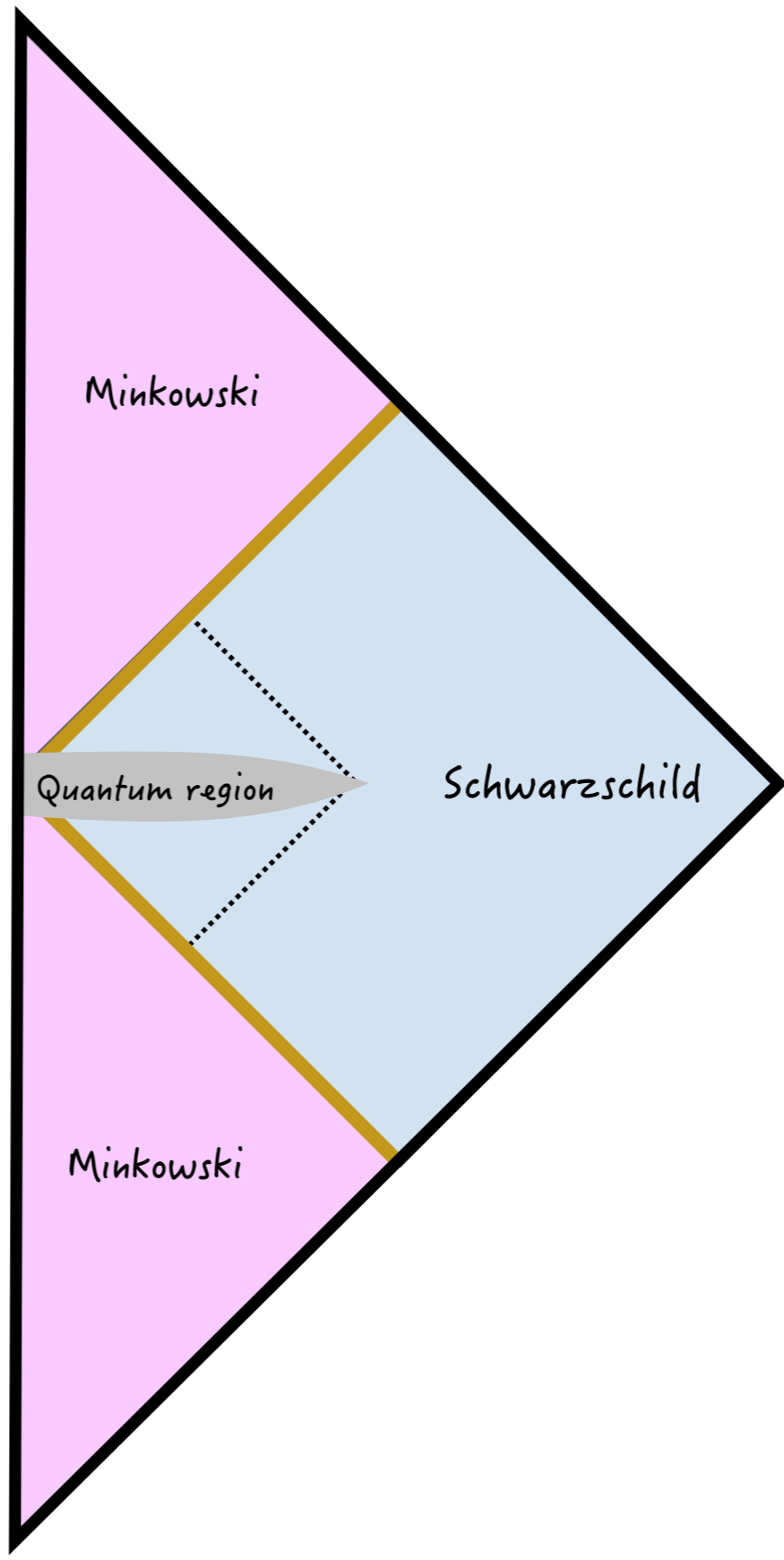


Imaging an Event Horizon: Mitigation of Scattering toward Sagittarius A* [Fish et al 2014](#)

- Fluctuations of the causal structure allowing black hole to decay



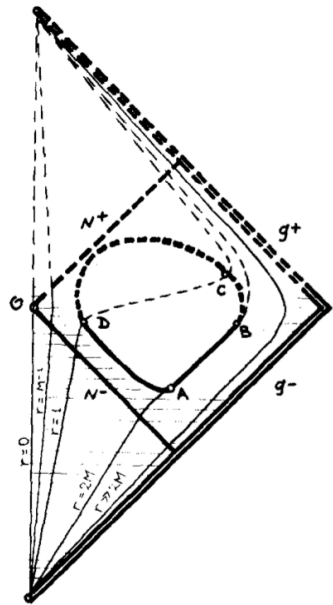
Exploding holes



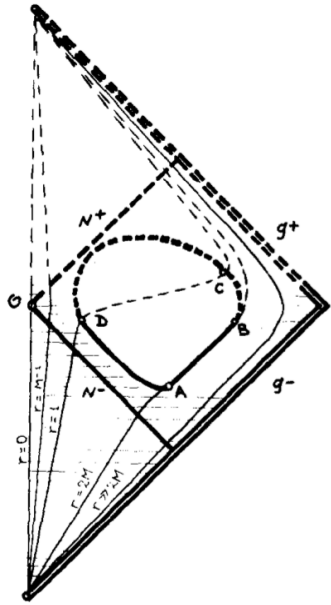
Exploding holes

Exploding holes

■ Frolov, Vilkovinski '79



Exploding holes



- Frolov, Vilkovinski '79

- Stephen, t'Hooft, Whithing '93

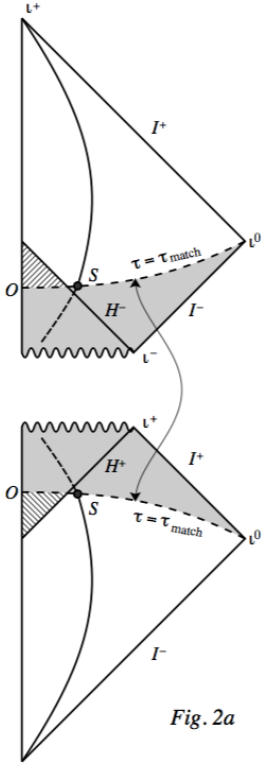


Fig. 2a

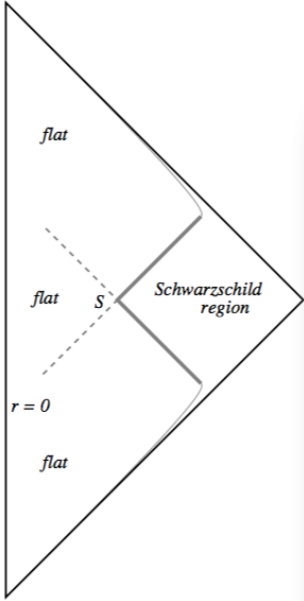


Fig. 2b

Exploding holes

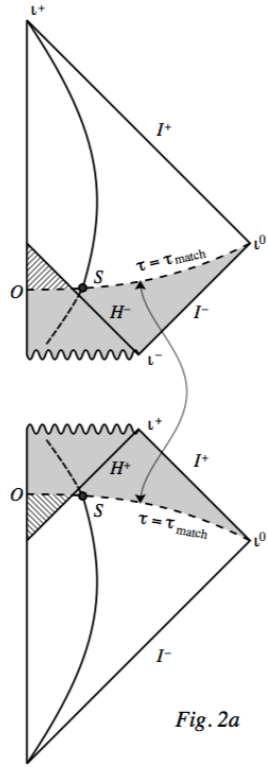
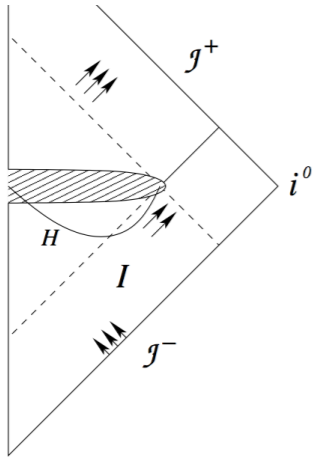
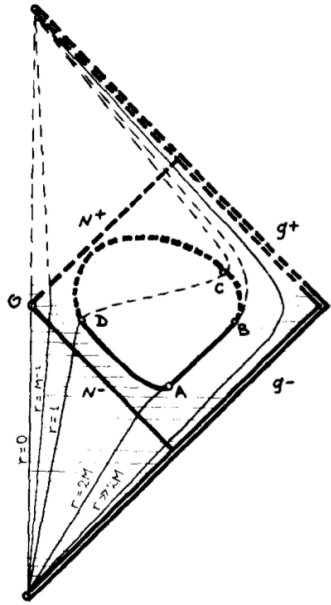


Fig. 2a

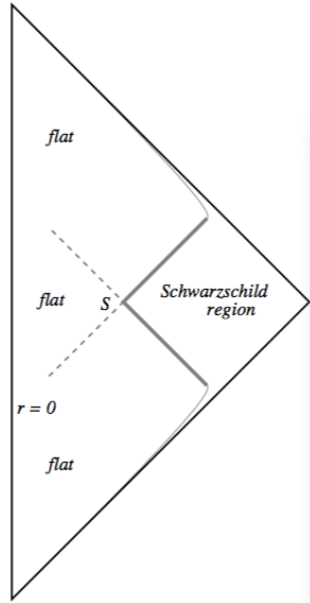
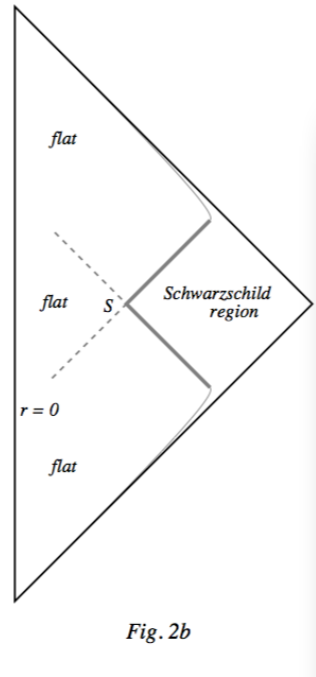
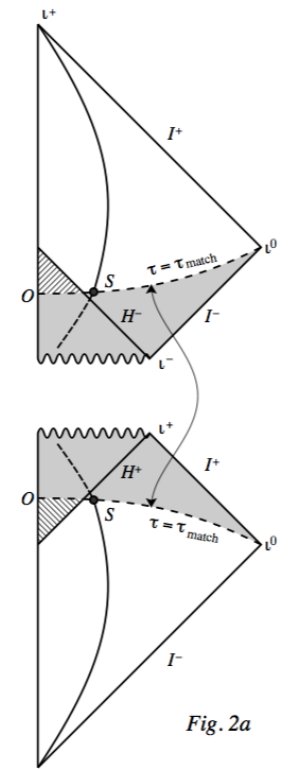
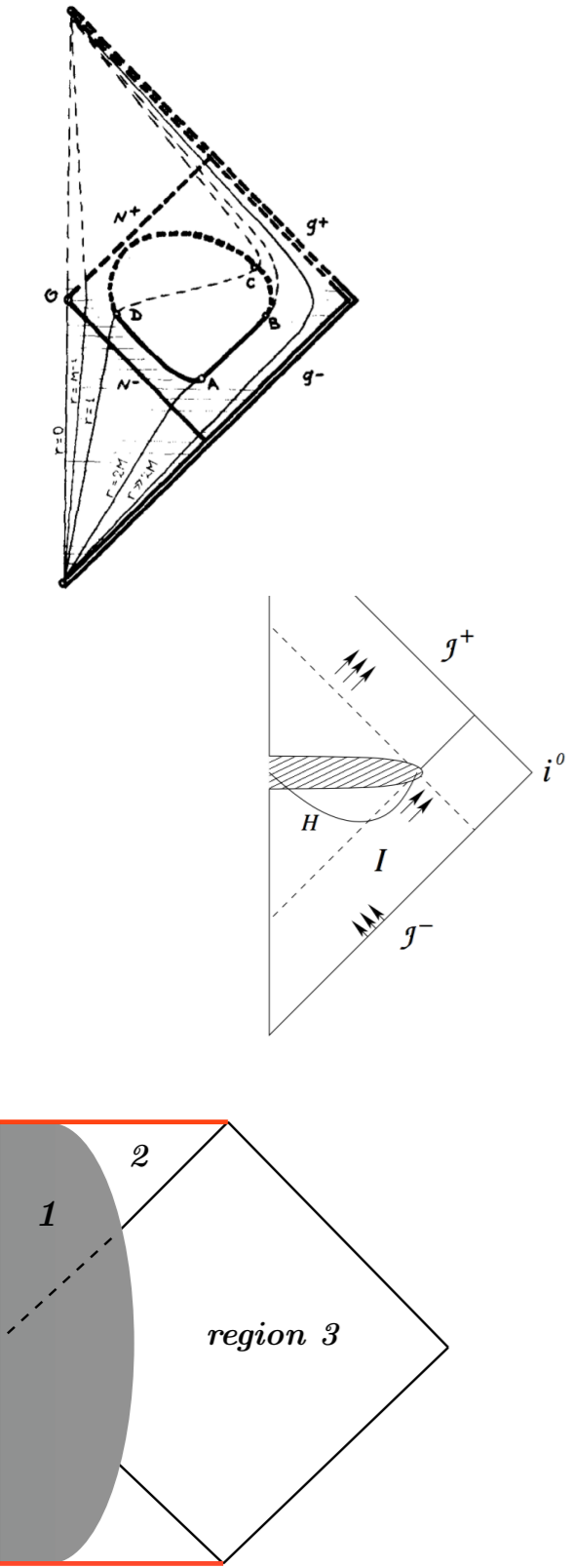


Fig. 2b

- Frolov, Vilkovinski '79
- Stephen, t'Hooft, Whithing '93
- Ashtekar, Bojowald '05

Exploding holes

- Frolov, Vilkovinski '79
- Stephen, t'Hooft, Whithing '93
- Ashtekar, Bojowald '05
- Modesto '06



Exploding holes

■ Frolov, Vilkovinski '79

■ Stephen, t'Hooft, Whiting '93

■ Ashtekar, Bojowald '05

■ Modesto '06

■ Hayward '06

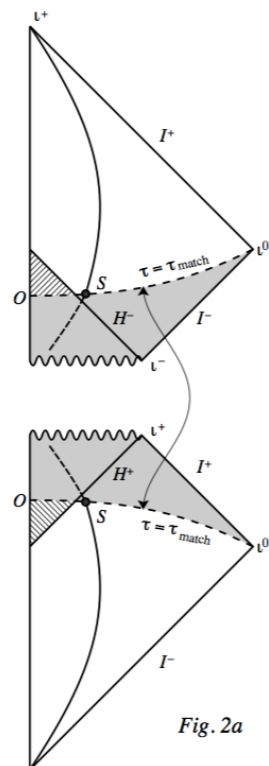
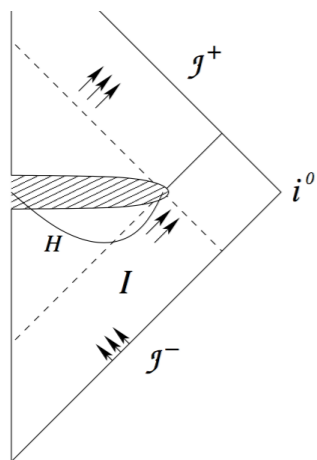
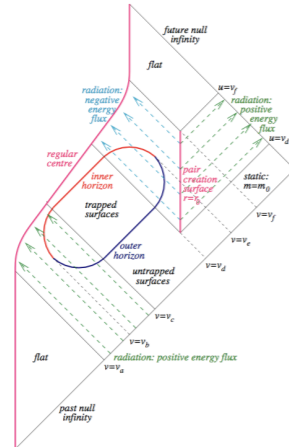
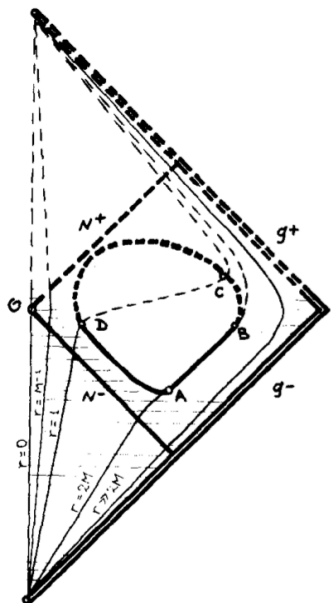


Fig. 2a

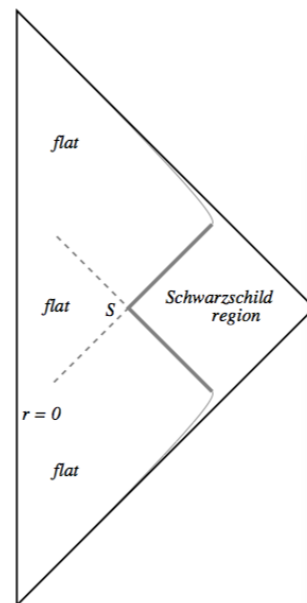
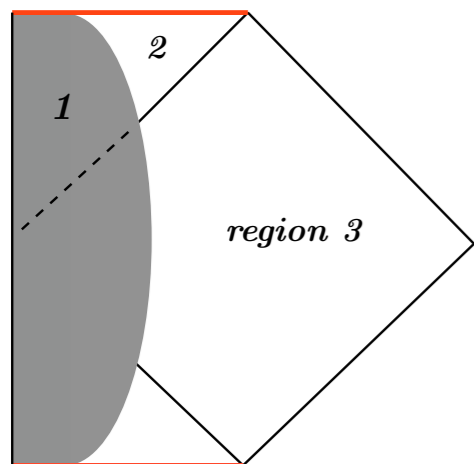


Fig. 2b



Exploding holes

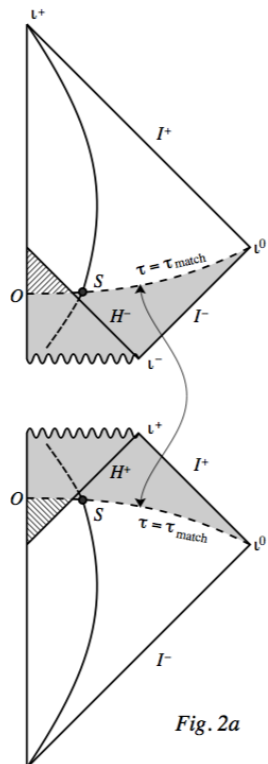
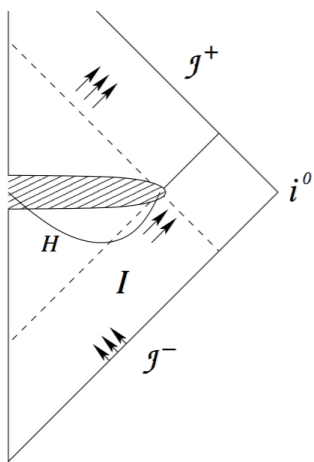
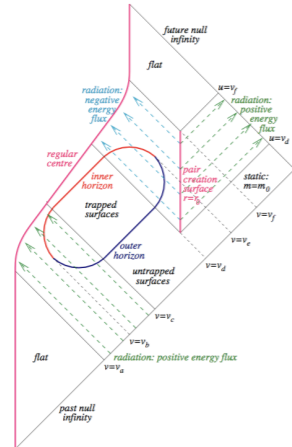
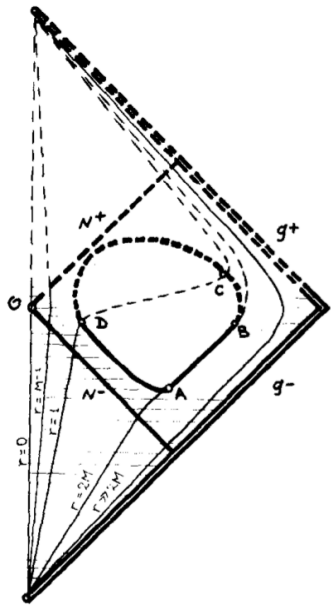


Fig. 2a

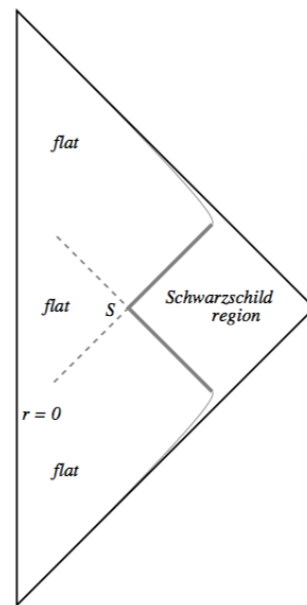
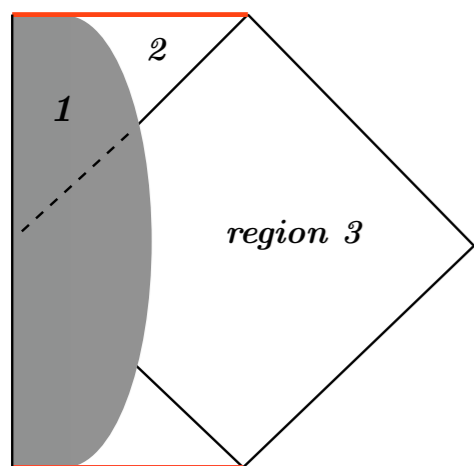


Fig. 2b



- Frolov, Vilkovinski '79
- Stephen, t'Hooft, Whiting '93
- Ashtekar, Bojowald '05
- Modesto '06
- Hayward '06
- Hajicek Kieffer '01

Exploding holes

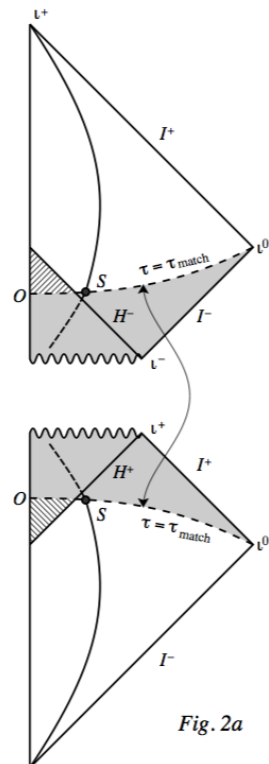
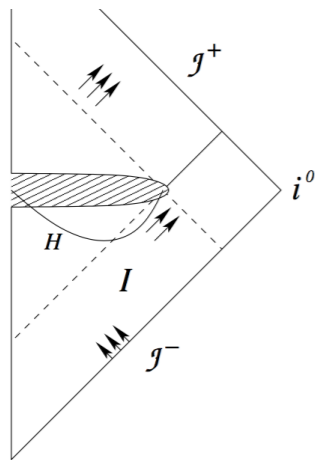
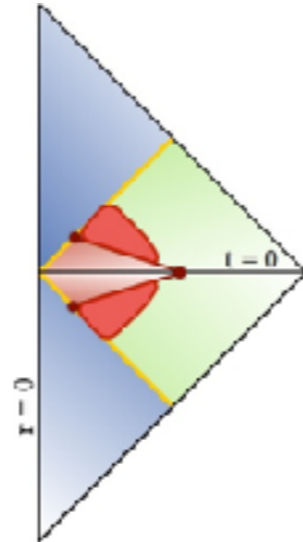
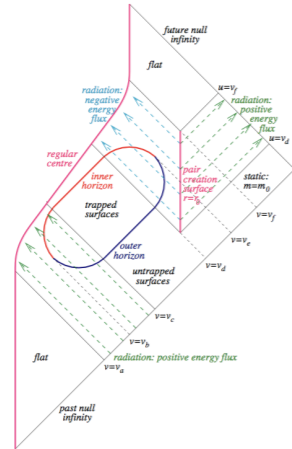
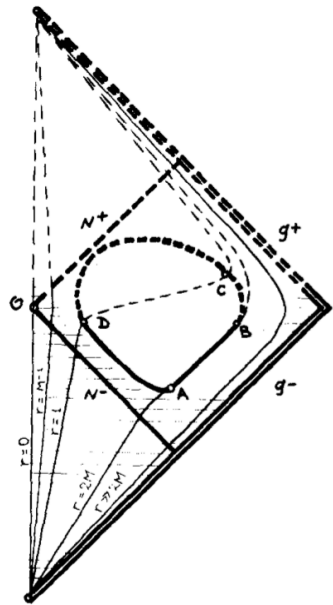


Fig. 2a

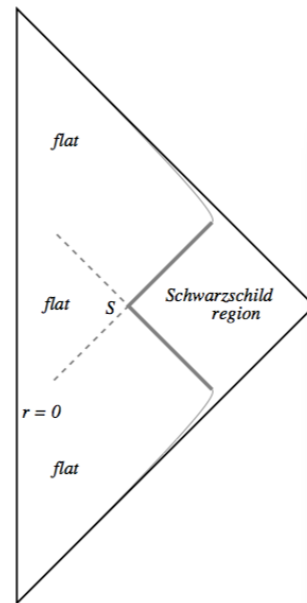
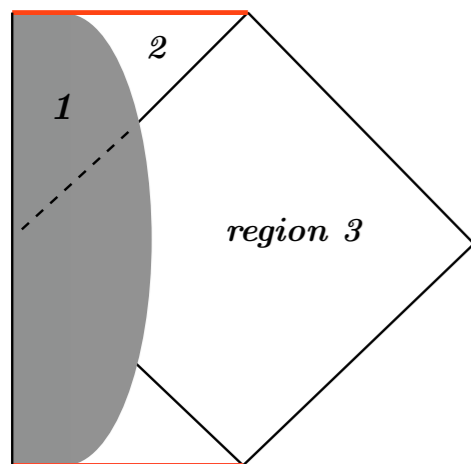


Fig. 2b



- Frolov, Vilkovinski '79
- Stephen, t'Hooft, Whiting '93
- Ashtekar, Bojowald '05
- Modesto '06
- Hayward '06
- Hajicek Kieffer '01
- Haggard, Rovelli '15

Covariant loop quantum gravity. Full definition.

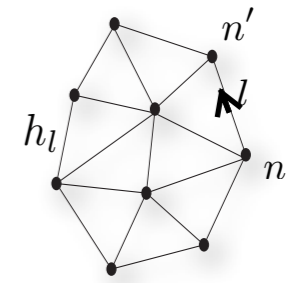
Kinematics
Boundary

State space

$$\mathcal{H}_\Gamma = L^2[SU(2)^L / SU(2)^N]_\Gamma \ni \psi(h_l) \quad \mathcal{H} = \lim_{\Gamma \rightarrow \infty} \mathcal{H}_\Gamma$$

Operators:

$$\vec{L}_l = \{L_l^i\}, i = 1, 2, 3 \quad \text{where } L_l^i \psi(h) \equiv \left. \frac{d}{dt} \psi(h e^{t\tau_i}) \right|_{t=0}$$



Γ spin network (nodes, links)

Dynamics
Bulk

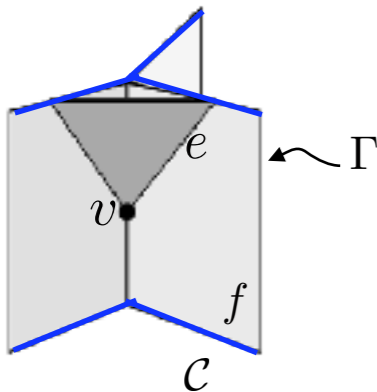
Transition amplitudes

$$W_C(h_l) = N_C \int_{SU(2)} dh_{vf} \prod_f \delta(h_f) \prod_v A(h_{vf})$$

$$h_f = \prod_v h_{vf}$$

Vertex amplitude

$$A(h_{vf}) = \int_{SL(2, \mathbb{C})} dg'_e \prod_f \sum_j (2j+1) D_{mn}^j(h_{vf}) D_{jmjn}^{\gamma(j+1)j}(g_e g_{e'}^{-1})$$



spinfoam (vertices, edges, faces)

$$W = \lim_{C \rightarrow \infty} W_C$$

$$8\pi\gamma\hbar G = 1$$



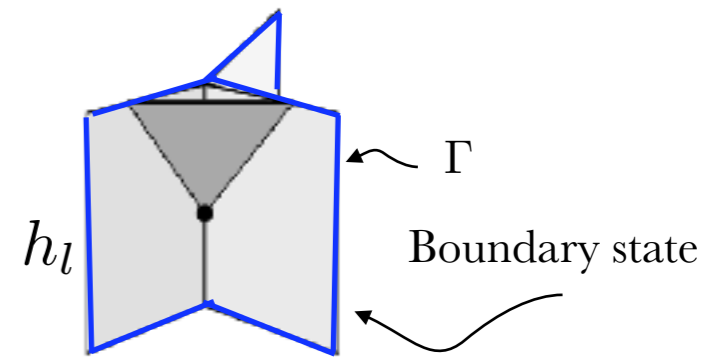
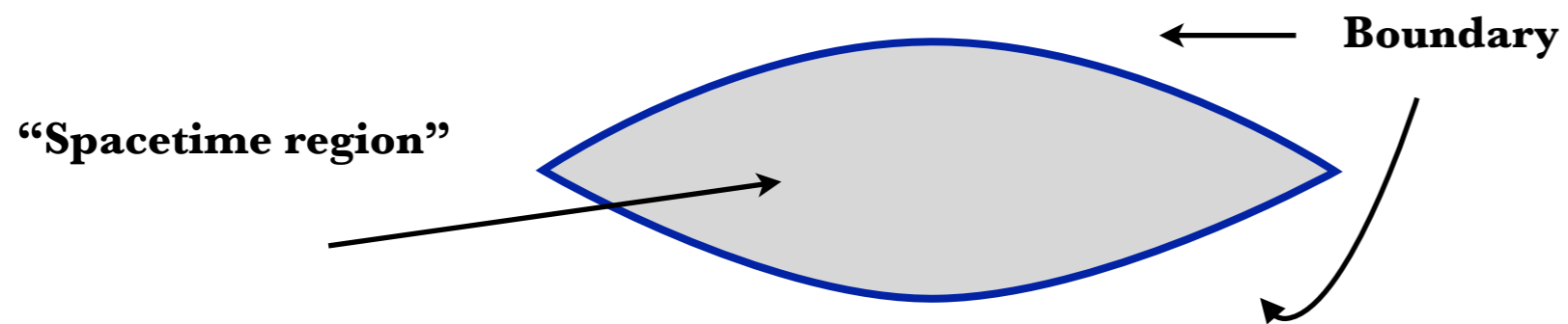
Satisfies all “requirements” listed by Steve Carlip in the last talk

- **Quantum theory** (Hilbert space, operators, transition amplitudes)
- **Classical limit: GR** (Barrett’s theorems)
- **UV qft divergences** (Han and Fairbairn’s theorems)
- **Minimum length** (Spectral properties of operators)
- **Black hole thermodynamics**

A process and its amplitude

Boundary state $\Psi = \psi_{in} \otimes \psi_{out}$

Amplitude $A = W(\Psi)$



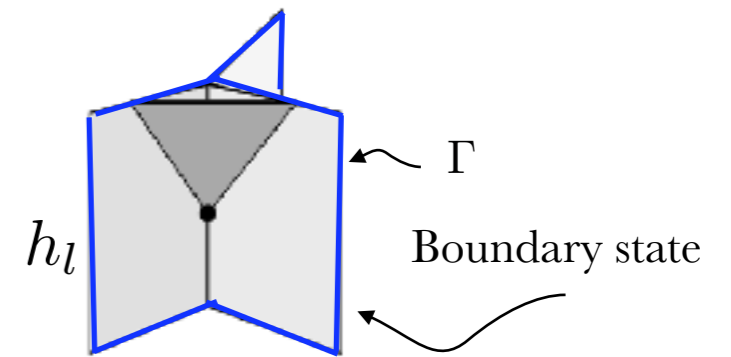
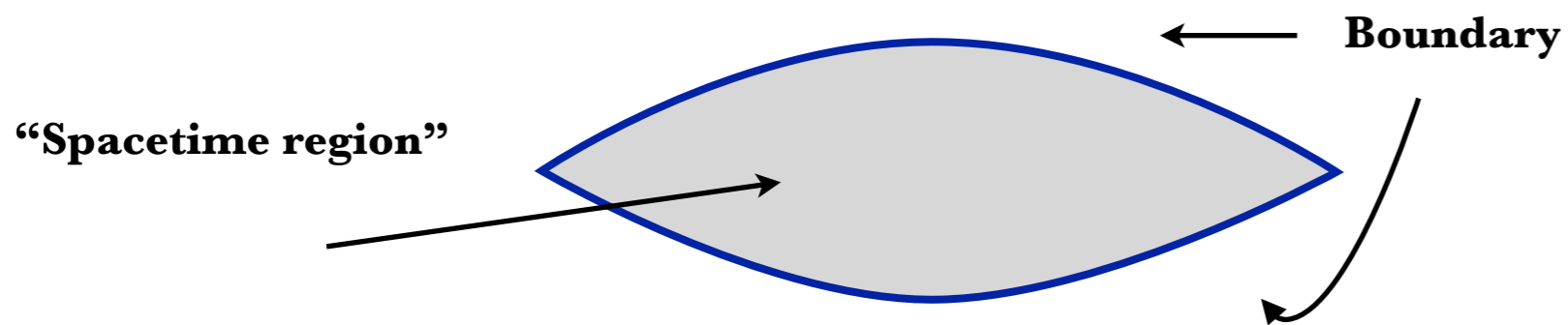
Quantum system
=
Spacetime region

→ Hamilton function: $S(q,t,q',t')$

A process and its amplitude

Boundary state $\Psi = \psi_{in} \otimes \psi_{out}$

Amplitude $A = W(\Psi)$



Quantum system
=
Spacetime region

→ Hamilton function: $S(q,t,q',t')$

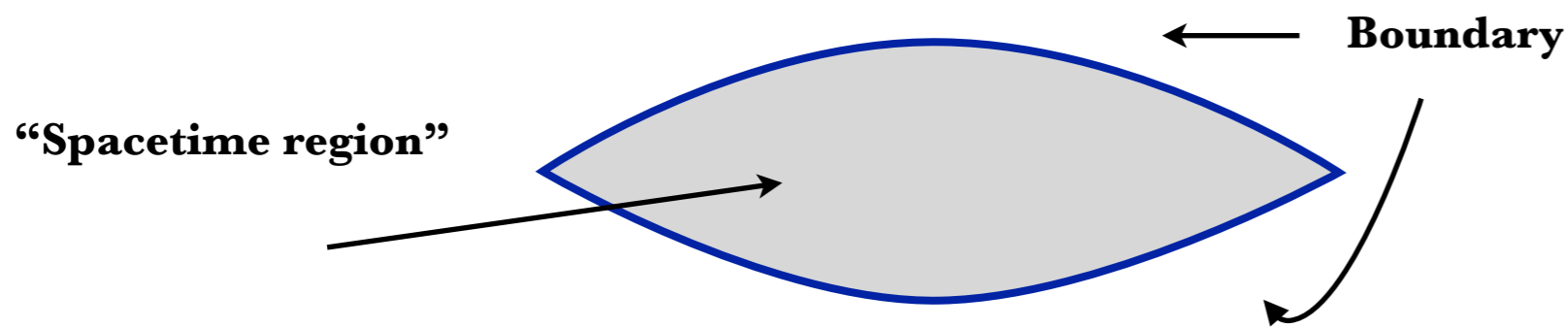
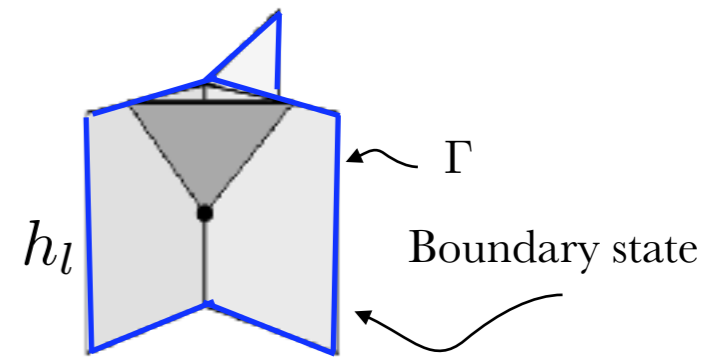
In GR, distance and time measurements are field measurements like any other one: they are part of the **boundary data** of the problem

Boundary values of the gravitational field = geometry of box surface = distance and time separation of measurements

A process and its amplitude

Boundary state $\Psi = \psi_{in} \otimes \psi_{out}$

Amplitude $A = W(\Psi)$



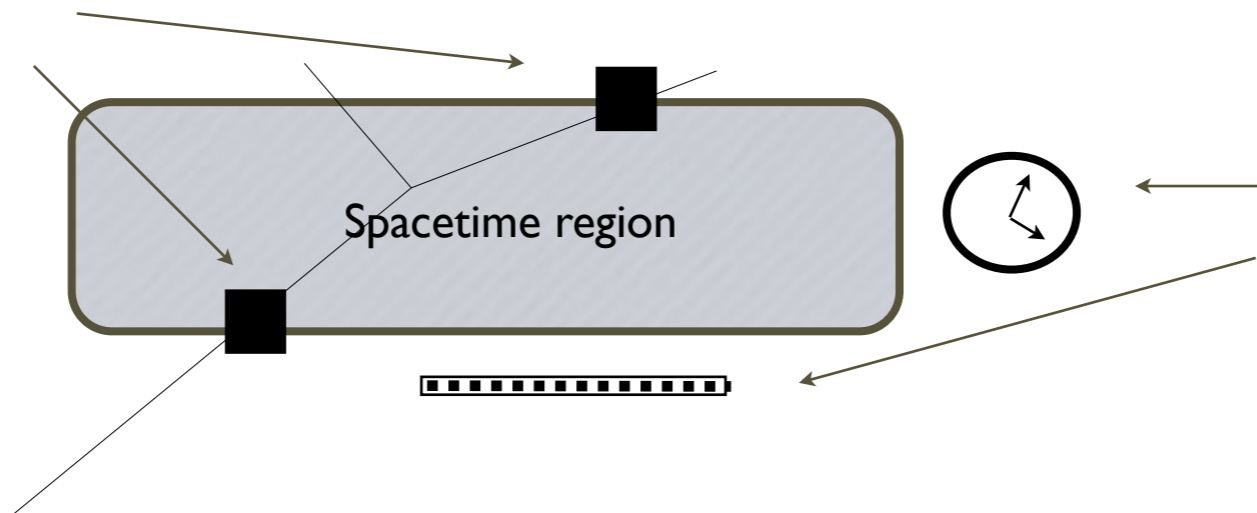
**Quantum system
=
Spacetime region**

→ Hamilton function: $S(q,t,q',t')$

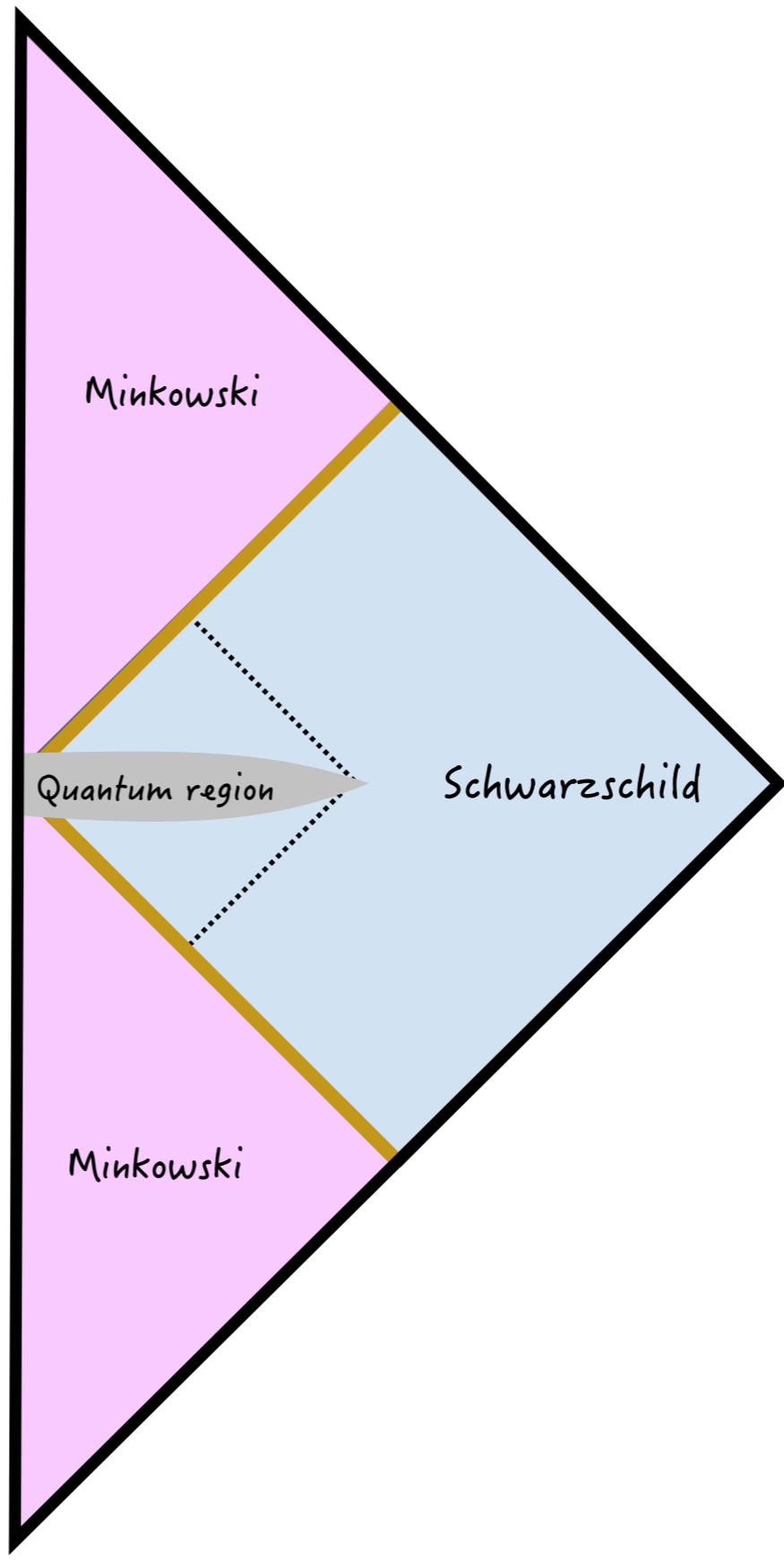
In GR, distance and time measurements are field measurements like any other one: they are part of the **boundary data** of the problem

Boundary values of the gravitational field = geometry of box surface = distance and time separation of measurements

Particle detectors
= field measurements



Distance and time measurements
= gravitational field measurements



Full expression for the transition amplitude:

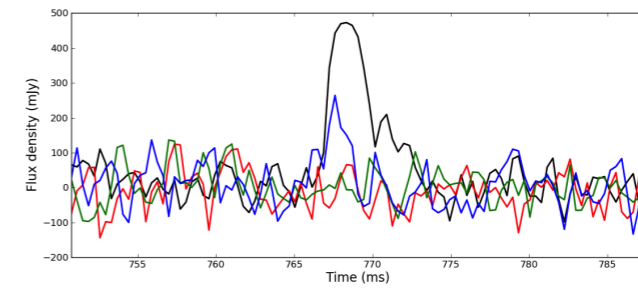
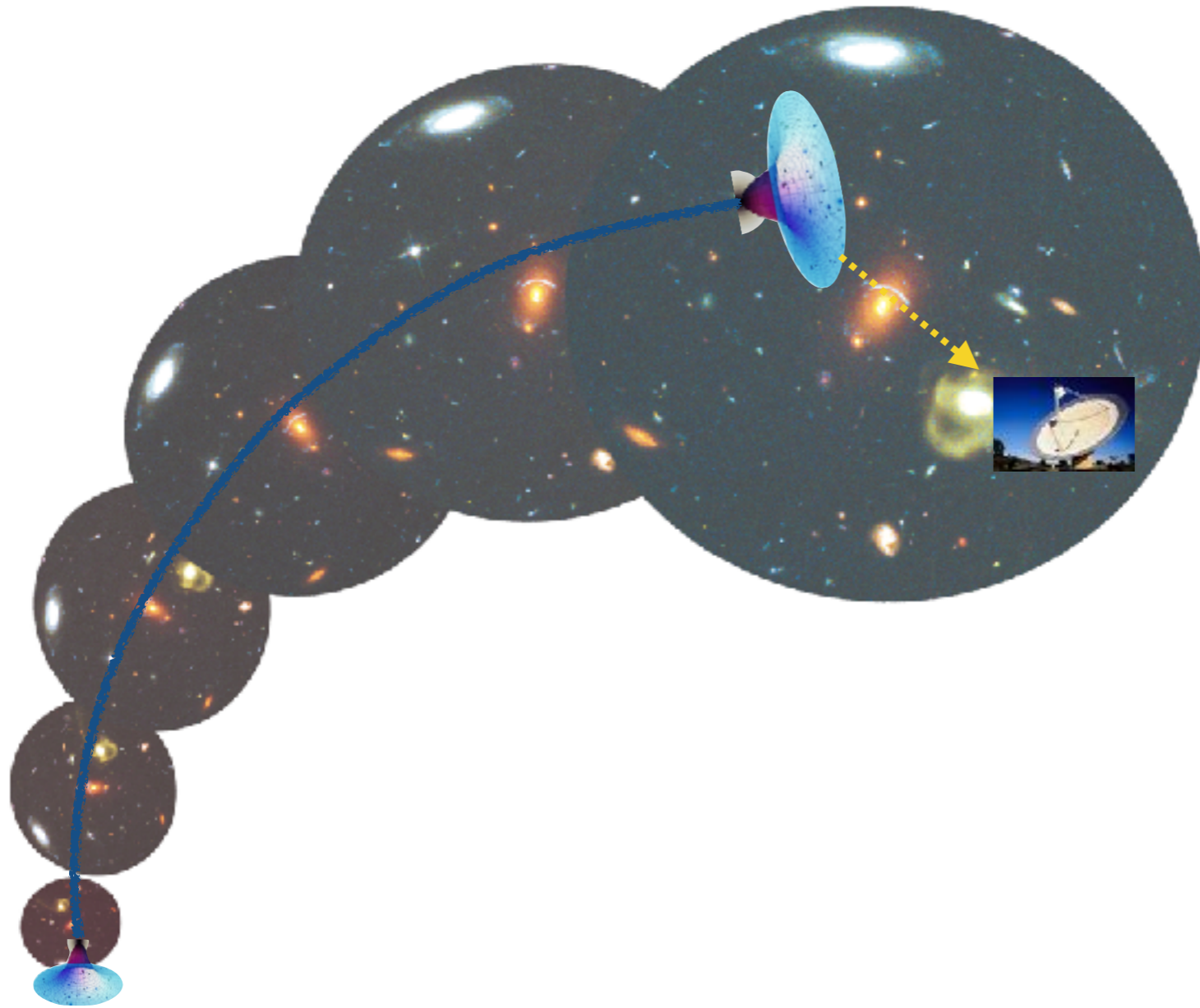
$$W(m, T) = \sum_{\{j_\ell\}} w(m, T, j_\ell) \sum_{\{J_n\}, \{K_n\}, \{l_\ell\}} \left(\bigotimes_n N_{\{j_n\}}^{J_n}(\{\nu_n\}, \{\alpha_n\}) f_{\{j_n\}\{l_n\}}^{J_n, K_n} \right) \left(\bigotimes_n i^{K_n, \{l_n\}} \right)_\Gamma$$

$$w(m, T, j_\ell) = c(m) \prod_\ell d_{j_\ell} e^{-\frac{1}{2\eta_\ell} \left(j_\ell - \frac{(2\eta_\ell^2 - 1)}{2} \right)^2} e^{i\gamma\zeta_\ell j_\ell}, \quad \eta_\ell^2 \sim m^2$$

$$N_{\{j_n\}}^{J_n} = \left(\bigotimes_{\ell \in n} \overleftarrow{D}_{m_\ell j_\ell}^{j_\ell}(\{\nu_n\}, \{\alpha_n\}) \right) i^{J_n, \{j_n\}}_{\{\vec{m}_n\}}$$

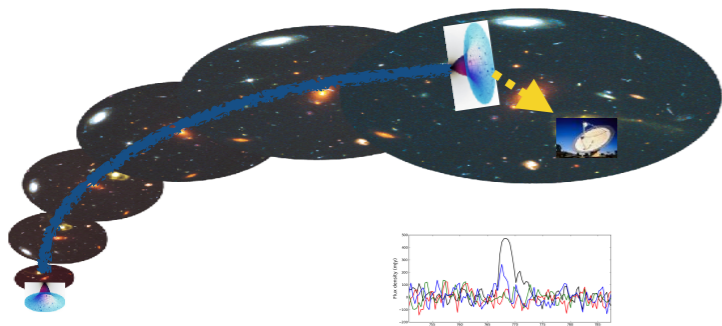
$$f_{\{j_n\}\{l_n\}}^{K_n, J_n} \equiv d_{J_n} i^{J_n, \{j_n\}}_{\{\vec{p}_n\}} \left(\int dr_n \frac{\sinh^2 r_n}{4\pi} \bigotimes_{\ell \in n} \overrightarrow{d}_{j_\ell l_\ell p_\ell}(r_n) \right) i^{K_n, \{l_n\}}_{\{\vec{p}_n\}} d_{K_n}$$

Primordial black holes!

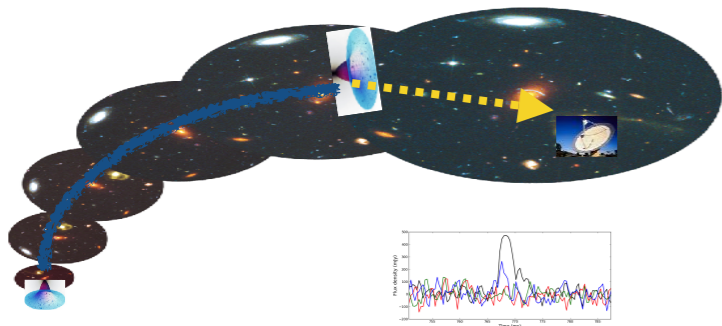


Signature: distance/energy relation

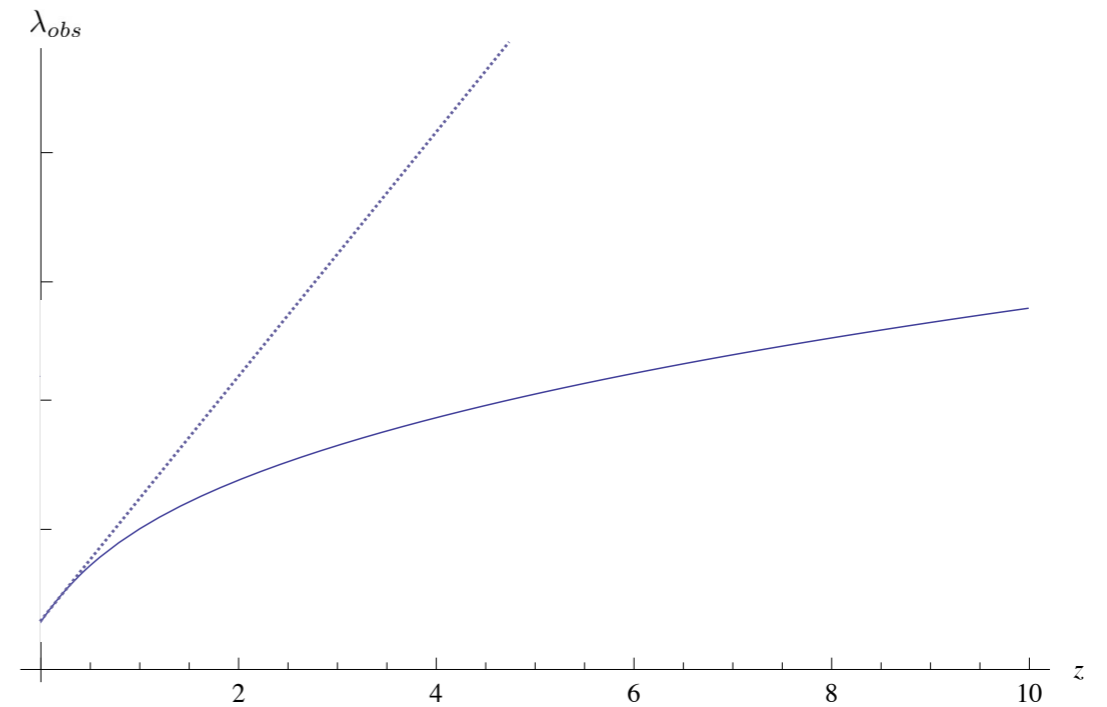
$$\lambda_{obs} \sim \frac{2Gm}{c^2} (1+z) \sqrt{\frac{H_0^{-1}}{6k\Omega_\Lambda^{1/2}} \sinh^{-1} \left[\left(\frac{\Omega_\Lambda}{\Omega_M} \right)^{1/2} (z+1)^{-3/2} \right]}$$



m_1



$m_2 < m_1$



Fundamental Physics with the SKA: Dark Matter and Astroparticles

K. Kelley¹, S. Riemer-Sørensen², E. Athanassoula³, C. Boehm^{4,5}, G. Bertone⁶, A. Bosma⁷, M. Brüggen⁸, C. Burigana^{9,10,11}, F. Calore^{6,12}, S. Camera^{13,14,15}, J.A.R. Cembranos¹⁶, R.M.T. Connors¹⁷, Á. de la Cruz-Dombriz^{18,19}, P.K.S Dunsby², N. Fornengo^{13,14,15}, D. Gaggero⁶, M. Méndez-Isla^{18,19}, Y. Ma^{21,22,23}, H. Padmanabhan²⁴, A. Pourtsidou²⁵, P.J. Quinn¹, M. Regis^{13,14}, M. Sahlén²⁶, M. Sakellariadou²⁷, L. Shao²⁸, J. Silk^{29,30,31,32}, T. Trombetti^{10,33,9}, F. Vazza^{34,35}, F. Vidotto³⁶, F. Villaescusa-Navarro³⁷, C. Weniger⁶ and L. Wolz³⁸

3.9.2 PBHs and Quantum Gravity

To date the search for PBHs and for constraints on PBHs has mostly considered an almost monochromatic mass spectrum, and the presence of Hawking evaporation for PBHs of small mass. Monochromatic mass spectrum has been challenged by different authors as unrealistic (for example Carr et al. (2017a)). An extended mass function is compatible with different PBHs formation mechanics, from critical collapse to cosmic strings. Hawking evaporation is a phenomenon that becomes relevant on a time scale that depends on the mass of the BH. Its time scale is M_{BH}^3 in Planck units. This implies that within the age of the universe only PBHs with mass smaller than 10^{12} kg could have evaporated, and possibly produced very high-energy cosmic rays (Barrau, 2000). As cosmic rays of such energies are rare, constraints are derived on the very-small-mass end of the PBH mass spectrum.

Hawking evaporation, however, is a phenomenon predicted in the context of quantum field theory on a fixed curved background. This is a theory with a regime of validity that may likely break down when approximately half of the mass of the hole has evaporated, as indicated for instance by the ‘firewall’ no-go theorem (Almheiri et al., 2013). The geometry around a BH can indeed undergo quantum fluctuations on a time scale shorter than M_{BH}^3 , when the effects of the Hawking evaporation have not yet significantly modified the size of the hole. As any classical system, the hole has a characteristic timescale after which the the departure from

of this discreteness on the dynamics can be modeled at the effective level by an effective potential that prevents the gravitational collapse from forming the singularity and triggers a bounce. The bounce connects a collapsing solution of the Einstein equation, that is the classical black hole, to an explosive expanding one, a white hole (Haggard & Rovelli, 2014), through an intermediate quantum region. This process is a typical quantum tunnelling event, and the characteristic time at which it takes place, the hole lifetime, can be as a decaying time, similar to the lifetime of conventional nuclear radioactivity. The resulting picture is conservative in comparison to other models of non-singular BHs. The collapse still produces a horizon, but it is now a dynamical horizon with a finite lifetime, rather than a perpetual event horizon. The collapsing matter continues its fall after entering the trapping region, forming a very dense object whose further collapse is prevented by quantum pressure (referred to with the suggesting name of *Planck Star* Rovelli & Vidotto (2014)).

The collapsing matter that forms PBHs in the radiation dominated epoch is mainly constituted by photons. Seen from the center of the hole, those photons collapse through the trapping region, then expands passing through an anti-trapping region and eventually exits the white-hole horizon, always at the speed of light, the process is thus extremely fast. On the other hand, for an observer sitting outside the horizon, a huge but finite redshift stretches this time to cosmological times. This time, properly called the hole lifetime, as discussed before has

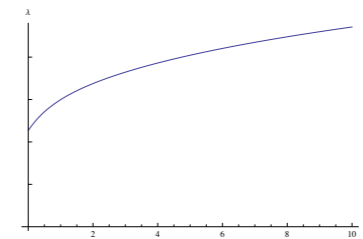
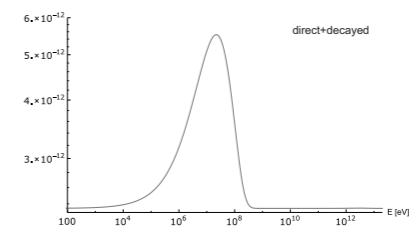
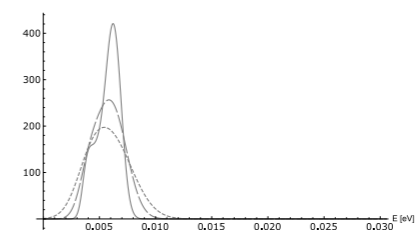


Figure 21. The expected wavelength (unspecified units) of the

DM and Astroparticles



Quantum Gravity observations are not absurd anymore.

- Decrease the Bayesian confidence of some current theories (absence of supersymmetry at energy where it was expected)
- Already rules out theories (Lorentz invariance)
- Suggested astrophysical observations motivating astronomers (Cosmology+Black holes)
- Interesting laboratory experiments (Entanglement via gravity)

(Result of a poll at a recent conference (3rd Karl Schwarzschild Meeting on Gravitational Physics and the Gauge/Gravity Correspondence, Frankfurt am Main, July 2017): 80% of the participants expect observational evidence for quantum gravity observations within the next decade.)

Progress has happened along some research directions


	Empirical success	Empirical failure	Theoretical success	Theoretical failure	Key open issue
Strings	None	Low-energy super symmetry	AdS-CFT Non-fundamental applications	Standard model parameters	Lack of predictivity
Loops	None		Low energy limit		Infinite graph limit
Hojava-Lifshitz	None	Lorentz violation at Planck scale	Renormalizability		Other scale?
Asymptotic savety	None		Increase evidence of fixed point		Computing amplitudes
Your favorite

**Loops
open problems
in the late 80'**



Loops open problems in the late 80'

- Definition of the Hilbert space
- Mathematical foundation
- Coupling matter
- Recovering low energy GR
- Problem of time
- Observables
- Application to early universe
- Produce verifiable physical predictions




Loops open problems in the late 80'

- Definition of the Hilbert space 
- Mathematical foundation
- Coupling matter
- Recovering low energy GR
- Problem of time
- Observables
- Application to early universe
- Produce verifiable physical predictions





Loops open problems in the late 80'

- Definition of the Hilbert space 
- Mathematical foundation 
- Coupling matter
- Recovering low energy GR
- Problem of time
- Observables
- Application to early universe
- Produce verifiable physical predictions






Loops open problems in the late 80'

- Definition of the Hilbert space 
- Mathematical foundation 
- Coupling matter 
- Recovering low energy GR
- Problem of time
- Observables
- Application to early universe
- Produce verifiable physical predictions







Loops open problems in the late 80'

- Definition of the Hilbert space 
- Mathematical foundation 
- Coupling matter 
- Recovering low energy GR 
- Problem of time
- Observables
- Application to early universe
- Produce verifiable physical predictions

Loops open problems in the late 80'

- Definition of the Hilbert space 
- Mathematical foundation 
- Coupling matter 
- Recovering low energy GR 
- Problem of time 
- Observables
- Application to early universe
- Produce verifiable physical predictions









Loops open problems in the late 80'

- Definition of the Hilbert space 
- Mathematical foundation 
- Coupling matter 
- Recovering low energy GR 
- Problem of time 
- Observables 
- Application to early universe
- Produce verifiable physical predictions

Loops open problems in the late 80'

- Definition of the Hilbert space ✓
- Mathematical foundation ✓
- Coupling matter ✓
- Recovering low energy GR ✓
- Problem of time ✓
- Observables ✓
- Application to early universe ✓
- Produce verifiable physical predictions

Loops open problems in the late 80'

- Definition of the Hilbert space 
- Mathematical foundation 
- Coupling matter 
- Recovering low energy GR 
- Problem of time 
- Observables 
- Application to early universe 
- Produce verifiable physical predictions 

Loops open problems in the late 80'

- Definition of the Hilbert space
- Mathematical foundation
- Coupling matter
- Recovering low energy GR
- Problem of time
- Observables
- Application to early universe
- Produce verifiable physical predictions



Strings open problems in the late 80'

Loops open problems in the late 80'

- Definition of the Hilbert space ✓
- Mathematical foundation ✓
- Coupling matter ✓
- Recovering low energy GR ✓
- Problem of time ✓
- Observables ✓
- Application to early universe ✓
- Produce verifiable physical predictions ✗

Strings open problems in the late 80'

- Finding a fundamental formulation of the theory
- Deriving $SU(3) \times SU(2) \times U(1)$ from first principles
- Computing the parameters of the standard model
- Supersymmetry breaking
- Compactification
- Extend to the non-perturbative regime
- Produce tentative verifiable physical predictions

Loops open problems in the late 80'

- Definition of the Hilbert space ✓
- Mathematical foundation ✓
- Coupling matter ✓
- Recovering low energy GR ✓
- Problem of time ✓
- Observables ✓
- Application to early universe ✓
- Produce verifiable physical predictions ✗

Strings open problems in the late 80'

- Finding a fundamental formulation of the theory ✗
- Deriving $SU(3) \times SU(2) \times U(1)$ from first principles
- Computing the parameters of the standard model
- Supersymmetry breaking
- Compactification
- Extend to the non-perturbative regime
- Produce tentative verifiable physical predictions

Loops open problems in the late 80'

- Definition of the Hilbert space ✓
- Mathematical foundation ✓
- Coupling matter ✓
- Recovering low energy GR ✓
- Problem of time ✓
- Observables ✓
- Application to early universe ✓
- Produce verifiable physical predictions ✗

Strings open problems in the late 80'

- Finding a fundamental formulation of the theory ✗
- Deriving $SU(3) \times SU(2) \times U(1)$ from first principles ✗
- Computing the parameters of the standard model
- Supersymmetry breaking
- Compactification
- Extend to the non-perturbative regime
- Produce tentative verifiable physical predictions

Loops open problems in the late 80'

- Definition of the Hilbert space ✓
- Mathematical foundation ✓
- Coupling matter ✓
- Recovering low energy GR ✓
- Problem of time ✓
- Observables ✓
- Application to early universe ✓
- Produce verifiable physical predictions ✗

Strings open problems in the late 80'

- Finding a fundamental formulation of the theory ✗
- Deriving $SU(3) \times SU(2) \times U(1)$ from first principles ✗
- Computing the parameters of the standard model ✗
- Supersymmetry breaking
- Compactification
- Extend to the non-perturbative regime
- Produce tentative verifiable physical predictions

Loops open problems in the late 80'

- Definition of the Hilbert space ✓
- Mathematical foundation ✓
- Coupling matter ✓
- Recovering low energy GR ✓
- Problem of time ✓
- Observables ✓
- Application to early universe ✓
- Produce verifiable physical predictions ✗

Strings open problems in the late 80'

- Finding a fundamental formulation of the theory ✗
- Deriving $SU(3) \times SU(2) \times U(1)$ from first principles ✗
- Computing the parameters of the standard model ✗
- Supersymmetry breaking ✗
- Compactification
- Extend to the non-perturbative regime
- Produce tentative verifiable physical predictions

Loops open problems in the late 80'

- Definition of the Hilbert space ✓
- Mathematical foundation ✓
- Coupling matter ✓
- Recovering low energy GR ✓
- Problem of time ✓
- Observables ✓
- Application to early universe ✓
- Produce verifiable physical predictions ✗

Strings open problems in the late 80'

- Finding a fundamental formulation of the theory ✗
- Deriving $SU(3) \times SU(2) \times U(1)$ from first principles ✗
- Computing the parameters of the standard model ✗
- Supersymmetry breaking ✗
- Compactification ✗
- Extend to the non-perturbative regime
- Produce tentative verifiable physical predictions

Loops open problems in the late 80'

- Definition of the Hilbert space ✓
- Mathematical foundation ✓
- Coupling matter ✓
- Recovering low energy GR ✓
- Problem of time ✓
- Observables ✓
- Application to early universe ✓
- Produce verifiable physical predictions ✗

Strings open problems in the late 80'

- Finding a fundamental formulation of the theory ✗
- Deriving $SU(3) \times SU(2) \times U(1)$ from first principles ✗
- Computing the parameters of the standard model ✗
- Supersymmetry breaking ✗
- Compactification ✗
- Extend to the non-perturbative regime ✓
- Produce tentative verifiable physical predictions ✗

Loops open problems in the late 80'

- Definition of the Hilbert space ✓
- Mathematical foundation ✓
- Coupling matter ✓
- Recovering low energy GR ✓
- Problem of time ✓
- Observables ✓
- Application to early universe ✓
- Produce verifiable physical predictions ✗

Strings open problems in the late 80'

- Finding a fundamental formulation of the theory ✗
- Deriving $SU(3) \times SU(2) \times U(1)$ from first principles ✗
- Computing the parameters of the standard model ✗
- Supersymmetry breaking ✗
- Compactification ✗
- Extend to the non-perturbative regime ✓
- Produce tentative verifiable physical predictions ✓

Loops open problems in the late 80'

- Definition of the Hilbert space ✓
- Mathematical foundation ✓
- Coupling matter ✓
- Recovering low energy GR ✓
- Problem of time ✓
- Observables ✓
- Application to early universe ✓
- Produce verifiable physical predictions ✗

Strings open problems in the late 80'

- Finding a fundamental formulation of the theory ✗
- Deriving $SU(3) \times SU(2) \times U(1)$ from first principles ✗
- Computing the parameters of the standard model ✗
- Supersymmetry breaking ✗
- Compactification ✗
- Extend to the non-perturbative regime ✓
- Produce tentative verifiable physical predictions ✓ ✗

3. There is some convergence on some issues of major philosophical relevance

i: nature of physical space

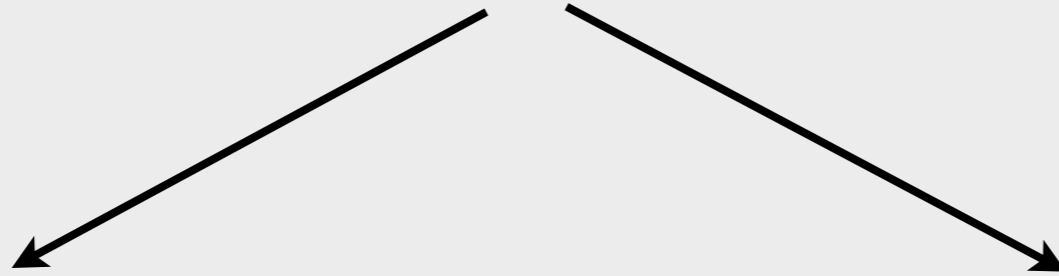
ii: nature of physical time

iii: relation between physical and experiential time

iv: connection between relational aspects of GR and QM

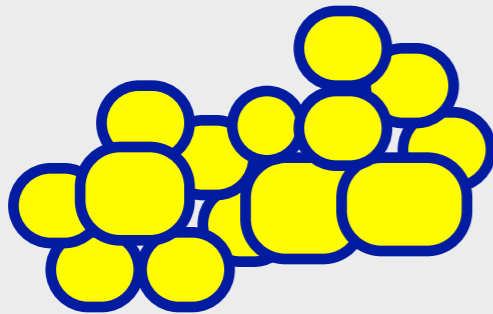
i: nature of physical space

space



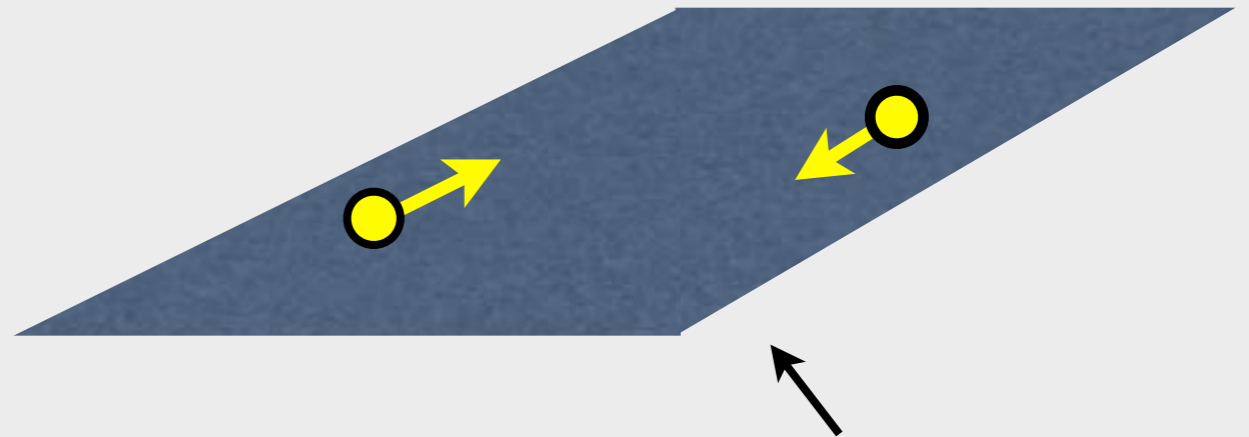
relation

Aristotle
Descartes



entity

Newton



Space is an entity independent
from its content

Aristotle :

Substance

Aristotle :

Substance

Descartes:

Res extensa

Aristotle :

Substance

Descartes:

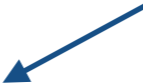
Res extensa

Newton:

Time

Space

Particles



Aristotle :

Substance

Descartes:

Res extensa

Newton:

Time

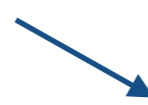
Space

Particles

Faraday-Maxwell:

Fields

Particles



Aristotle :

Substance

Descartes:

Res extensa

Newton:

Time

Space

Particles

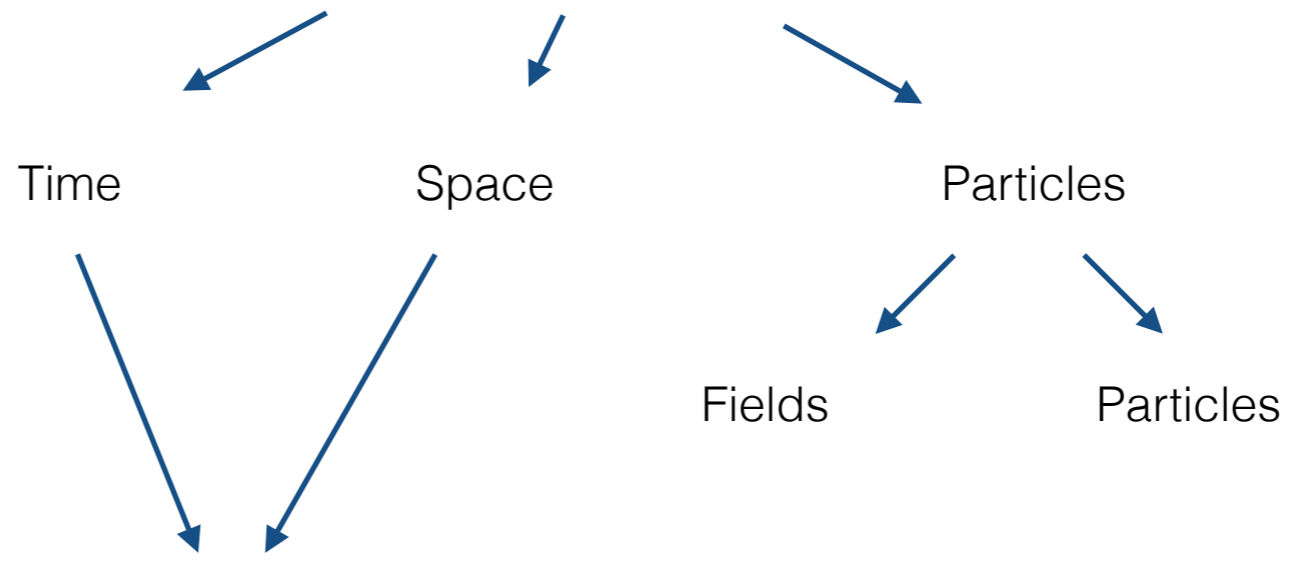
Faraday-Maxwell:

Fields

Particles

Einstein 05:

Spacetime



Aristotle :

Substance

Descartes:

Res extensa

Newton:

Time

Space

Particles

Faraday-Maxwell:

Fields

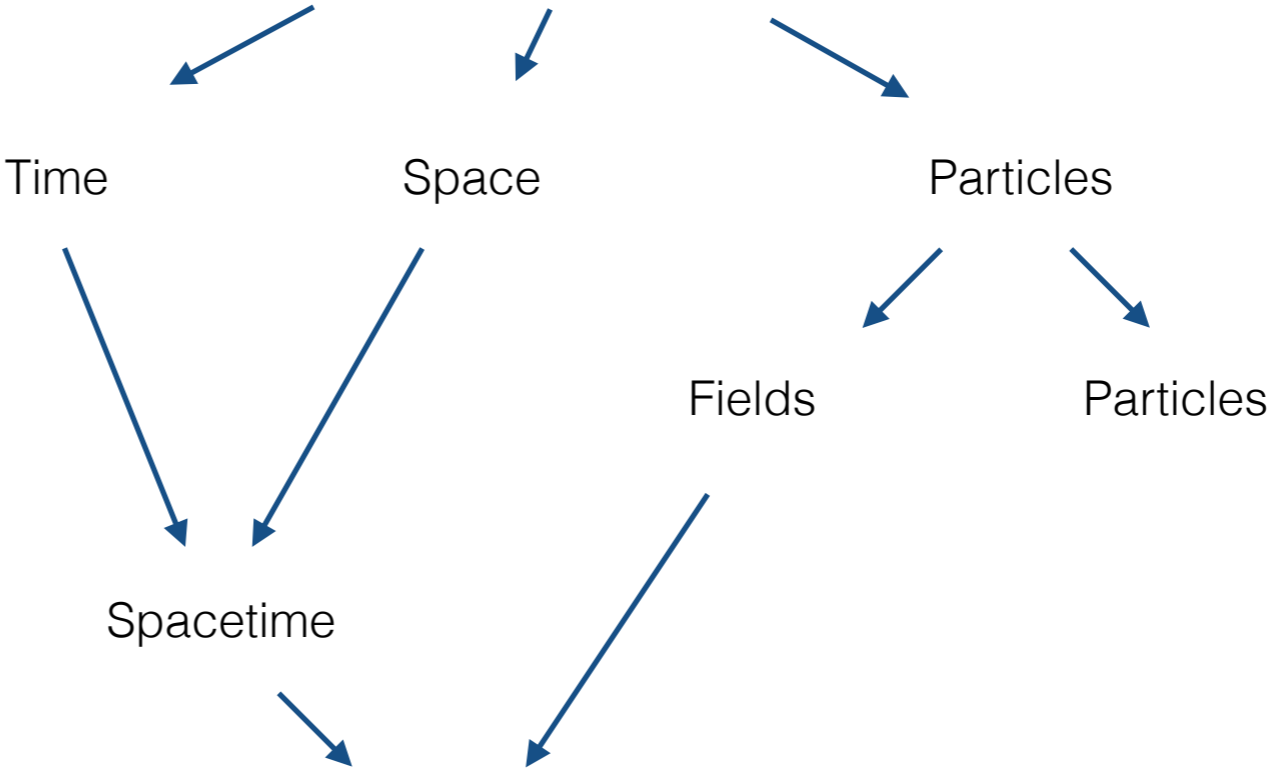
Particles

Einstein 05:

Spacetime

Einstein 15:

Covariant Fields



Aristotle :

Substance

Descartes:

Res extensa

Newton:

Time

Space

Particles

Faraday-Maxwell:

Fields

Particles

Einstein 05:

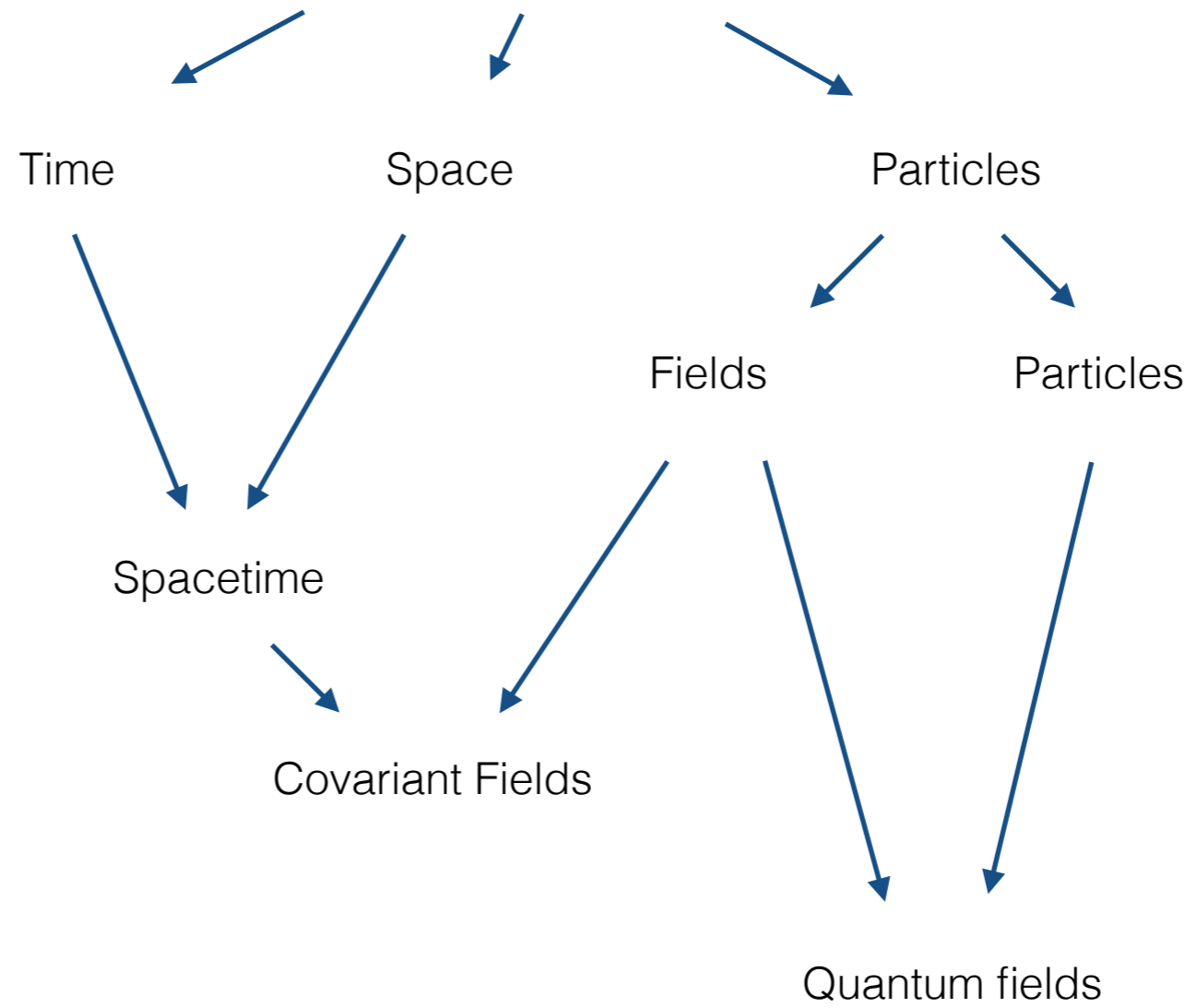
Spacetime

Einstein 15:

Covariant Fields

Quanta:

Quantum fields



Aristotle :

Substance

Descartes:

Res extensa

Newton:

Time

Space

Particles

Faraday-Maxwell:

Fields

Particles

Einstein 05:

Spacetime

Einstein 15:

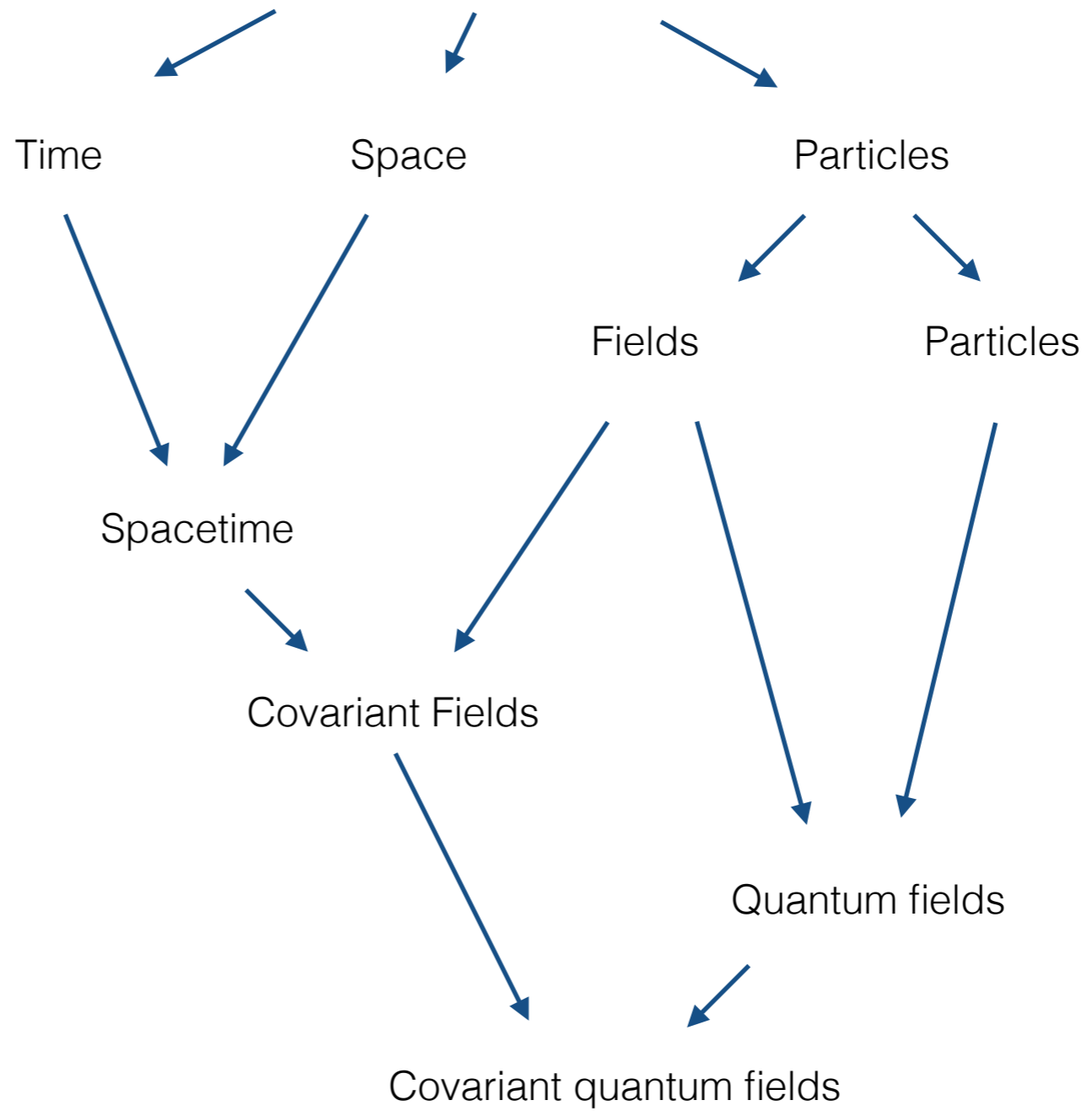
Covariant Fields

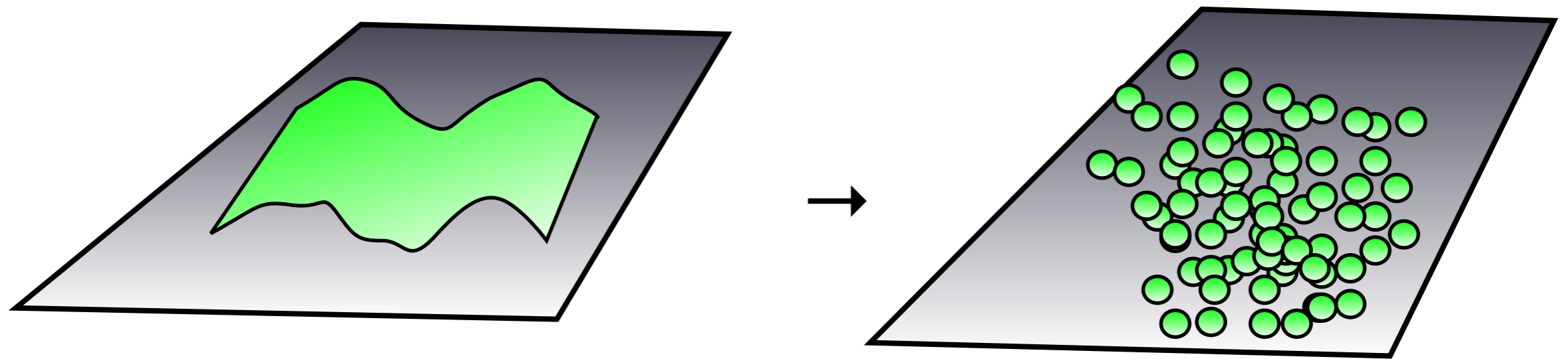
Quanta:

Quantum fields

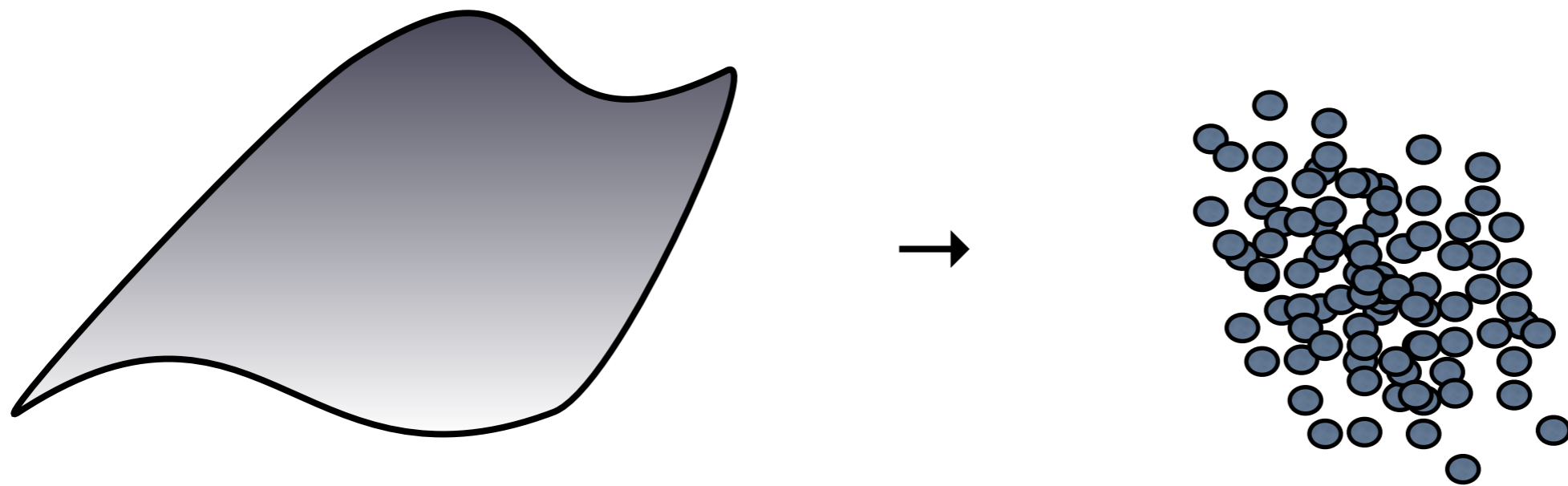
Quantum gravity:

Covariant quantum fields





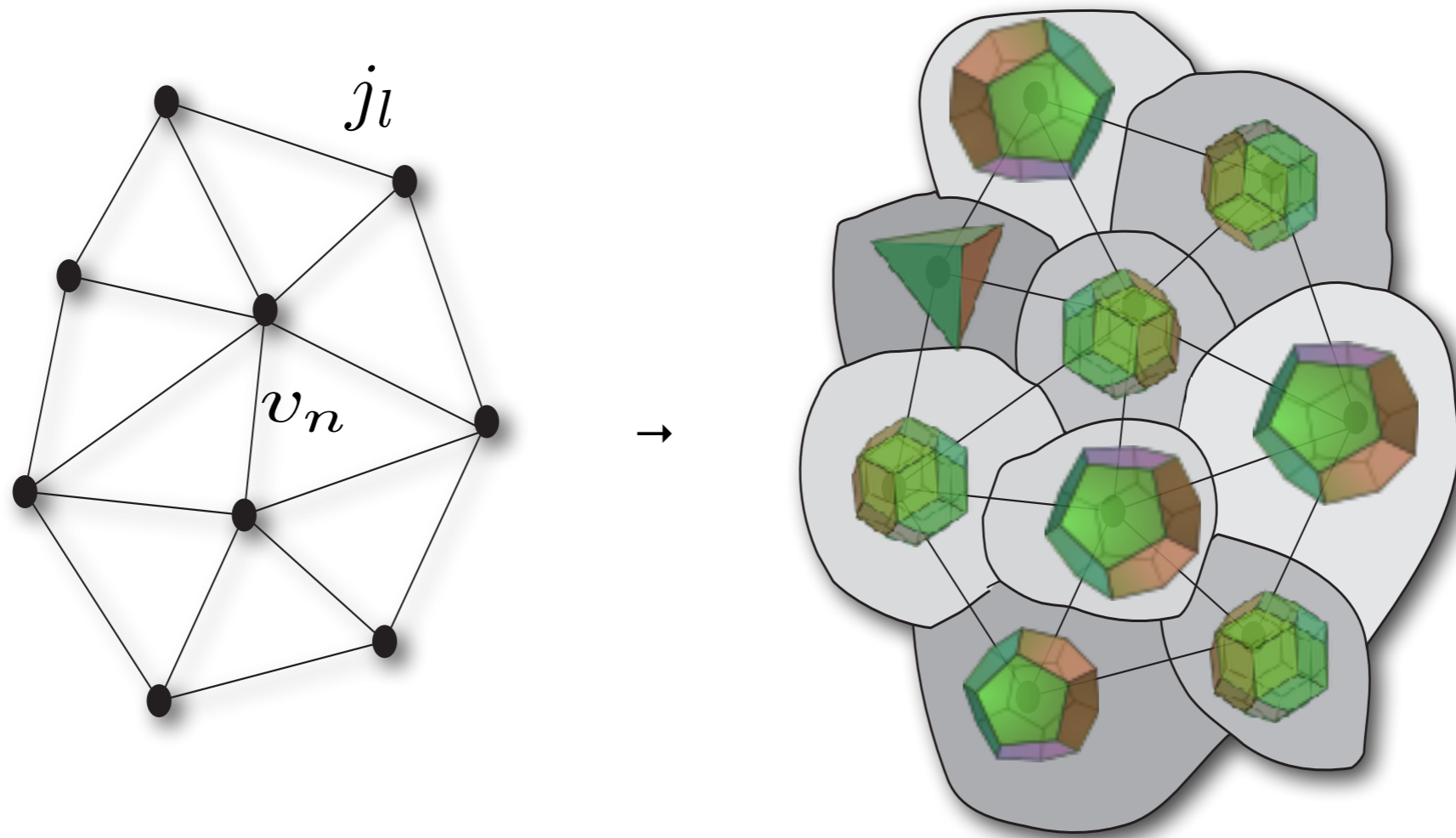
Quantum electromagnetic field: quanta of light



Quantum gravitational space: quanta of space

“Emergence” ←

physical space is granular



Space has a granular structure: spin networks

Covariant loop quantum gravity. Full definition.

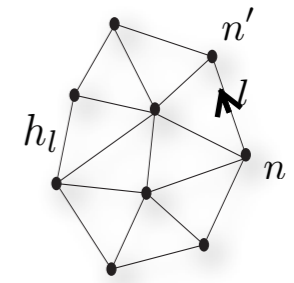
Kinematics
Boundary

State space

$$\mathcal{H}_\Gamma = L^2[SU(2)^L / SU(2)^N]_\Gamma \ni \psi(h_l) \quad \mathcal{H} = \lim_{\Gamma \rightarrow \infty} \mathcal{H}_\Gamma$$

Operators:

$$\vec{L}_l = \{L_l^i\}, i = 1, 2, 3 \quad \text{where } L_l^i \psi(h) \equiv \left. \frac{d}{dt} \psi(h e^{t\tau_i}) \right|_{t=0}$$



Γ spin network
(nodes, links)

Dynamics
Bulk

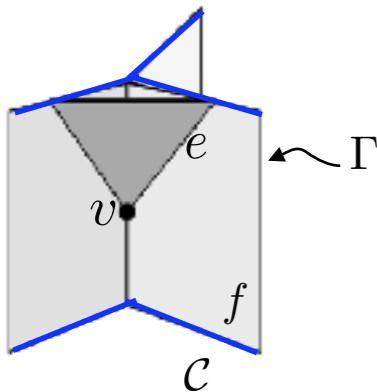
Transition amplitudes

$$W_C(h_l) = N_C \int_{SU(2)} dh_{vf} \prod_f \delta(h_f) \prod_v A(h_{vf})$$

$$h_f = \prod_v h_{vf}$$

Vertex amplitude

$$A(h_{vf}) = \int_{SL(2, \mathbb{C})} dg'_e \prod_f \sum_j (2j+1) D_{mn}^j(h_{vf}) D_{jmjn}^{\gamma(j+1)j}(g_e g_{e'}^{-1})$$



$$W = \lim_{C \rightarrow \infty} W_C$$

spinfoam
(vertices, edges, faces)

$$8\pi\gamma\hbar G = 1$$

ii: nature of physical time

Covariant loop quantum gravity. Full definition.

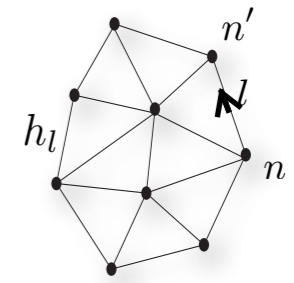
Kinematics
Boundary

State space

$$\mathcal{H}_\Gamma = L^2[SU(2)^L / SU(2)^N]_\Gamma \ni \psi(h_l) \quad \mathcal{H} = \lim_{\Gamma \rightarrow \infty} \mathcal{H}_\Gamma$$

Operators:

$$\vec{L}_l = \{L_l^i\}, i = 1, 2, 3 \quad \text{where } L^i \psi(h) \equiv \left. \frac{d}{dt} \psi(h e^{t\tau_i}) \right|_{t=0}$$



Γ spin network
(nodes, links)

Dynamics
Bulk

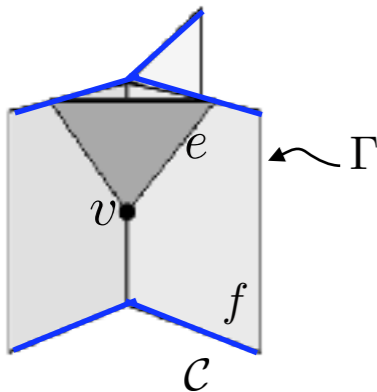
Transition amplitudes

$$W_C(h_l) = N_C \int_{SU(2)} dh_{vf} \prod_f \delta(h_f) \prod_v A(h_{vf})$$

$$h_f = \prod_v h_{vf}$$

Vertex amplitude

$$A(h_{vf}) = \int_{SL(2, \mathbb{C})} dg'_e \prod_f \sum_j (2j+1) D_{mn}^j(h_{vf}) D_{jmjn}^{\gamma(j+1)j}(g_e g_{e'}^{-1})$$

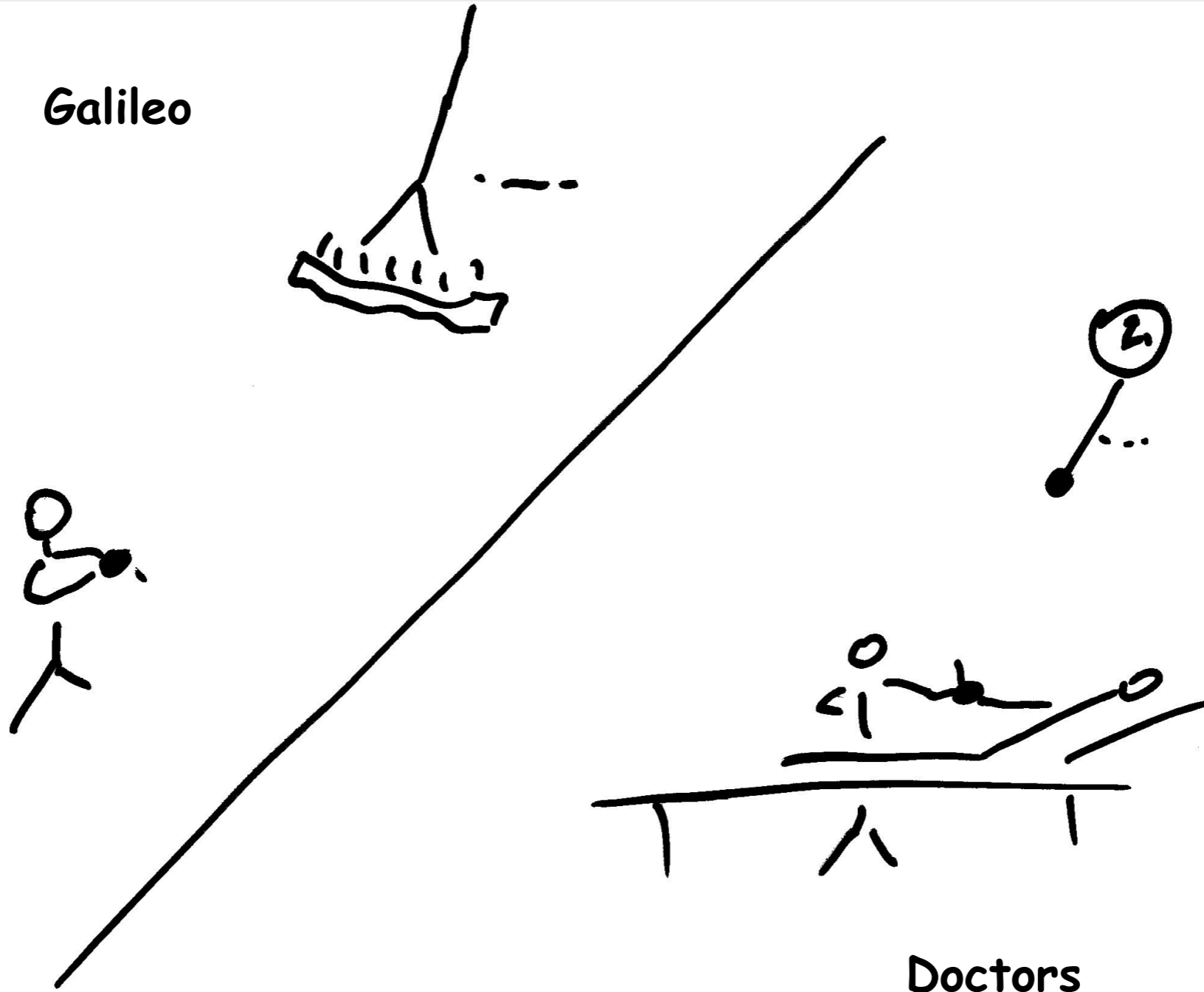


spinfoam
(vertices, edges, faces)

$$W = \lim_{C \rightarrow \infty} W_C$$

$$8\pi\gamma\hbar G = 1$$

Galileo



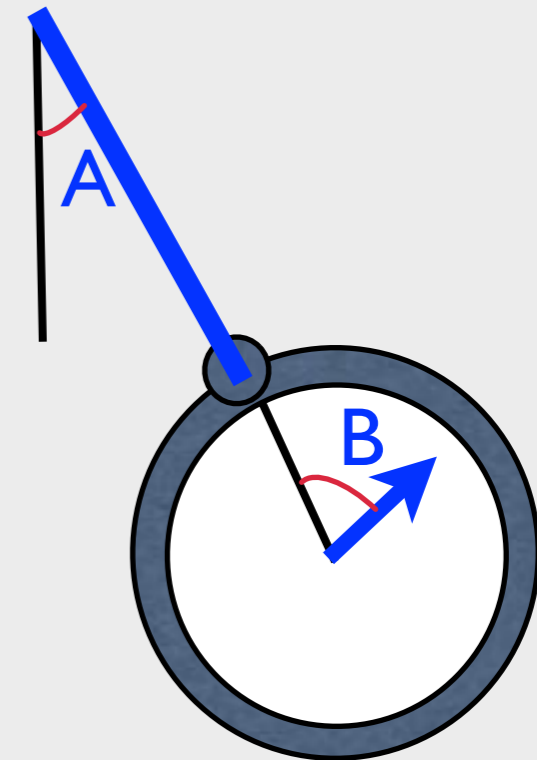
Doctors

Time: t

What we observe: A, B, \dots

Newton: $A(t), B(t), \dots$

But we actually only see: $A(B), B(A) \dots$



→ the world can be described without t

At the elementary level Nature is not organized in terms of evolution in time.

iii: relation between physical and experiential time

“Experiential Time” is a complex layer of structures and properties which appear (emerge) at different level of approximations and descriptions of the world:

Internal flow: Diff invariant qft

Connes’s flow of the vonN. algebra

Connes, A. & Rovelli, C. Von Neumann algebra automorphisms and time thermodynamics relation in general covariant quantum theories. *Class.Quant. Grav.* 11, 2899–2918 (1994).

Flow and energy: Statistical mechanics:

Thermal time (Tomita flow)

Unicity, present:

Non-relativistic limit

Orientation -> traces

(Perspectival?) low past entropy:

Price, H. *Time’s Arrow*. (Oxford University Press, 1996).

Rovelli, C. *Is Time’s Arrow Perspectival?* 6 (2015).

Replicating systems with memory:

Biology and evolution

“Sense of flowing time”, identity as narration

Brain’s memory and anticipation

D. Buonomano, *Your Brain is a Time Machine: The Neuroscience and Physics of Time*, Norton, New York, 2017.

Moral:

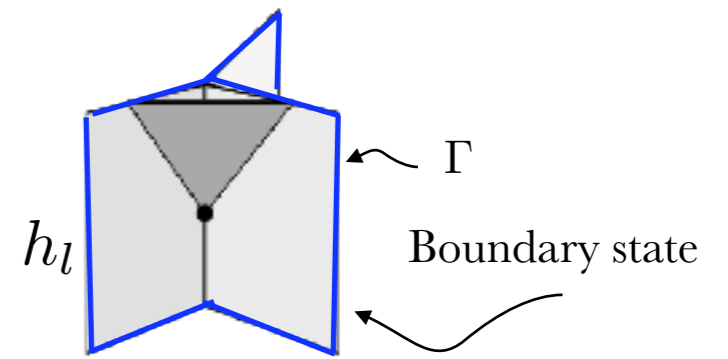
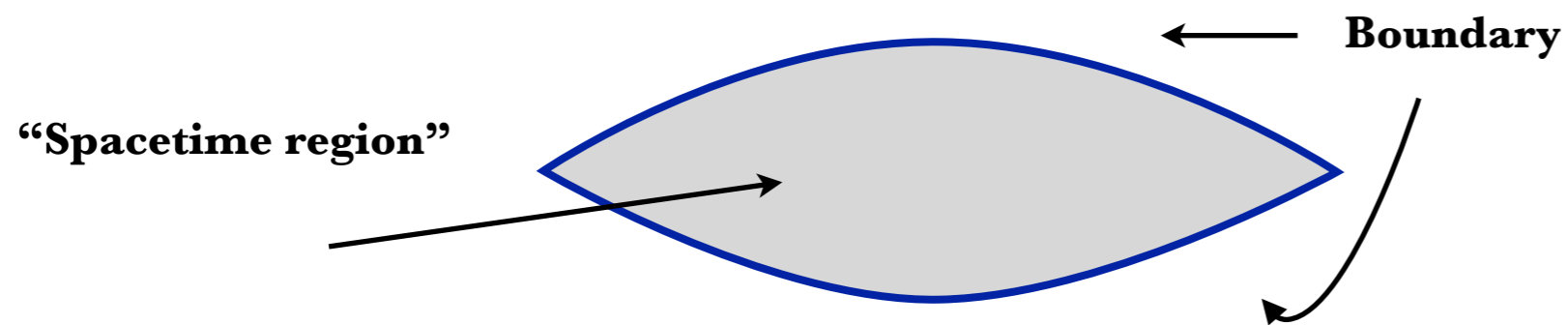
**do not search in elementary physics what is
not
in elementary physics**

iv: connection between relational aspects of GR and QM

A process and its amplitude

Boundary state $\Psi = \psi_{in} \otimes \psi_{out}$

Amplitude $A = W(\Psi)$



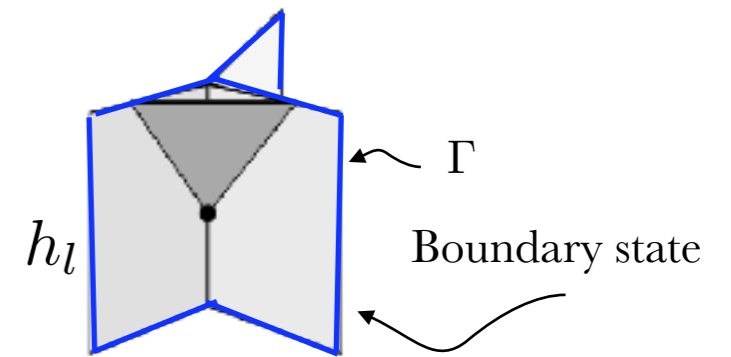
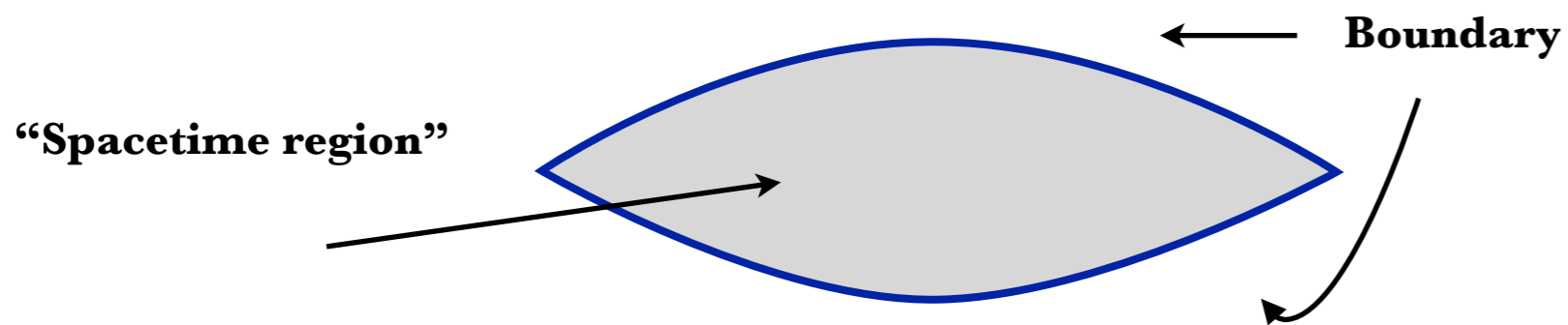
**Quantum system
=
Spacetime region**

→ Hamilton function: $S(q,t,q',t')$

A process and its amplitude

Boundary state $\Psi = \psi_{in} \otimes \psi_{out}$

Amplitude $A = W(\Psi)$



Quantum system
=
Spacetime region

→ Hamilton function: $S(q,t,q',t')$

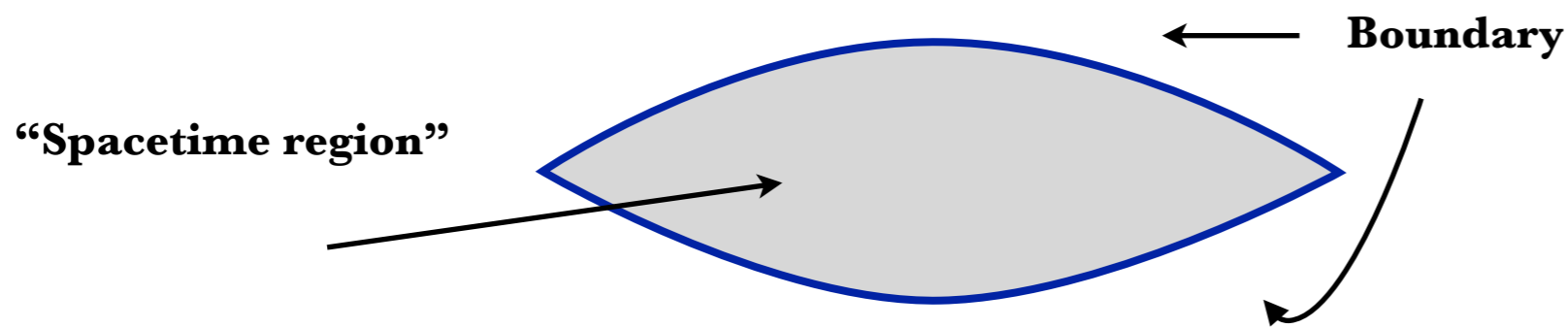
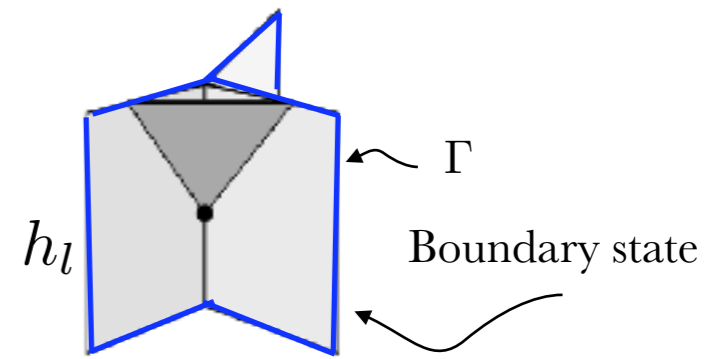
In GR, distance and time measurements are field measurements like any other one: they are part of the **boundary data** of the problem

Boundary values of the gravitational field = geometry of box surface = distance and time separation of measurements

A process and its amplitude

Boundary state $\Psi = \psi_{in} \otimes \psi_{out}$

Amplitude $A = W(\Psi)$



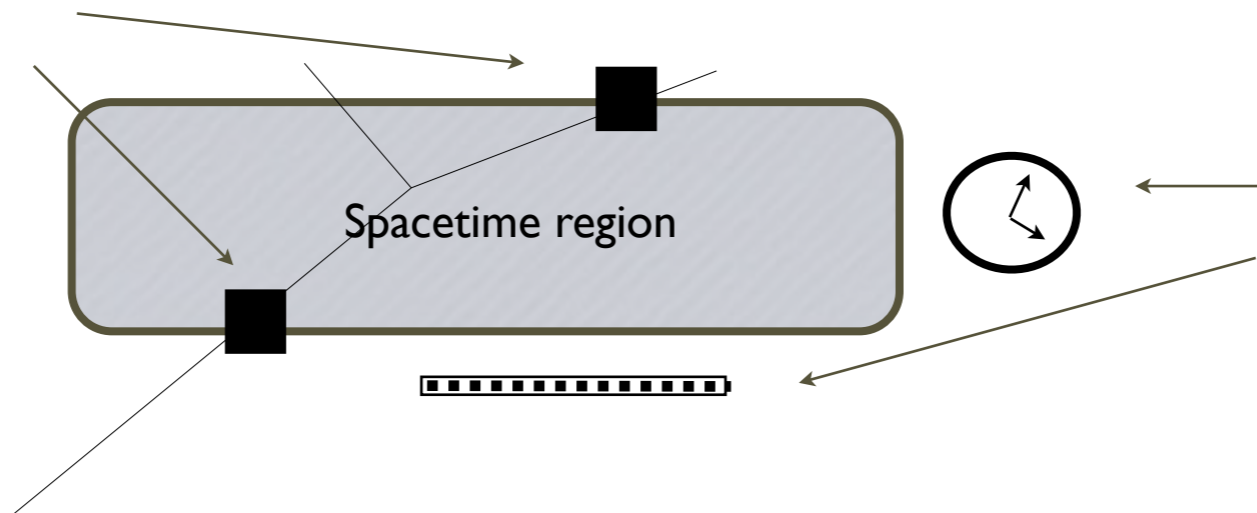
**Quantum system
=
Spacetime region**

→ Hamilton function: $S(q,t,q',t')$

In GR, distance and time measurements are field measurements like any other one: they are part of the **boundary data** of the problem

Boundary values of the gravitational field = geometry of box surface = distance and time separation of measurements

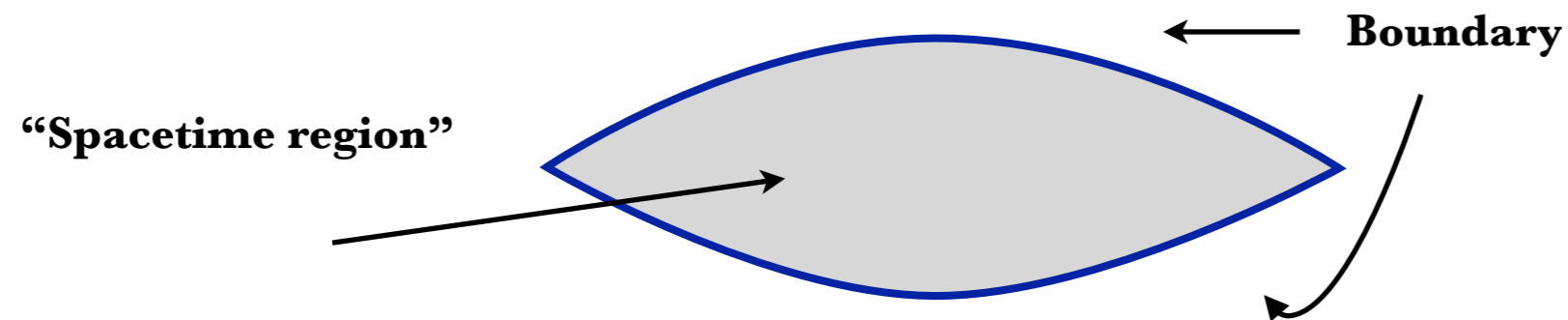
Particle detectors
= field measurements



Distance and time measurements
= gravitational field measurements

GR: localization of a system is relative to other systems

QM: events for system isare relative to other systems



Understand the world as a net of system influencing one another

1. We do have quantum gravity theories

(We do not know which is right)

2. We do have have some empirical evidence

(Rule out and/or disfavour some theories, but no positive support yet)

3. Results towards issues of philosophical relevance

i: nature of physical space

(“material”, discrete)

ii: nature of physical time

(absent at the fundamental level)

iii: relation between physical and experiential time

(Time is a multilayer concept, Experiential time is understood in terms of thermodynamics, biology, brain sciences.)

iv: connection between relational aspects of GR and QM

(Understand the world as a net of system influencing one another.)